

Graphene plasmonics

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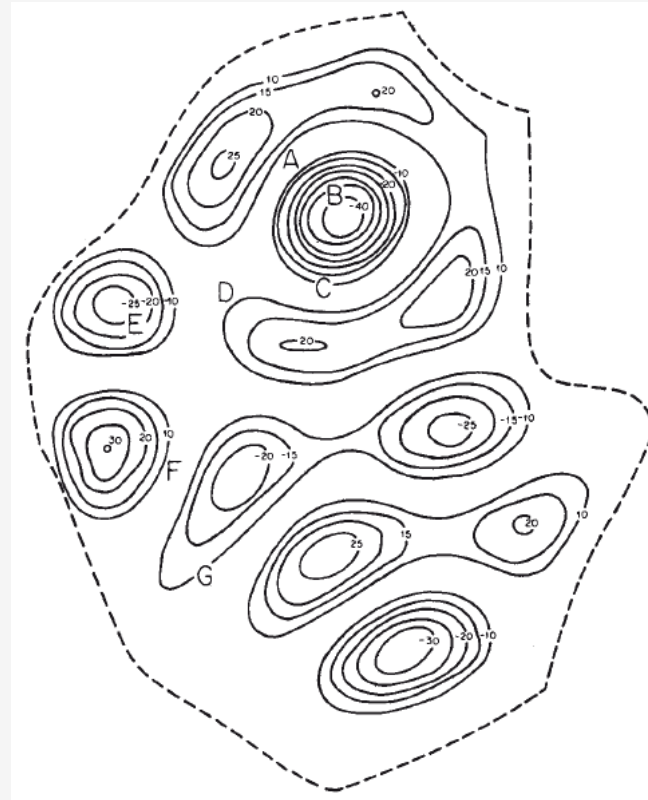
**New Materials and New Concepts
for Controlling Light and Waves**

3 - 7 October 2012
The Hong Kong University of Science and Technology

A collage of scientific images including a diffraction pattern, a 3D structure of a material, a graph showing a red line on a grid, and a microscopic view of a device.

Surface plasmons are surface waves ...



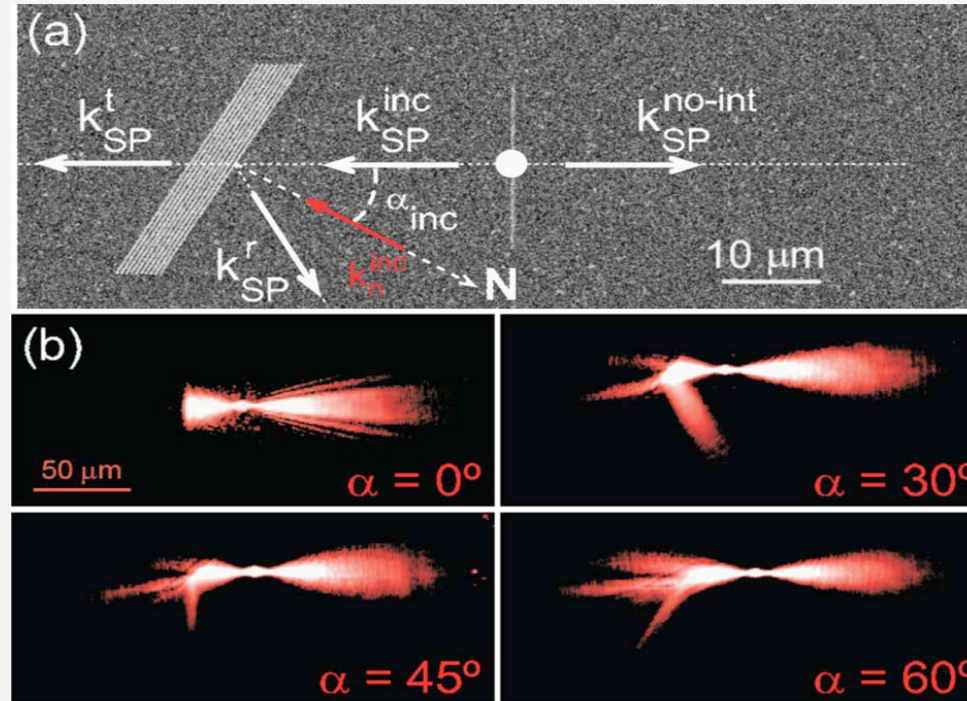


Flores *et al.*, Nature (1987)

Level of damage in the 1985 earthquake in Mexico city

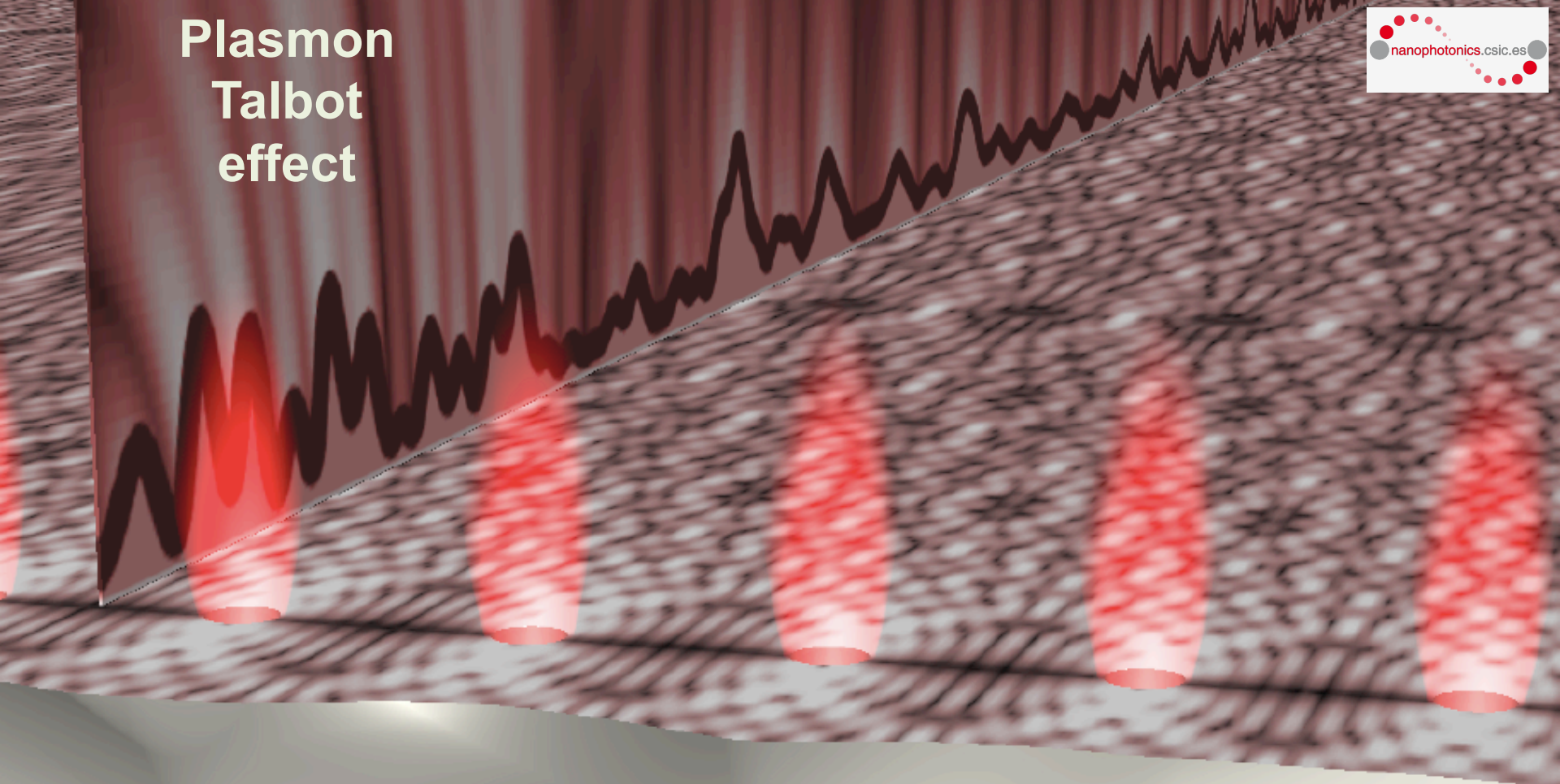
Surface plasmons are surface waves
involving collective electron motion
and propagating on metal surfaces ...

Plasmon Bragg mirrors

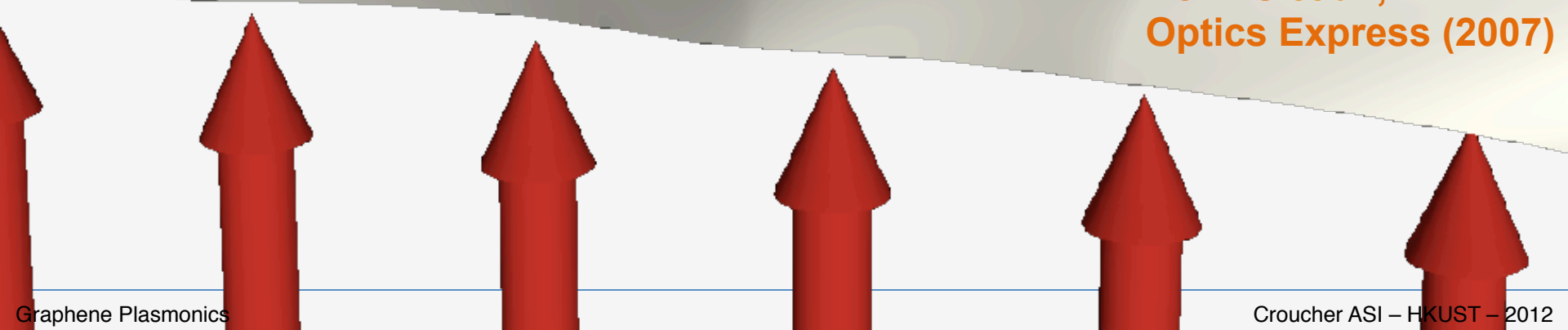


González, ..., Dereux, Quidant, Krenn, Opt. Lett. (2007)

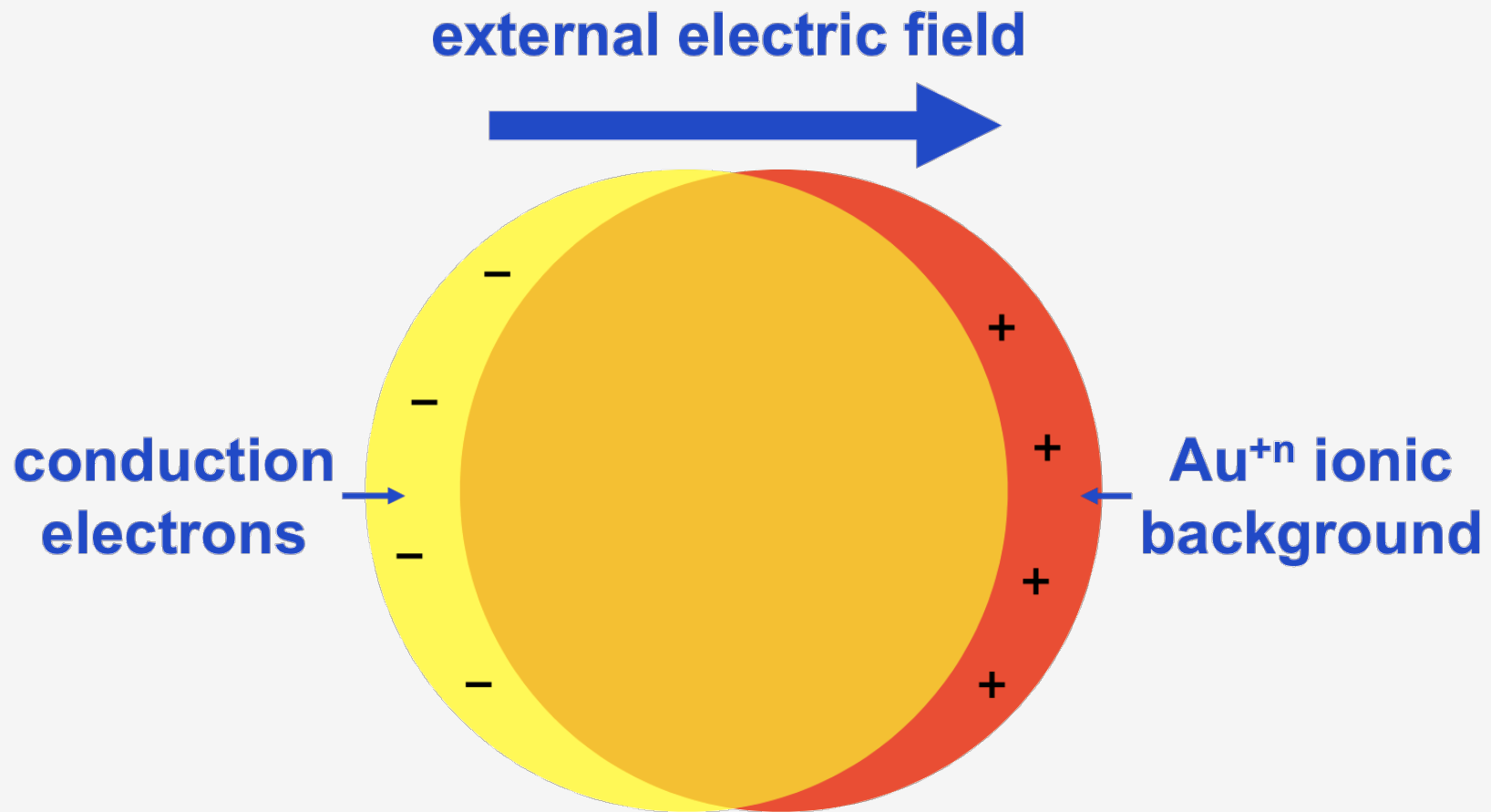
Plasmon Talbot effect



Dennis *et al.*,
Optics Express (2007)



Surface plasmons are surface waves involving collective electron motion and propagating on metal surfaces or localized in metal (nano)structures (e.g., nanoparticles), where they couple efficiently to light ...



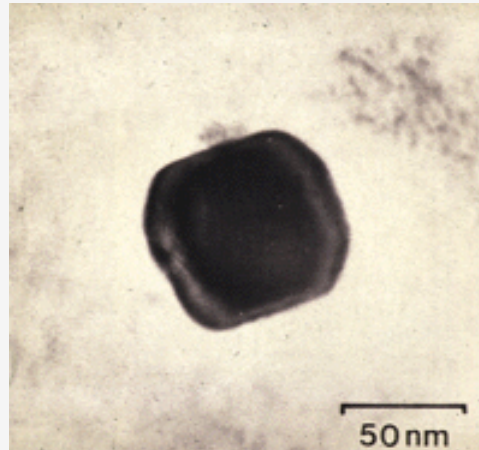
Plasmons in metallic nanoparticles

Romans played empirically with nanoparticle plasmons: the Licurgo cup dating from the IV century

In reflection



In transmission

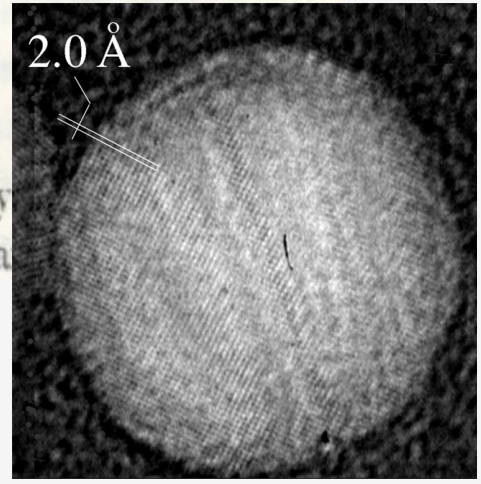
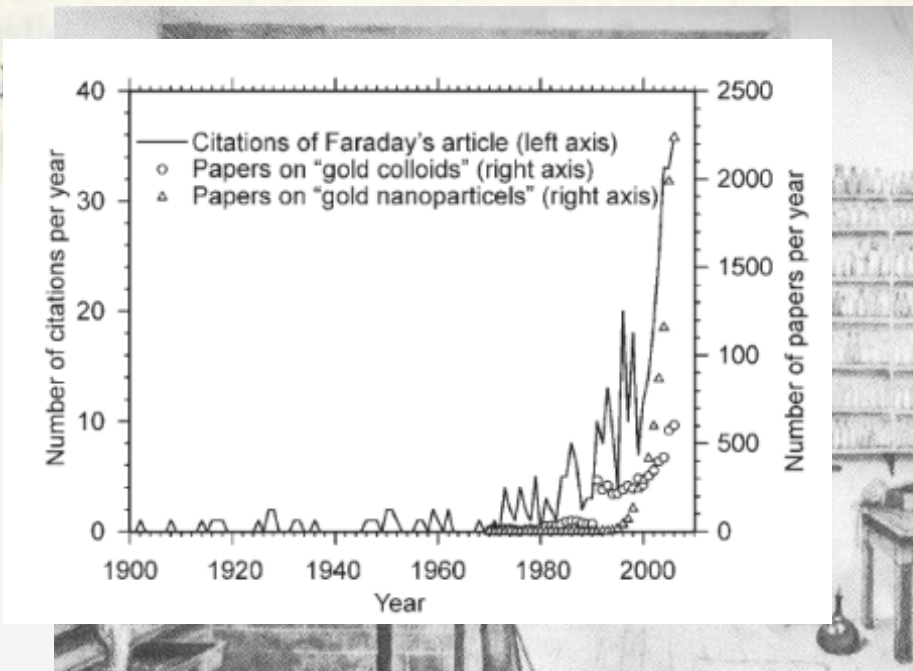


An electron microscope image shows 70-nm
Au-Ag nanoparticles inside the glass

Plasmons in metallic nanoparticles

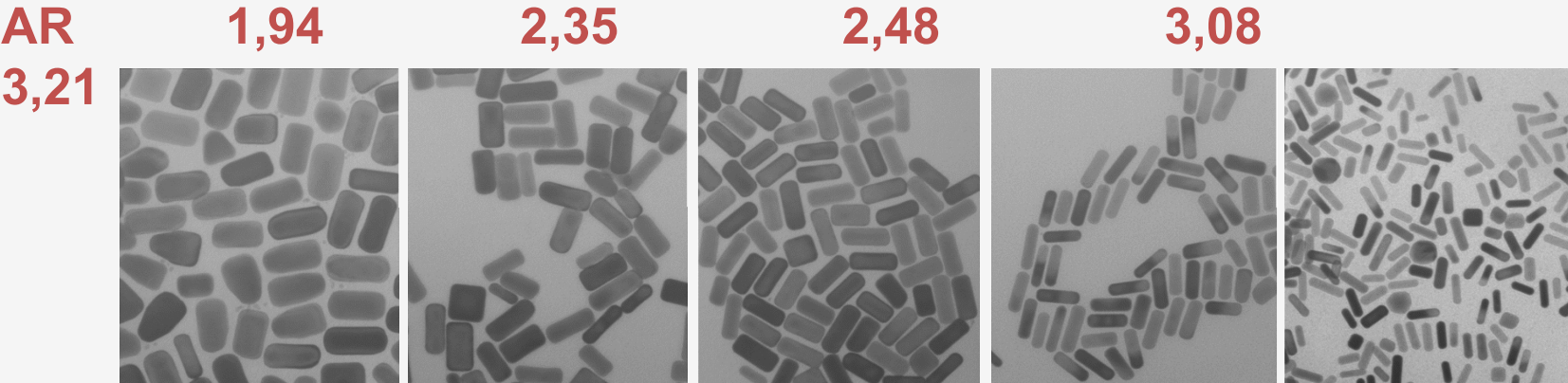
X. THE BAKERIAN LECTURE.—*Experimental Relations of Gold (and other Metals) to Light.* By MICHAEL FARADAY, Esq., D.C.L., F.R.S., Fullerian Prof. Chem. Royal Institution, Foreign Associate of the Acad. Sciences, Paris, Ord. Boruss. pour le Mérite, Eq., Memb. Royal and Imp. Acadd. of Sciences, Petersburg, Florence, Copenhagen, Berlin, Göttingen, Modena, Stockholm, Munich, Bruxelles, Vienna, Bologna, Commander of the Legion of Honour, &c. &c.

Received November 15, 1856,—Read February 5, **1857**

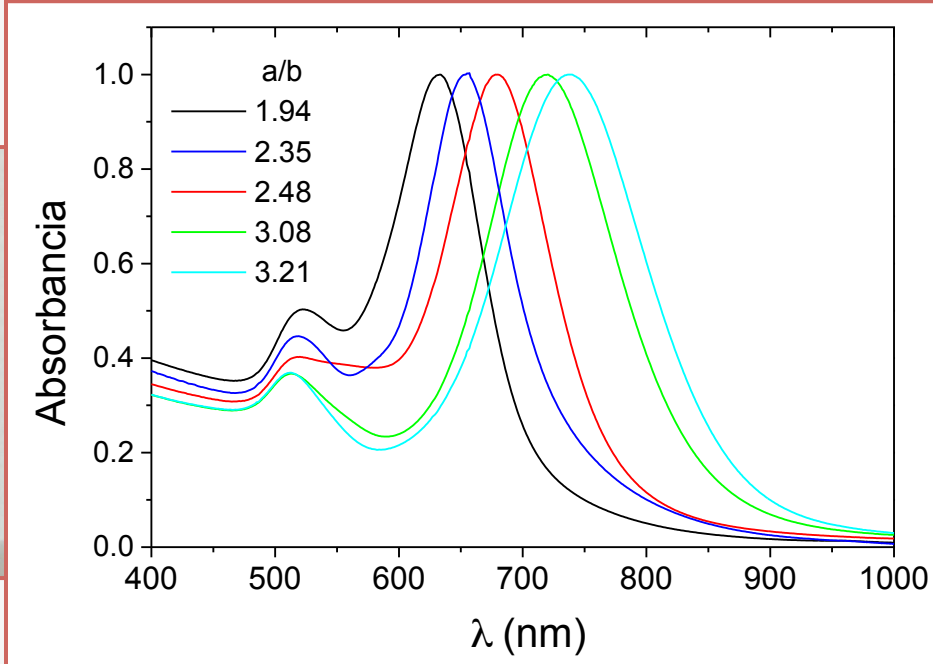


Plasmons in metallic nanoparticles

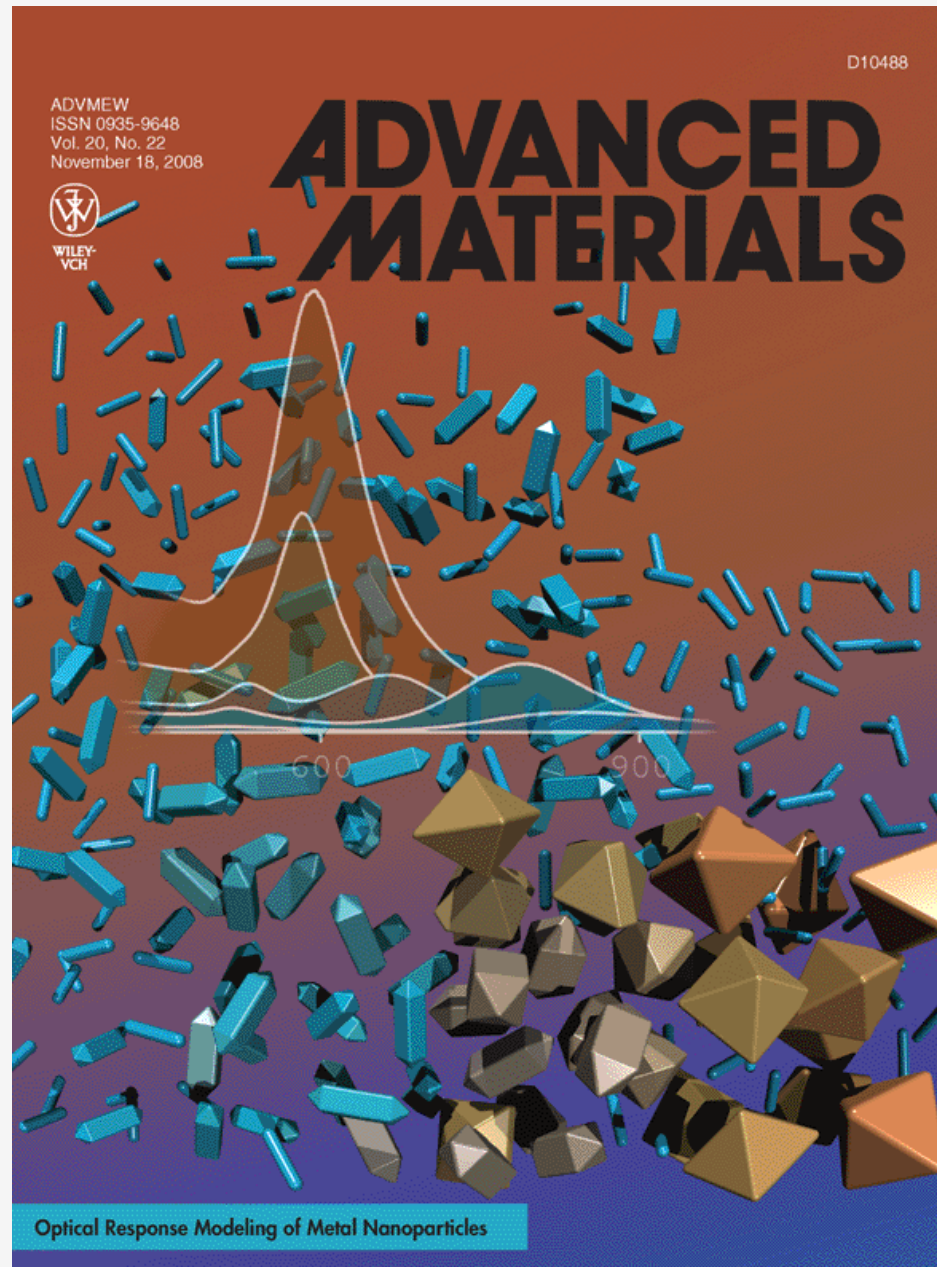
The colors of gold nanorods



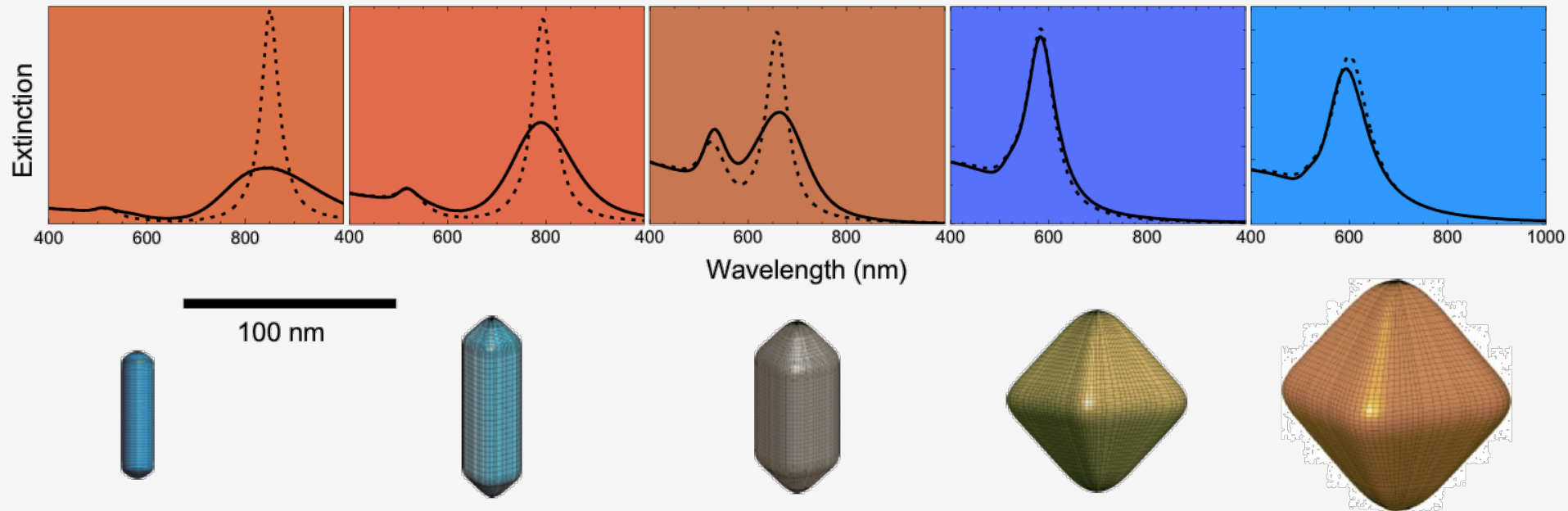
100 nm



Pérez-Juste *et al.*, *Appl. Surf. Sci.* (2004)



Artificial colors through tailored plasmons in nanoparticles

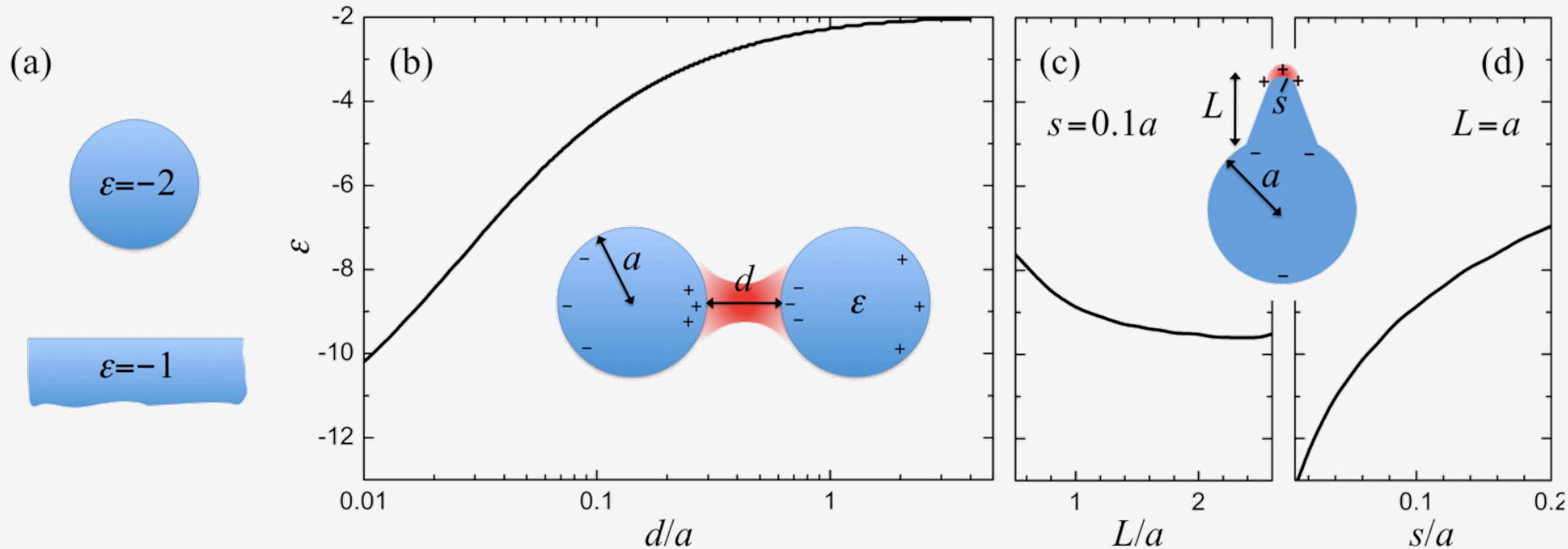


Myroshnychenko *et al.*, *Advanced Materials* (2008)

Surface plasmons are surface waves involving collective electron motion and propagating on metal surfaces or localized in metal (nano)structures (e.g., nanoparticles), where they couple efficiently to light, **they produce strong confinement of the electromagnetic field (size \ll wavelength) ...**

Why do we need metals?

Plasmons in the long wavelength limit (Poisson equation) are scale-invariant, and therefore, they exist for structures down to a few nm.



Localized excitations require negative permittivity

Álvarez-Puebla *et al.*, J. Phys. Chem. Lett. (2010)

Maxwell equations ...

$$\mathbf{E}(\mathbf{r}, t) = \mathbf{E}(\mathbf{r}, \omega) e^{-i\omega t} + \mathbf{E}^*(\mathbf{r}, \omega) e^{i\omega t}$$

$$\nabla \cdot \varepsilon(\mathbf{r}, \omega) \mathbf{E}(\mathbf{r}, \omega) = 0$$

$$\nabla \times \mathbf{E}(\mathbf{r}, \omega) = i \frac{\omega}{c} \mathbf{B}(\mathbf{r}, \omega)$$

$$\nabla \cdot \mathbf{B}(\mathbf{r}, \omega) = 0$$

$$\nabla \times \mathbf{H}(\mathbf{r}, \omega) = -i \frac{\omega}{c} \varepsilon(\mathbf{r}, \omega) \mathbf{E}(\mathbf{r}, \omega)$$

Why do we need metals?

Maxwell equations for small particles

$$\mathbf{E}(\mathbf{r}, t) = \mathbf{E}(\mathbf{r}, \omega) e^{-i\omega t} + \mathbf{E}^*(\mathbf{r}, \omega) e^{i\omega t}$$

$$\nabla \cdot \varepsilon(\mathbf{r}, \omega) \mathbf{E}(\mathbf{r}, \omega) = 0 \quad \nabla \times \mathbf{E}(\mathbf{r}, \omega) = 0$$

$$\nabla \cdot \mathbf{B}(\mathbf{r}, \omega) = 0 \quad \nabla \times \mathbf{H}(\mathbf{r}, \omega) = 0$$

Electricity and magnetism are decoupled
in the long-wavelength ($c \rightarrow \infty$) limit*

*Except in regions of very index of refraction, $|n|a \sim \lambda$

Maxwell equations for small particles

$$\mathbf{E}(\mathbf{r}, t) = \mathbf{E}(\mathbf{r}, \omega) e^{-i\omega t} + \mathbf{E}^*(\mathbf{r}, \omega) e^{i\omega t}$$

$$\nabla \cdot \varepsilon(\mathbf{r}, \omega) \mathbf{E}(\mathbf{r}, \omega) = 0 \quad \nabla \times \mathbf{E}(\mathbf{r}, \omega) = 0$$

$$\mathbf{H} = 0$$

Maxwell equations for small particles

$$\mathbf{E}(\mathbf{r}, t) = \mathbf{E}(\mathbf{r}, \omega) e^{-i\omega t} + \mathbf{E}^*(\mathbf{r}, \omega) e^{i\omega t}$$

$$\nabla \cdot \varepsilon(\mathbf{r}, \omega) \mathbf{E}(\mathbf{r}, \omega) = 0 \quad \mathbf{E}(\mathbf{r}, \omega) = -\nabla \phi(\mathbf{r}, \omega)$$

$$\nabla \cdot \varepsilon(\mathbf{r}, \omega) \nabla \phi(\mathbf{r}, \omega) = 0$$

The Poisson equation also describes stationary heat transport:

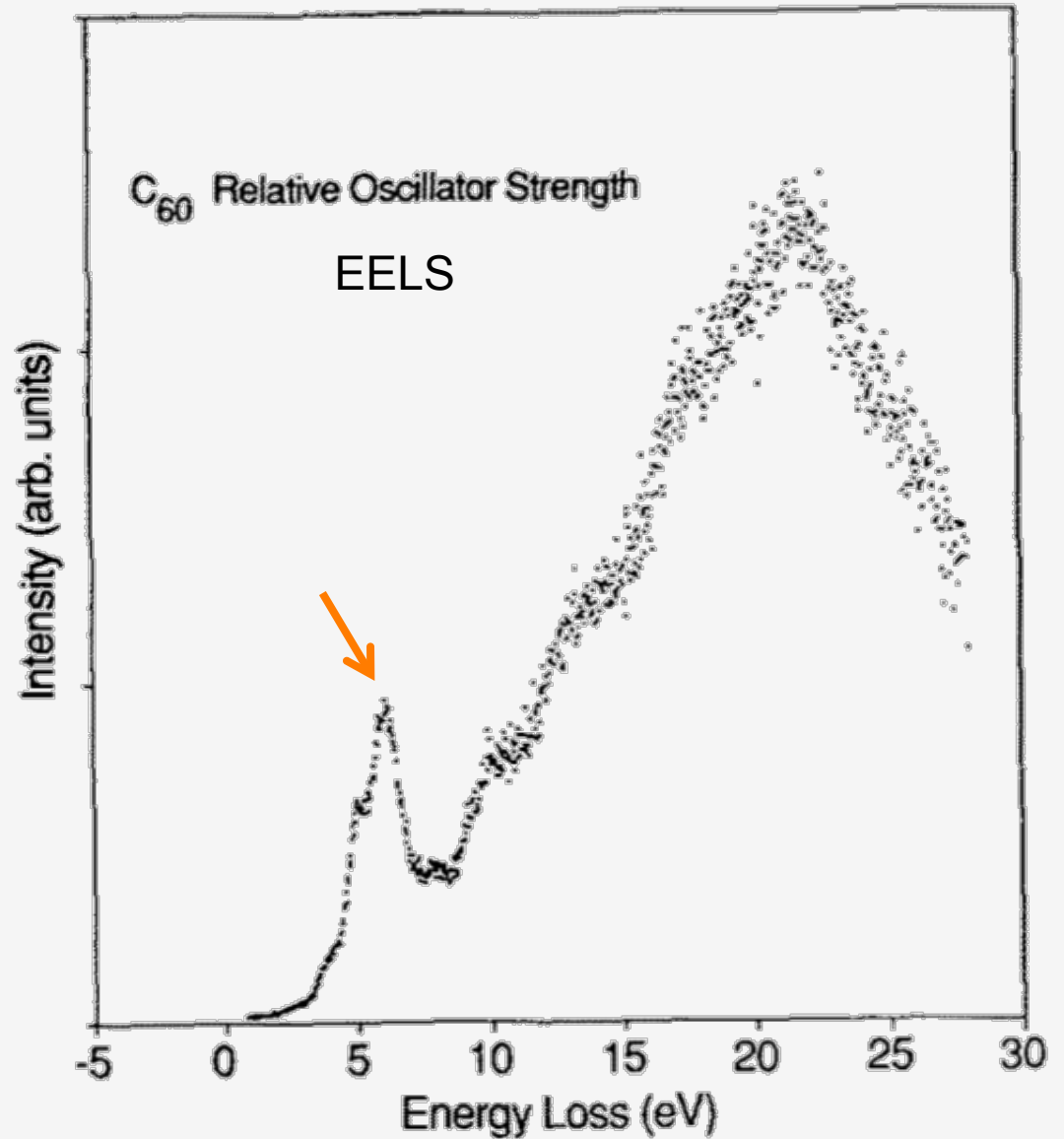
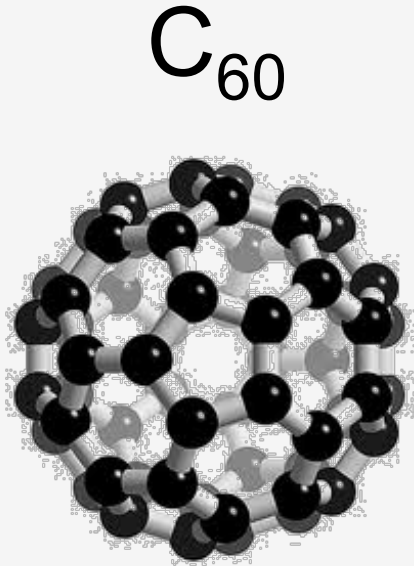
$\varepsilon \rightarrow k$, thermal conductivity

$\phi \rightarrow$ temperature

Thermodynamics:

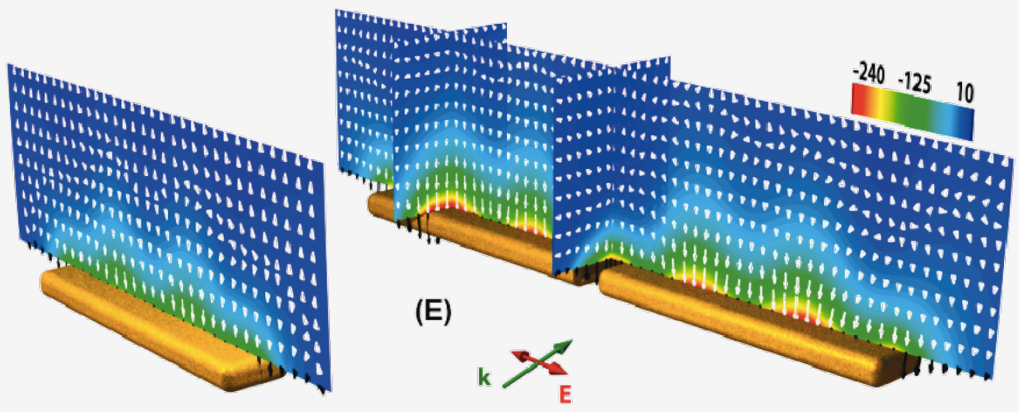
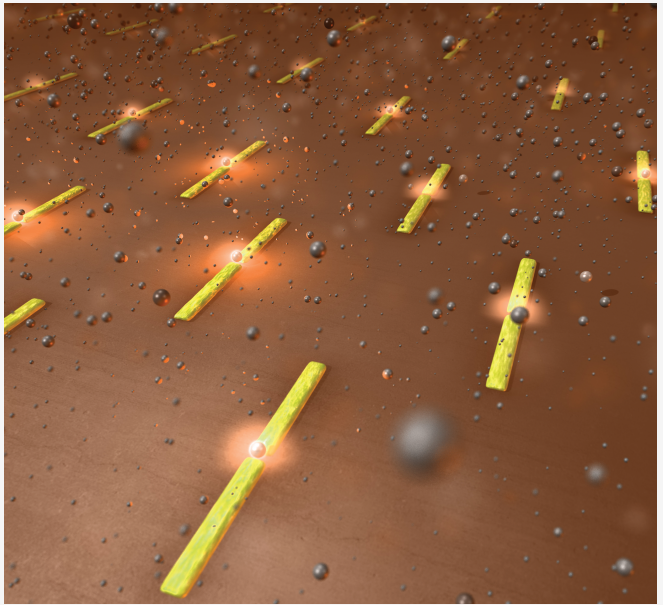
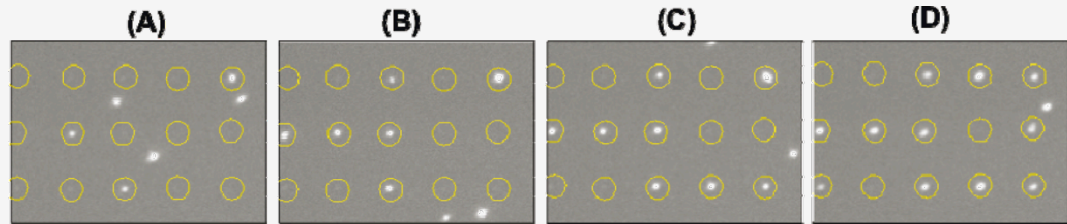
- flow towards lower temperature regions $\rightarrow k > 0$
- absence of trapped thermal energy

But ... do we really need metals?



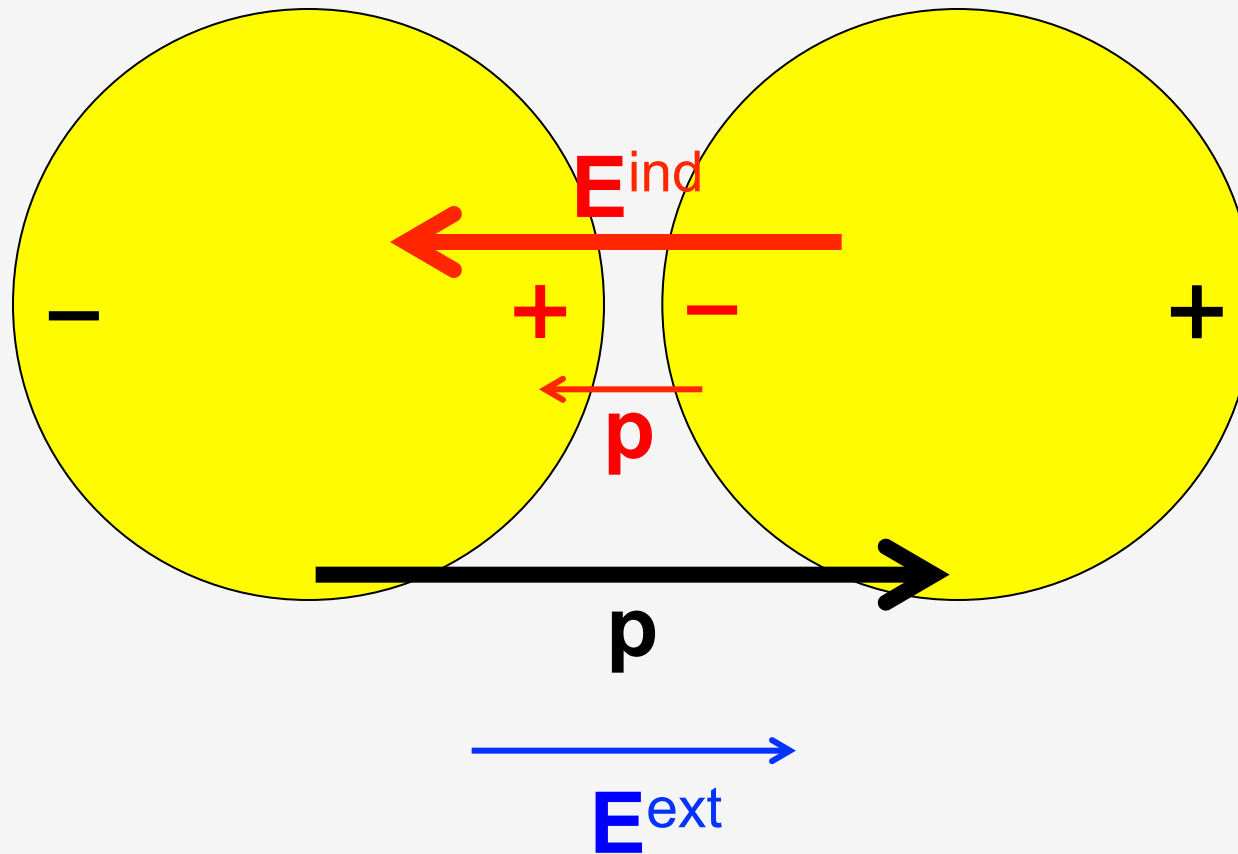
Keller and Coplan, Chem. Phys. Lett. (1992)

Optical trapping - nanotweezers



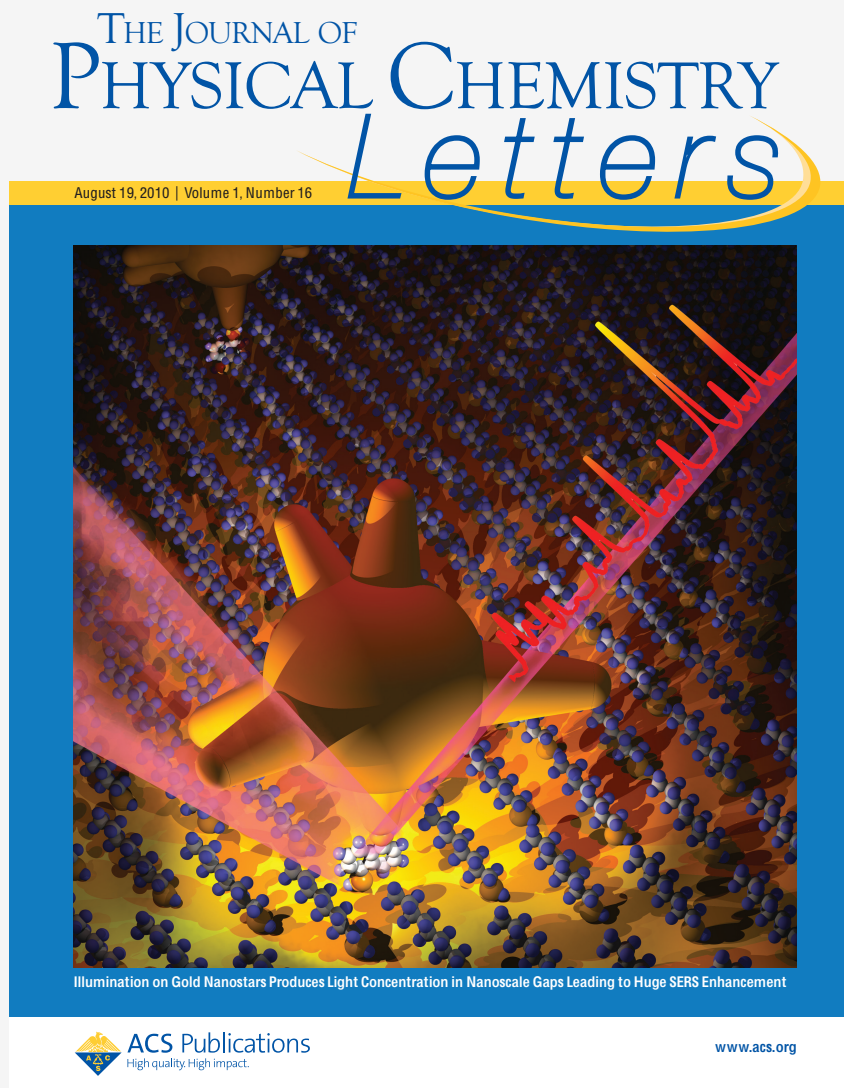
Nanoparticle Trapping with Resonant Optical Antennas

Surface plasmons are surface waves involving collective electron motion and propagating on metal surfaces or localized in metal (nano)structures (e.g., nanoparticles), where they couple efficiently to light, they produce strong confinement of the electromagnetic field (size \ll wavelength), **and they generate huge enhancement of the optical electric-field intensity.**



Charge neutrality \rightarrow strong coupling to light through p , strong enhancement in the gap

Controlled 10^{10} SERS enhancement \rightarrow 10^5 intensity enhancement



Álvarez-Puebla,
Liz-Marzán, G. de Abajo,
JACS (2009), JPCL (2010)

Plasmon simulation has become a simple task in most systems of current interest, for example to understand confined gap plasmons in nanoparticle dimers (e.g., check the widgets of our website for this and other applications).

Dimer widget at <http://www.nanophotonics.es>

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Light scattering by a sphere dimer: Extinction cross section

Contact [F. J. García de Abajo](#) for details

Extinction (= light scattering + absorption) cross section of a sphere dimer for polarization along the dimer axis, calculated from MESME (see [F. J. García de Abajo](#), *Phys. Rev. B* 60, 6086 (1999)) up to an iteration order as specified below.

Dielectric function $\epsilon(\omega)$

Drude
 Ag (J & C 1972)
 Al (Palik)
 Al2O3 (Palik)
 Au (J & C 1972)

For Drude:
 $\epsilon(\omega) = 1 - \omega_p^2 / (\omega(\omega + i\eta))$
 ω_p (eV) =
 η (eV) =

nanophotonics.csic.es
Calculate

Light scattering cross section of a dimer (Click on the figure to access data file.)
 $l_{max}=20$, $a=10.00$ nm, $b=10.00$ nm, and $d=1.00$ nm.

Geometrical parameters

Left-sphere radius a (nm) =

Right-sphere radius b (nm) =

Surface-to-surface distance d (nm) =

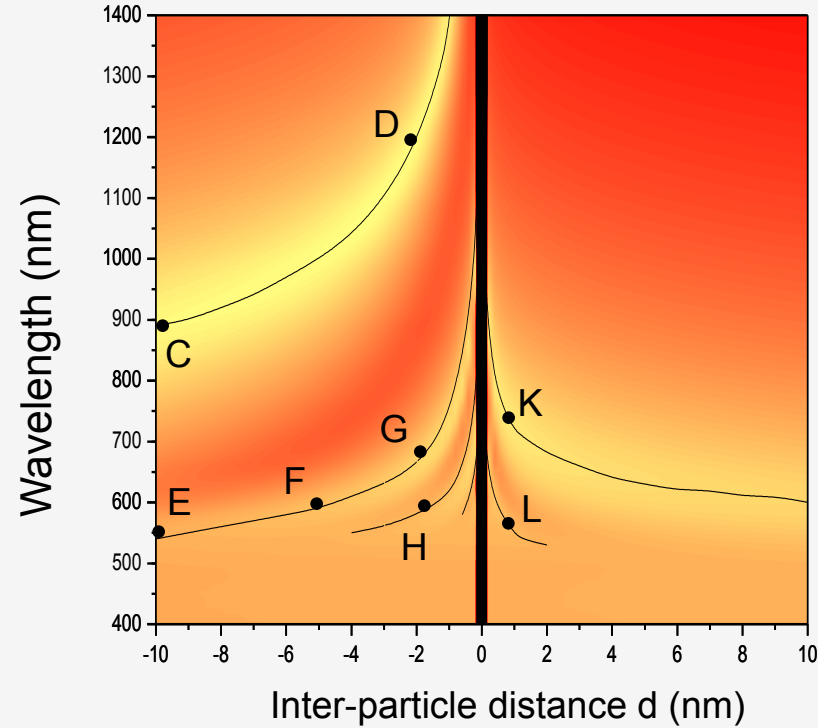
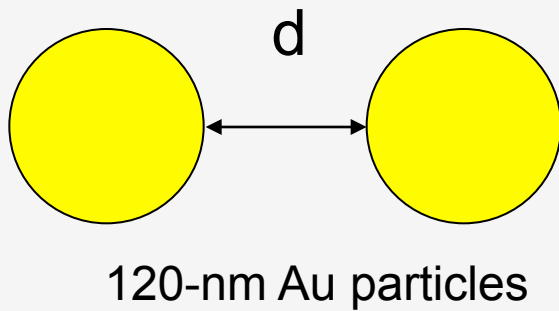
Other parameters

ω_{min} (eV) = ω_{max} (eV) =

l_{max} = (maximum orbital momentum no. ≤ 20)

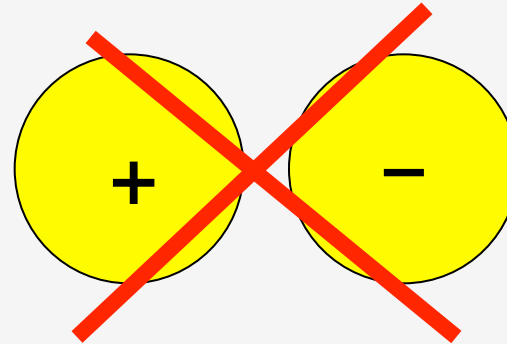
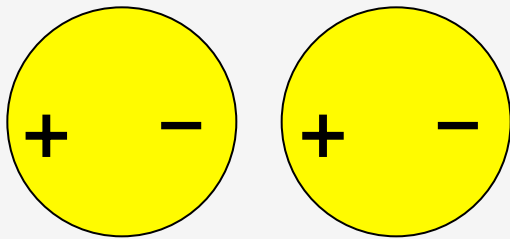
Iteration order =

Particle dimer: transition from touching to non-touching

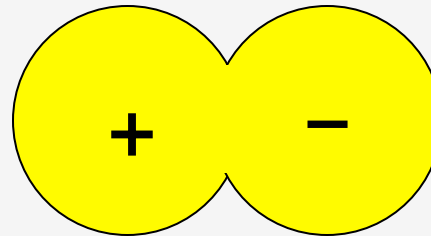
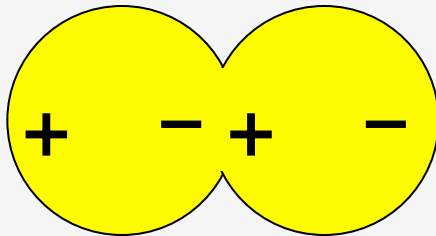


Romero *et al.*, Optics Express (2006)

Particle dimer: transition from touching to non-touching



unphysical

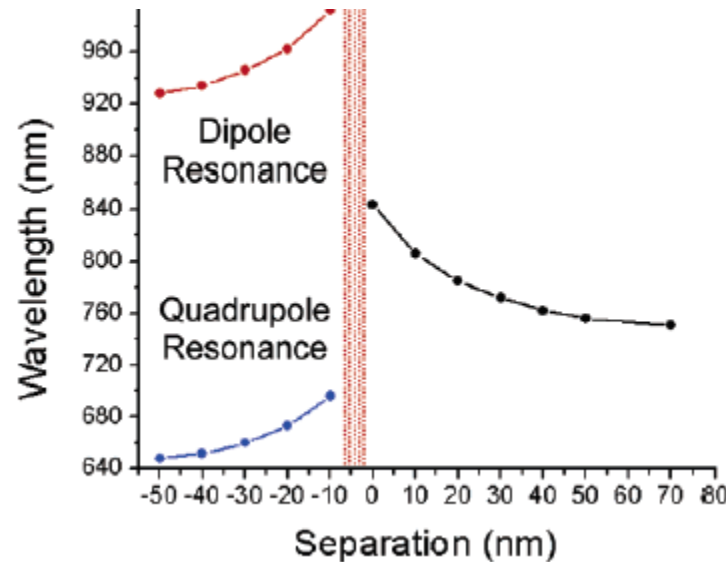
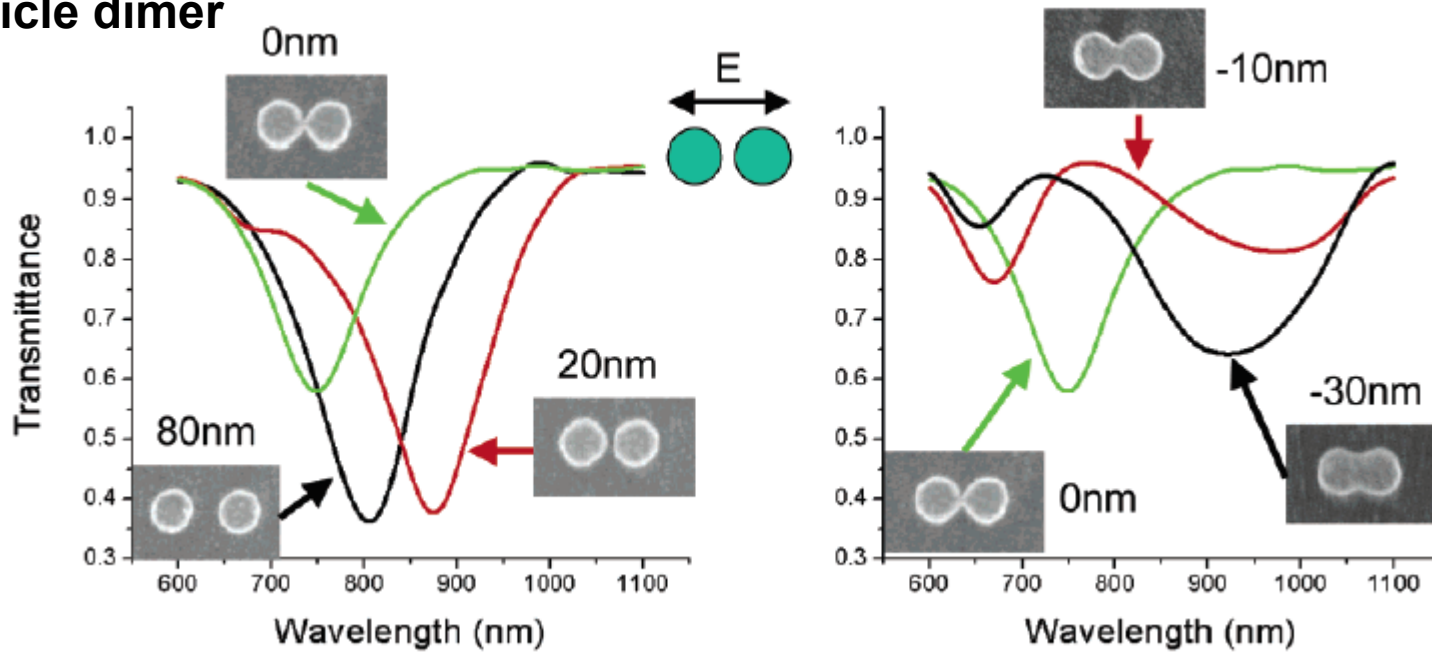


physical

Romero *et al.*, Optics Express (2006)

Transition between touching and non-touching

Particle dimer



Atay *et al.*,
Nano Letters (2004)

Plasmons can be imaged with nanometer precision using electron microscopes via EELS and cathodoluminescence.

See García de Abajo, *Rev. Mod. Phys.* **82**, 209 (2010).

Spectral Imaging of Individual Split-Ring Resonators

Guillaume Boudarham,¹ Nils Feth,² Viktor Myroshnychenko,³ Stefan Linden,^{2,4} Javier García de Abajo,³ Martin Wegener,^{2,4,5} and Mathieu Kociak^{1,*}

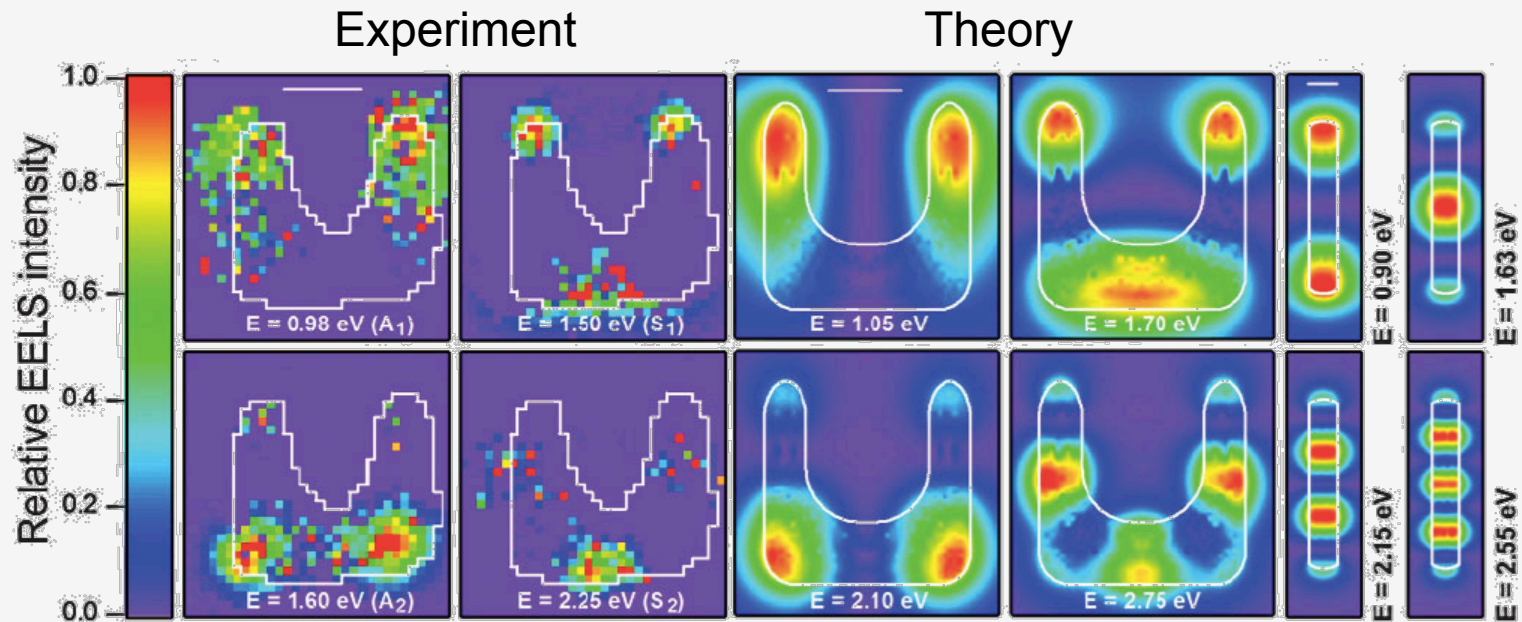
¹Laboratoire de Physique des Solides CNRS/UMR8502, Bâtiment 510, Univ. Paris-Sud, Orsay, 91405, France

²Institut für Angewandte Physik, Universität Karlsruhe, Wolfgang-Gaede-Strasse 1, 76131 Karlsruhe, Germany

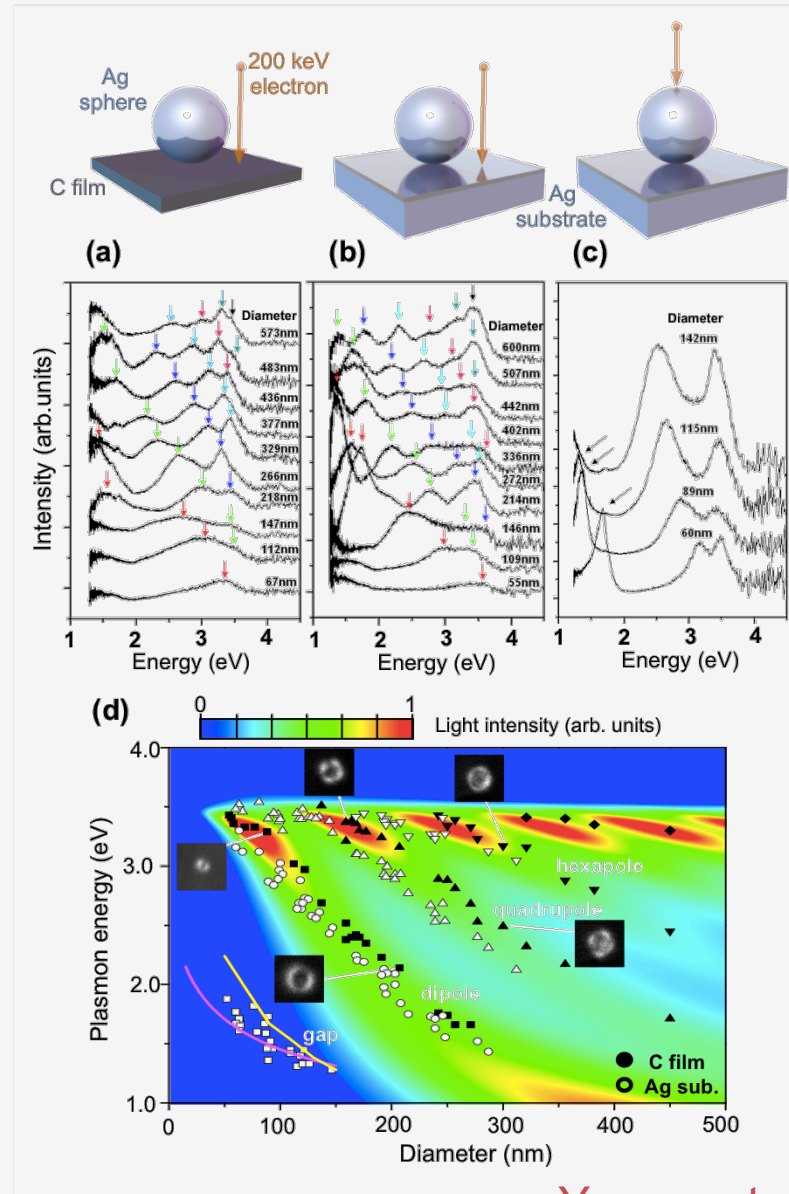
³Instituto de Óptica, CSIC, Serrano 121, 28006 Madrid, Spain

⁴Institut für Nanotechnologie, Karlsruhe Institute of Technology, 76021 Karlsruhe, Germany

⁵DFG-Center for Functional Nanostructures, Karlsruhe Institute of Technology, 76021 Karlsruhe, Germany



Mie and gap plasmons imaged by cathodoluminescence

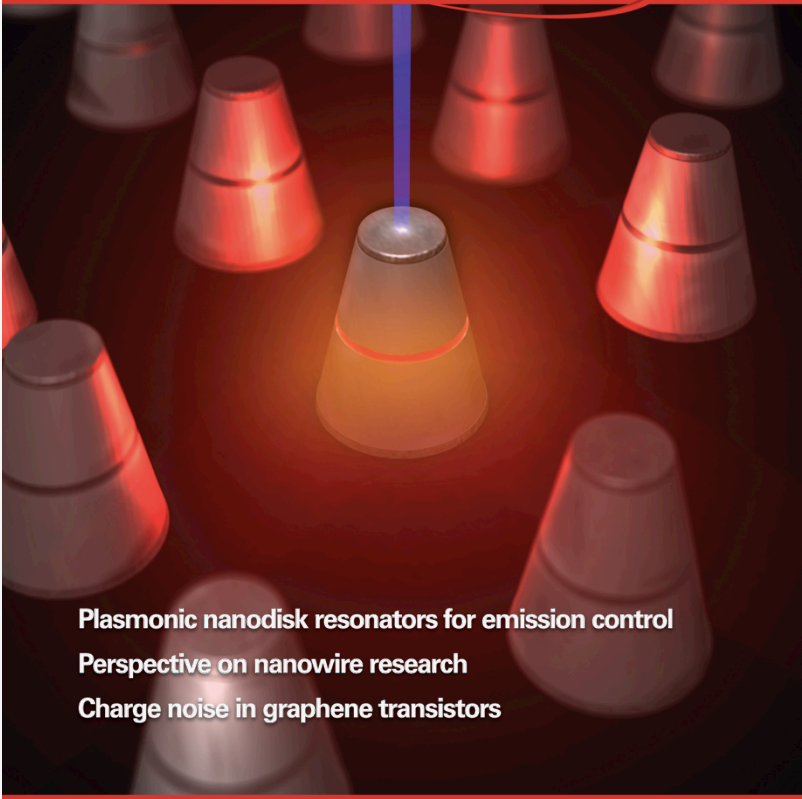


Yamamoto *et al.*, Nano Letters (2010)

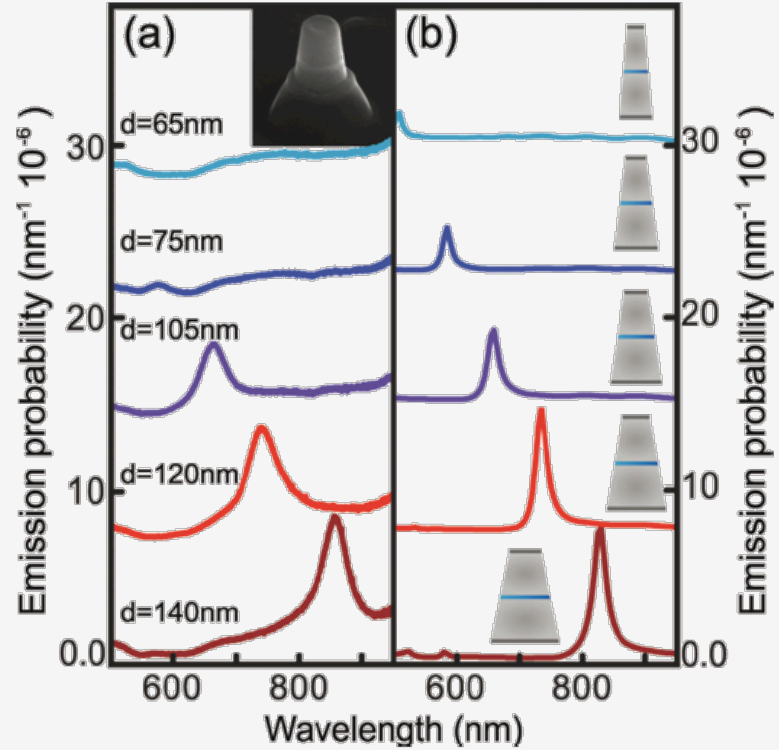
Ultrasmall Mode Volume Plasmonic Nanodisk Resonators

Martin Kuttge,^{*,†} F. Javier García de Abajo,[†] and Albert Polman[†]

NANO LETTERS **10TH**
ANNIVERSARY
 May 2010
 Volume 10, Number 5
pubs.acs.org/Nanolett

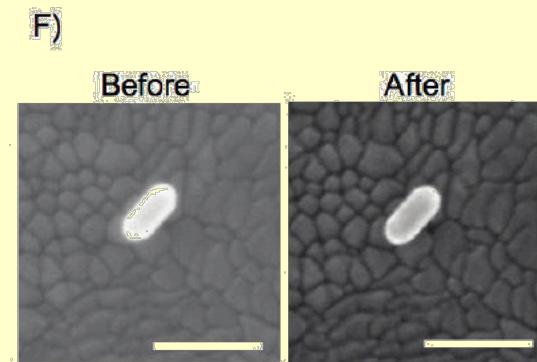
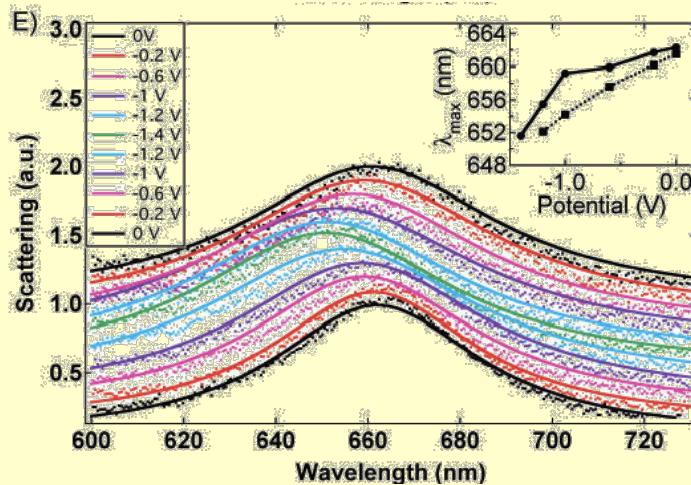


Plasmonic nanodisk resonators for emission control
 Perspective on nanowire research
 Charge noise in graphene transistors



Appealing properties of plasmons

- ✓ ■ Light concentration: plasmon size \ll wavelength
- ✓ ■ Field enhancement: induced field \gg external field
- ✗ ■ Fast tunability: electric doping

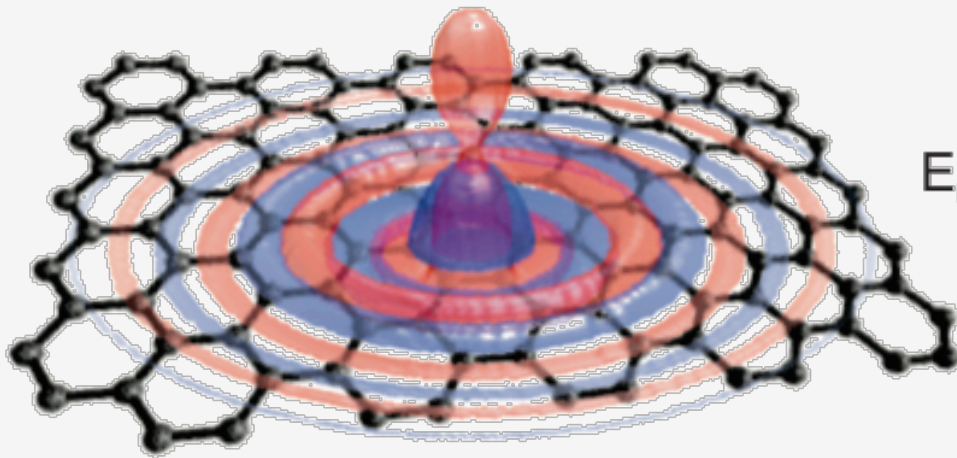


Novo, Funston, Gooding, Mulvaney, JACS (2009)

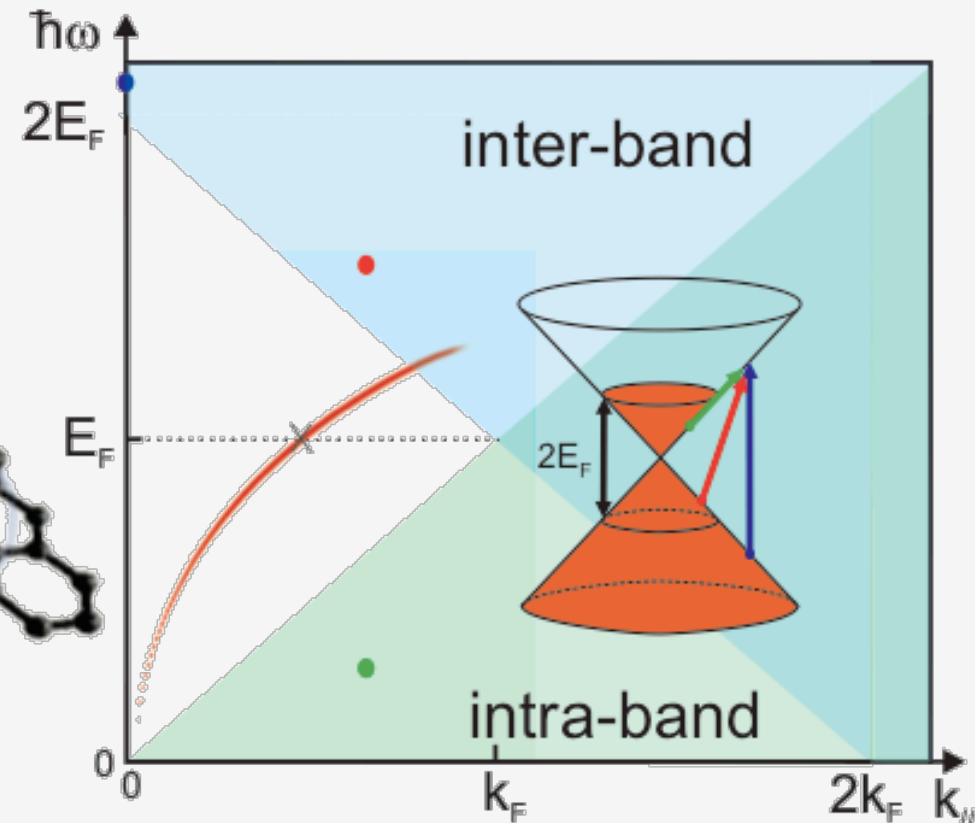
Graphene is a tunable plasmonic material that produces unprecedented confinement and strong light-matter interaction in a robust, solid-state environment

Koppens, Chang, García de Abajo,
Nano Letters **11**, 3370 (2011)

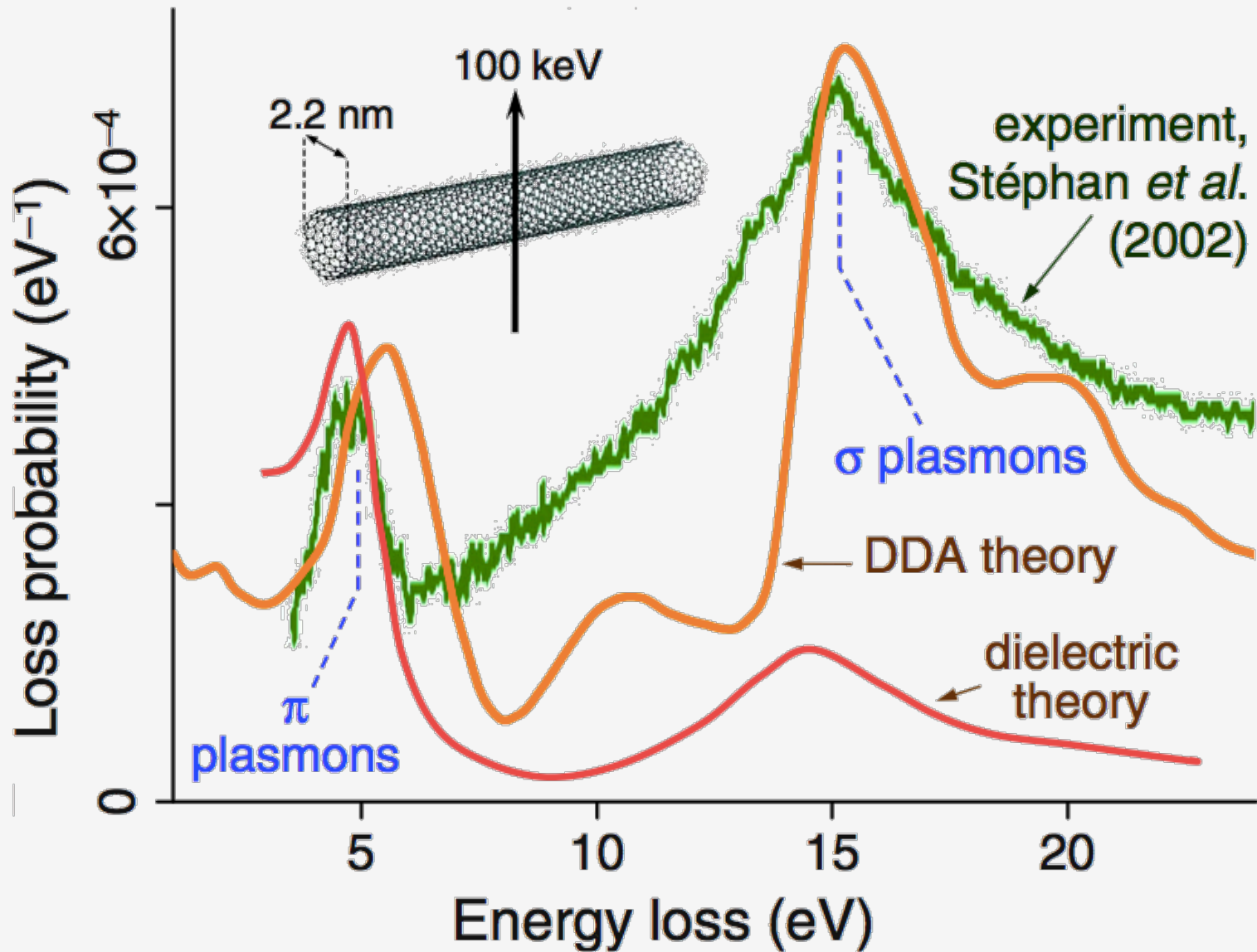
a



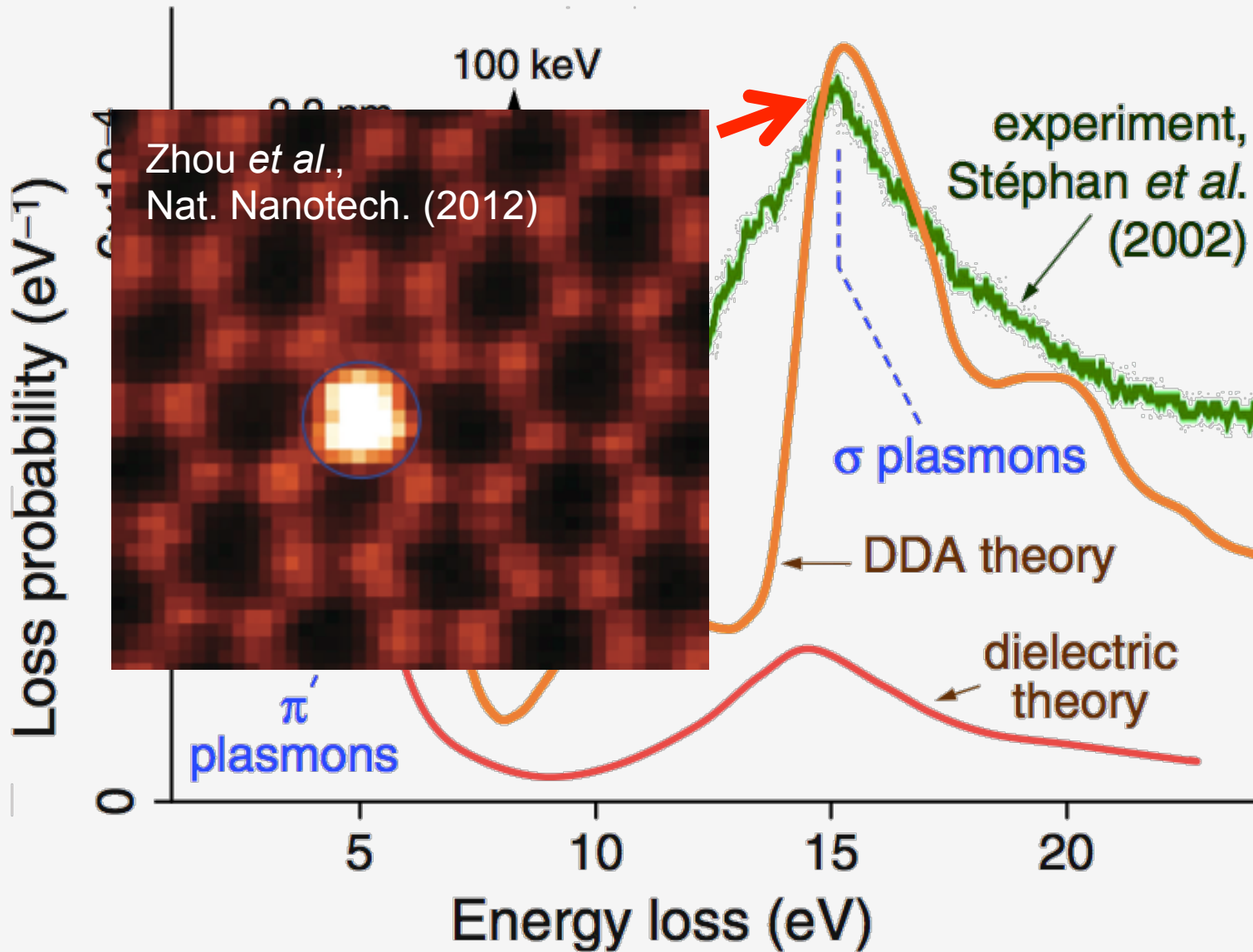
b



Koppens *et al.*, Nano Lett. (2011)

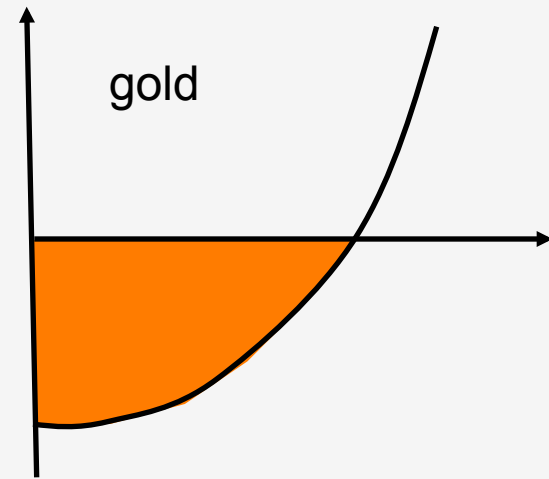
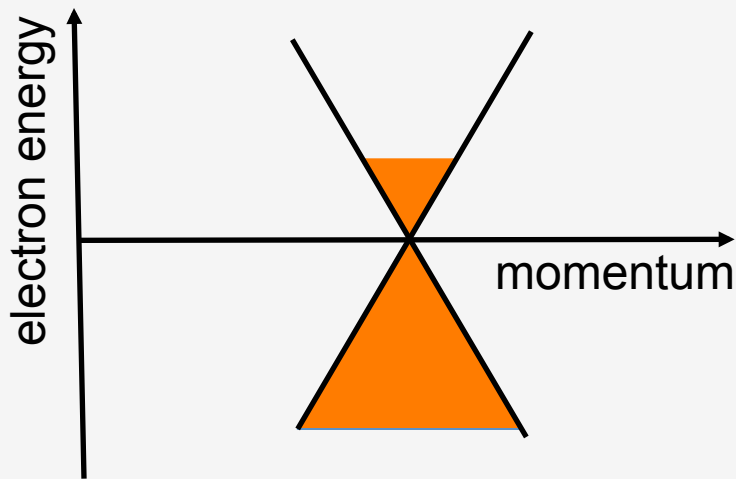
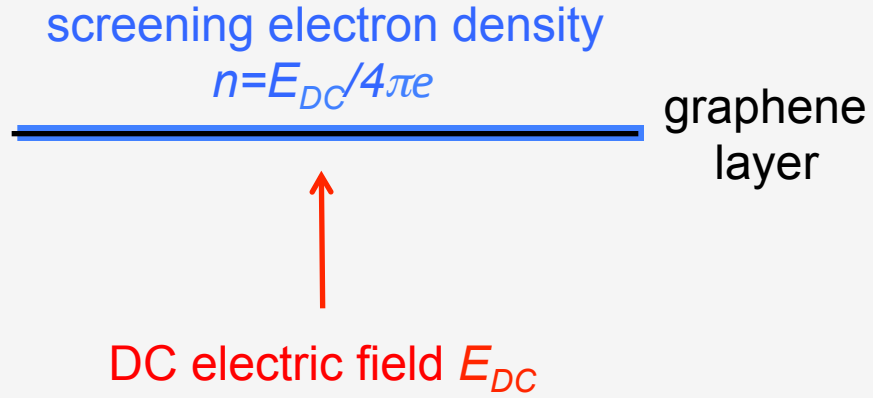


García de Abajo, *Rev. Mod. Phys.* (2010)

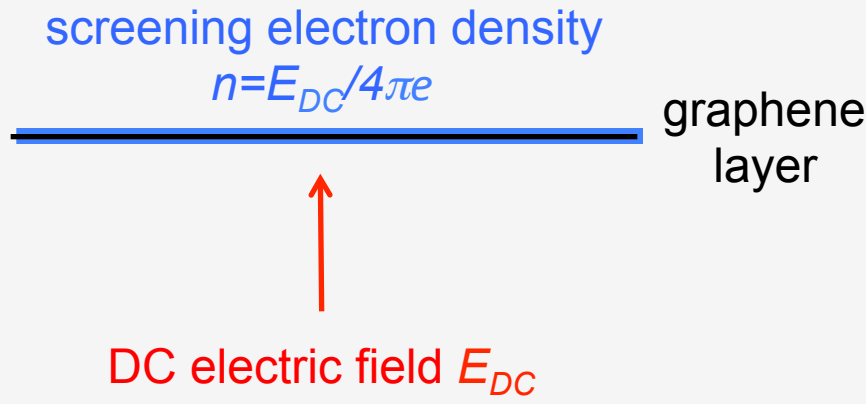


García de Abajo, Rev. Mod. Phys. (2010)

Electrostatic doping of graphene

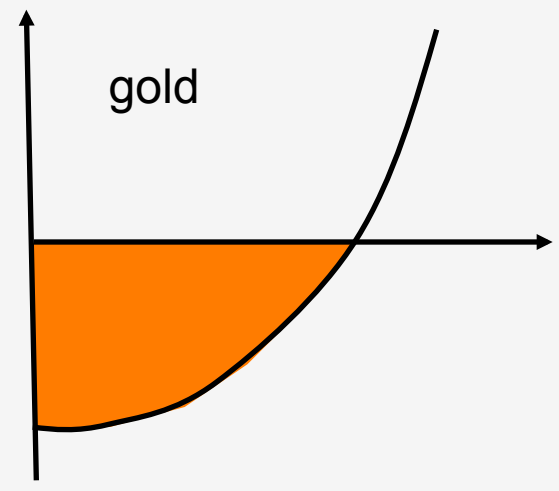


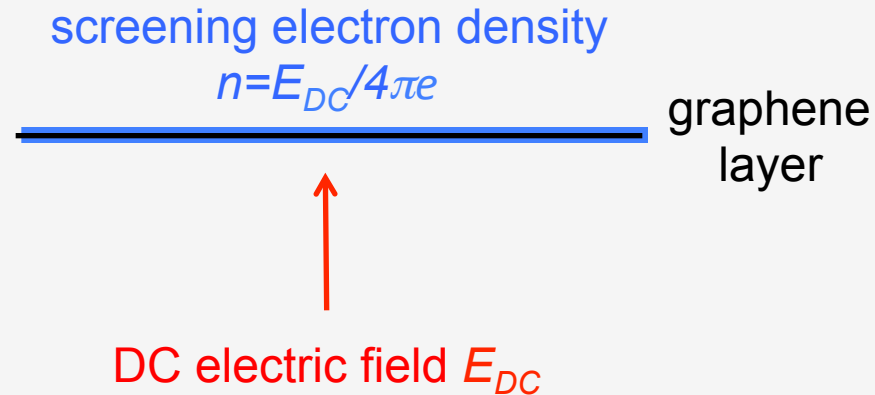
Electrostatic doping of graphene



$$\sigma(\omega) = \frac{ie^2 v_F}{\pi^{3/2} \hbar} \frac{\sqrt{n}}{(\omega + i\tau^{-1})}$$

$$v_F = 10^6 \text{ m / s}$$





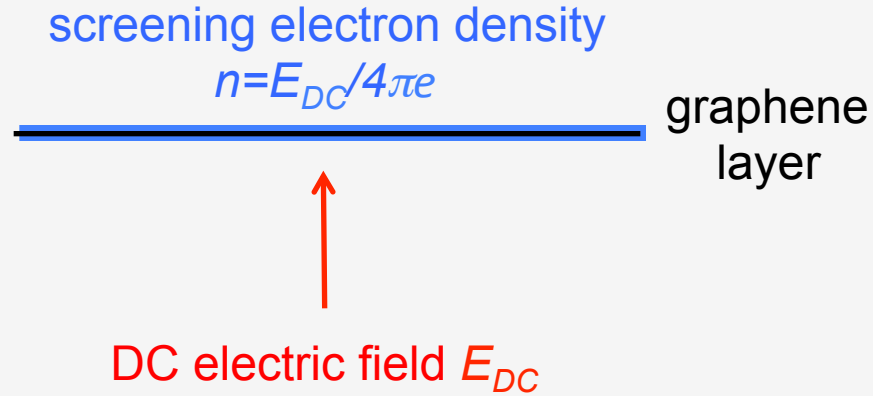
$$\sigma(\omega) = \frac{ie^2 v_F}{\pi^{3/2} \hbar} \frac{\sqrt{n}}{(\omega + i\tau^{-1})}$$

gold

$$v_F = 10^6 \text{ m/s}$$

$$\epsilon(\omega) = 1 - \frac{4\pi e^2}{m} \frac{n}{\omega(\omega + i\tau^{-1})}$$

Electrostatic doping of graphene

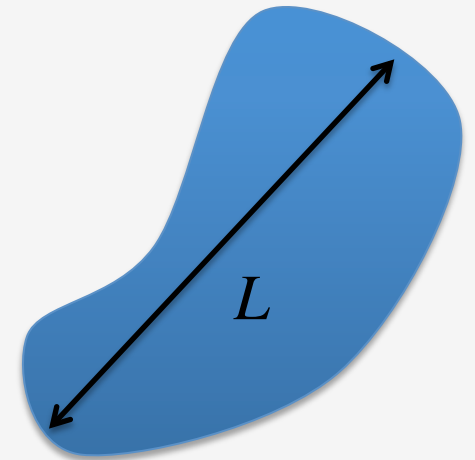


$$\sigma(\omega) = \frac{ie^2 v_F}{\pi^{3/2} \hbar} \frac{\sqrt{n}}{(\omega + i\tau^{-1})}$$

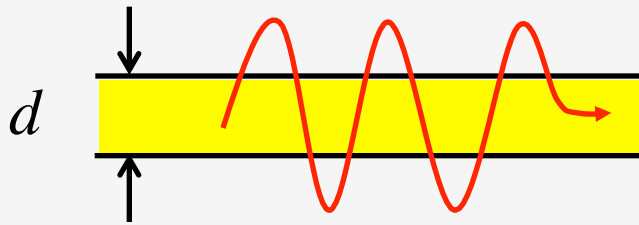
$$v_F = 10^6 \text{ m/s}$$

$$\omega_p = C \frac{E_{DC}^{1/4}}{\sqrt{L}}$$

$$a \ll \lambda$$



Graphene vs gold



$$\epsilon(\omega) \approx \frac{4\pi i \sigma}{\omega d}$$

$$\lambda_{sp} = \frac{-2\pi i \sigma}{c} \lambda$$

1 ML of gold
($d=0.24$ nm)

graphene
($d=0.3$ nm)

$$\sigma \approx \frac{\omega_p^2 d}{4\pi} \frac{i}{\omega + i\tau^{-1}}$$

$$\sigma \approx \frac{e^2 E_F}{\hbar^2 \pi} \frac{i}{\omega + i\tau^{-1}}$$

$$\tau = 1.0 \times 10^{-14} \text{ s}$$

$$\tau = 0 \quad \lambda_{sp} = 2\alpha \frac{E_F}{\hbar\omega} \lambda_0 \quad \text{eV}$$

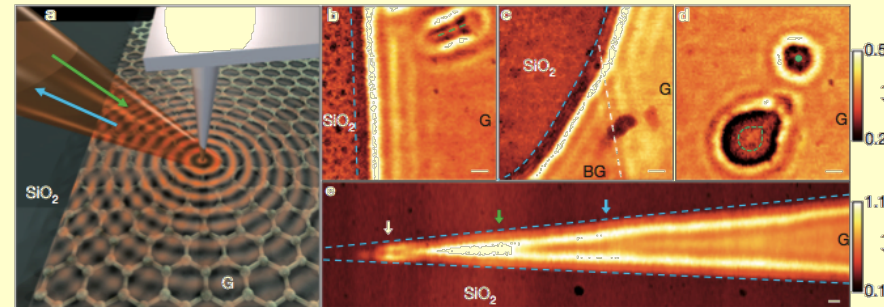
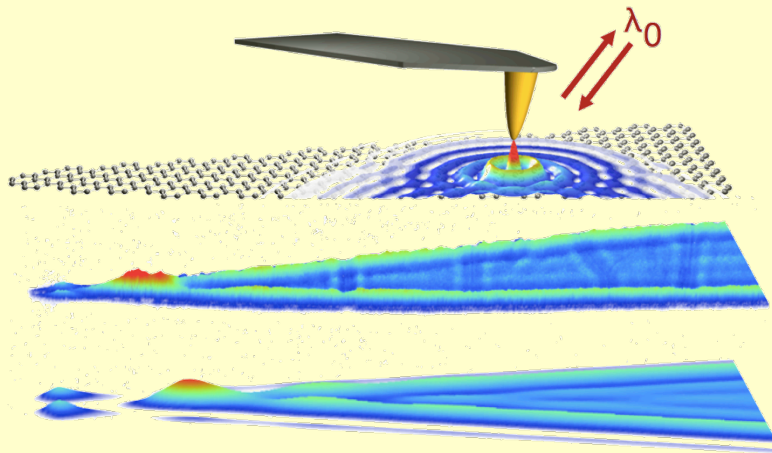
$$\hbar\omega = 0.5 \text{ eV} \rightarrow \lambda_{sp} = 240 \text{ nm}$$

decay length: $\sim (\omega\tau / 2\pi) \lambda_{sp}$

$$1.2 \lambda_{sp}$$

$$60 \lambda_{sp}$$

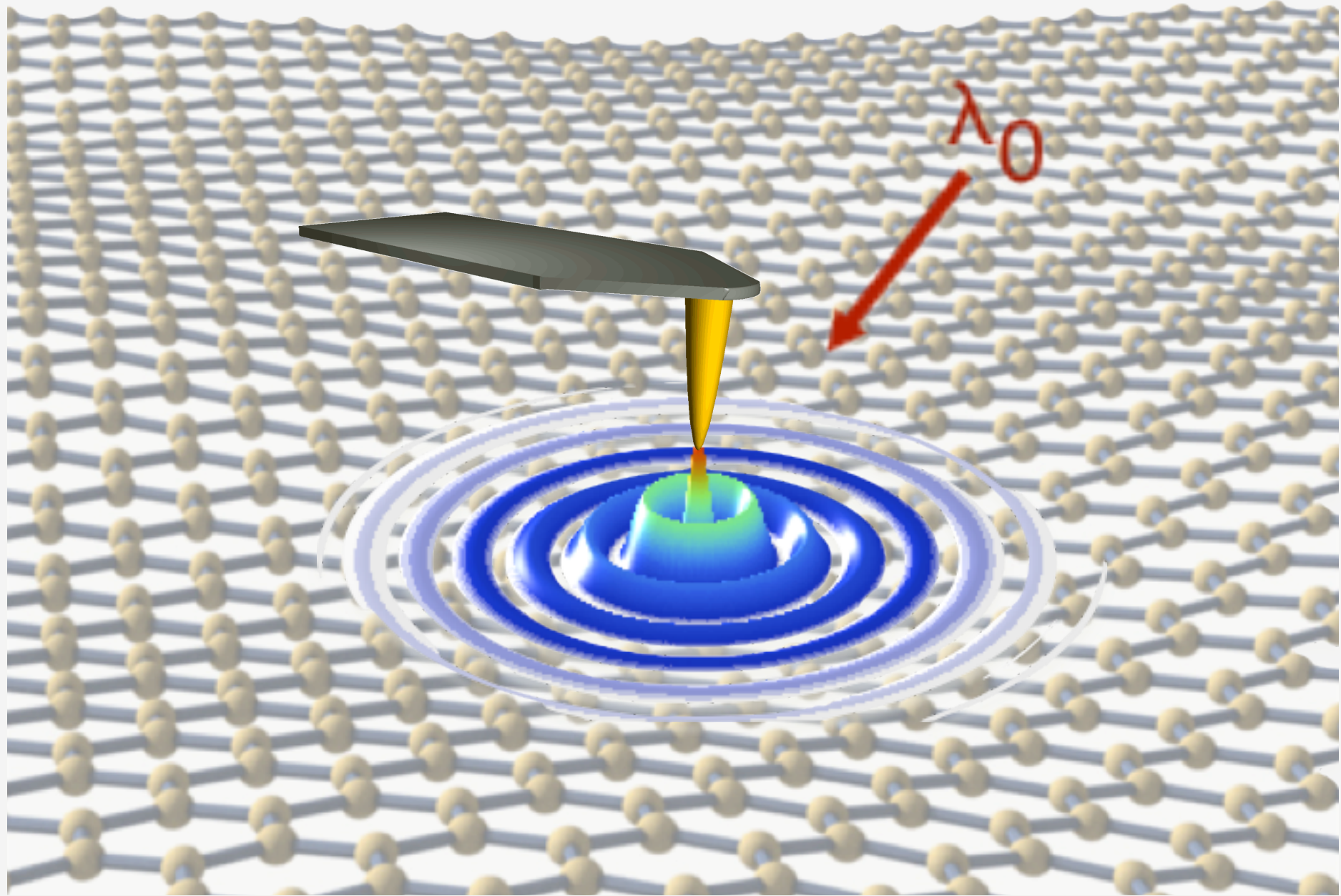
Experimental demonstration of spatial mapping and electrical tunability of graphene plasmons



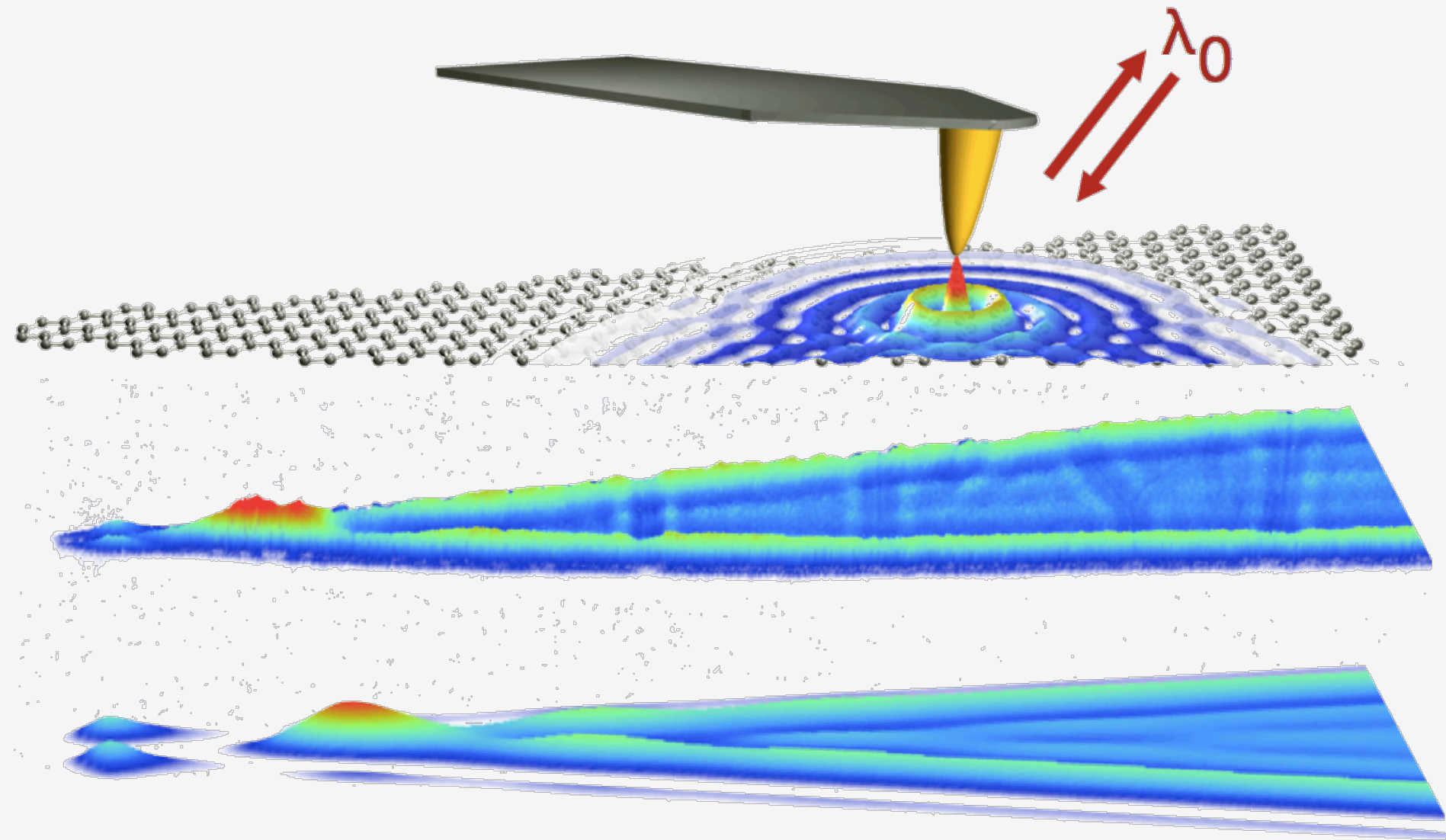
Chen *et al.*, Nature (2012)

Basov's group Fei *et al.*, Nature (2012)

Experimental mapping of graphene plasmons

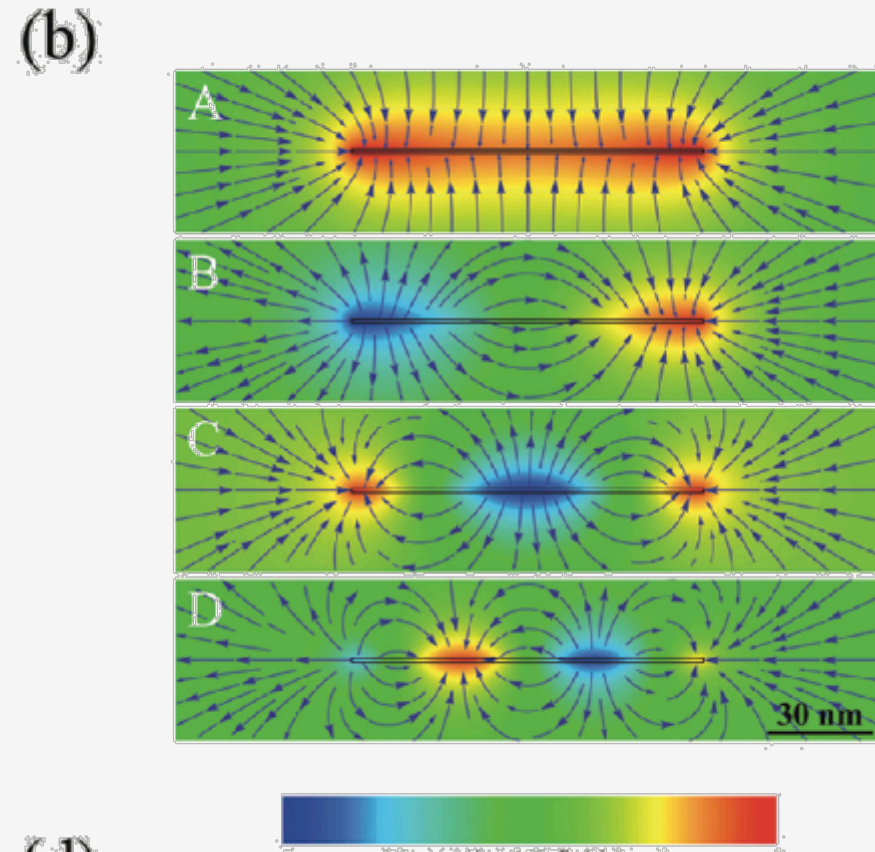
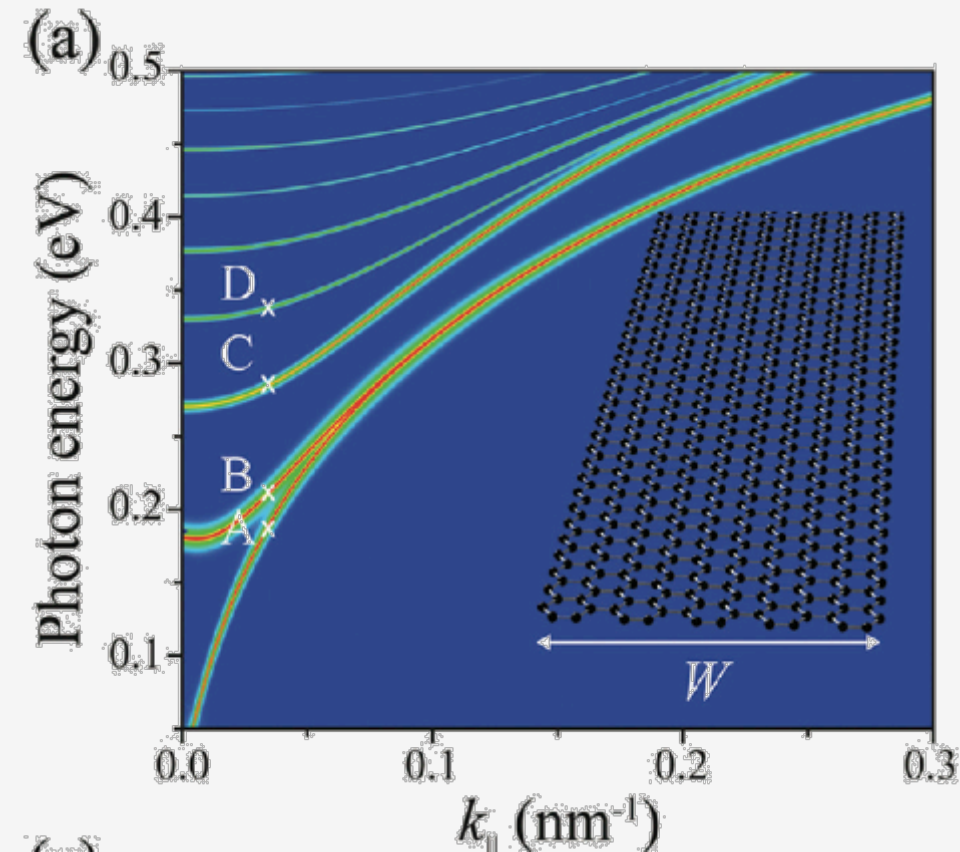


Experimental mapping of graphene plasmons



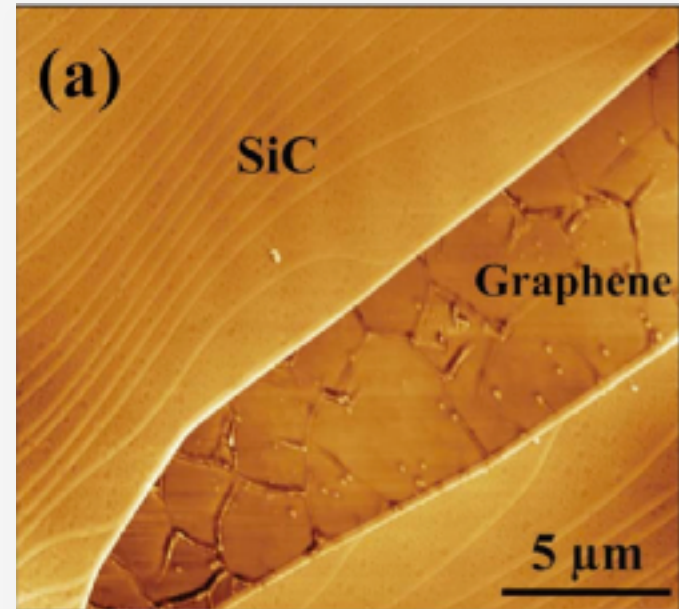
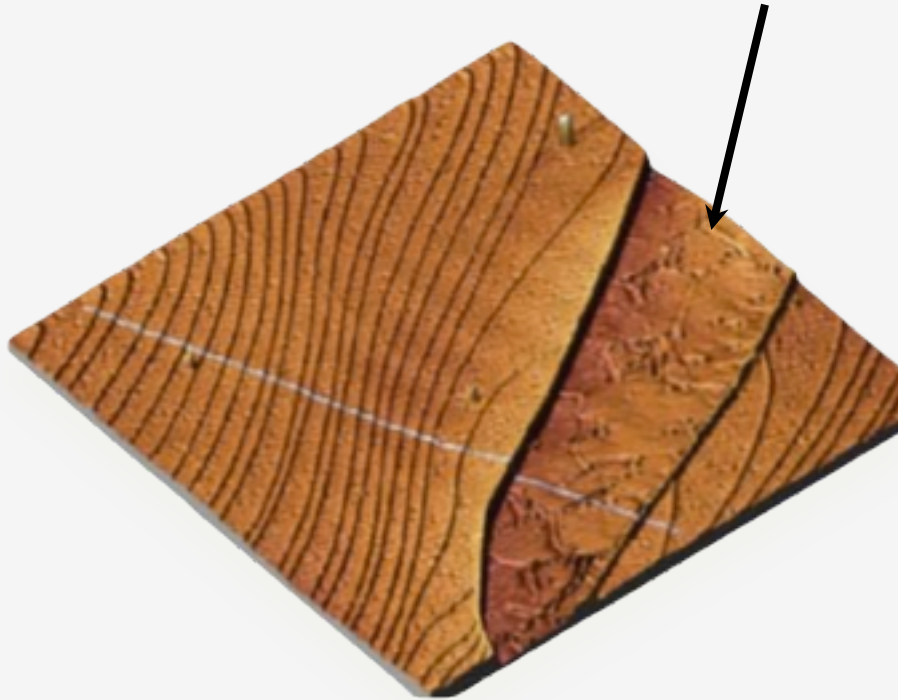
Chen *et al.*, Nature (2012)

Also, Basov's group Fei *et al.*, Nature (2012)



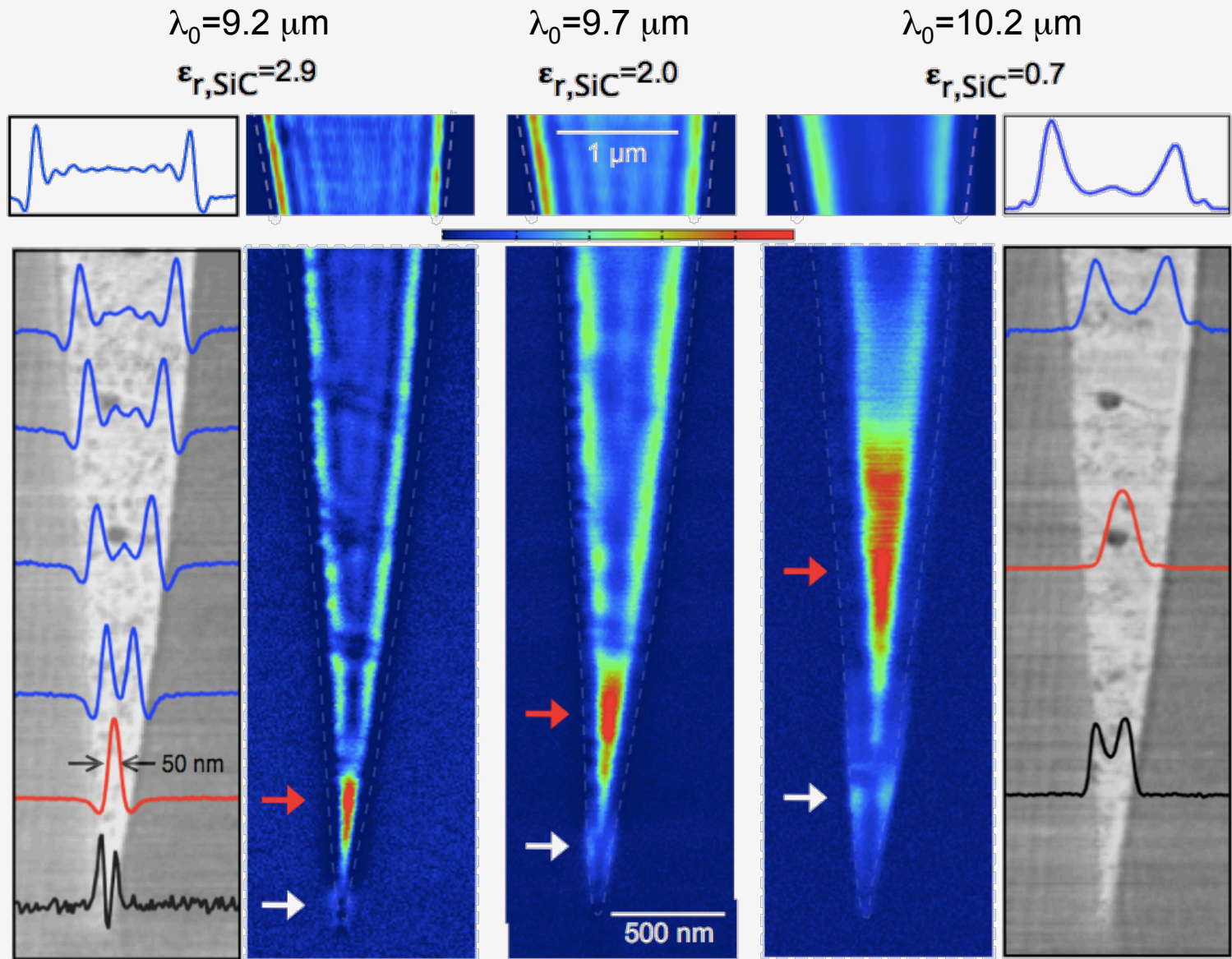
Christensen *et al.*, ACS Nano (2012)

Graphene on SiC



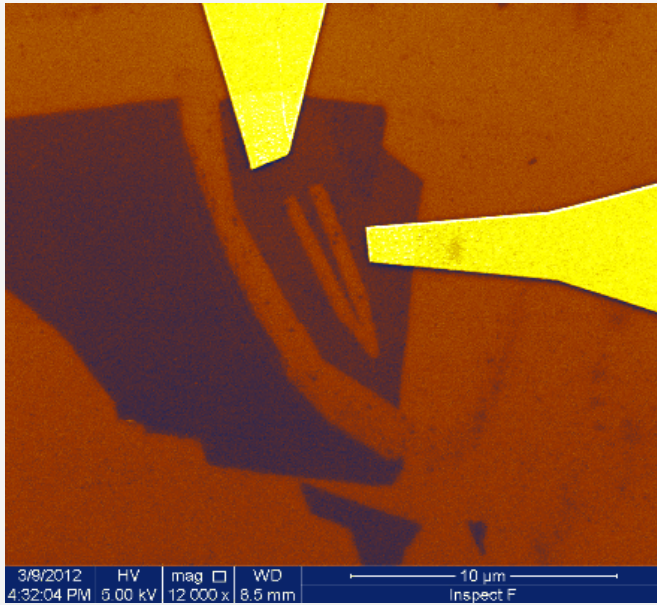
Camara *et al.*, PRB (2009)

Experimental mapping of graphene plasmons

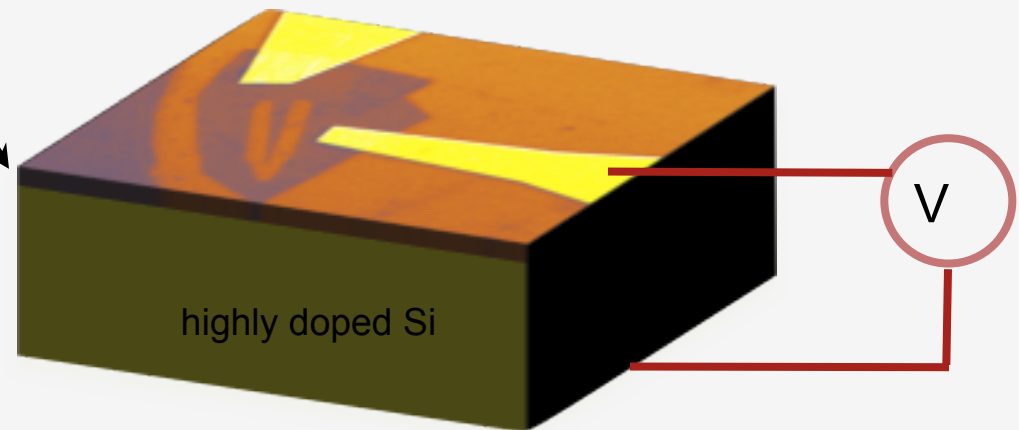


Chen *et al.*, Nature (2012)

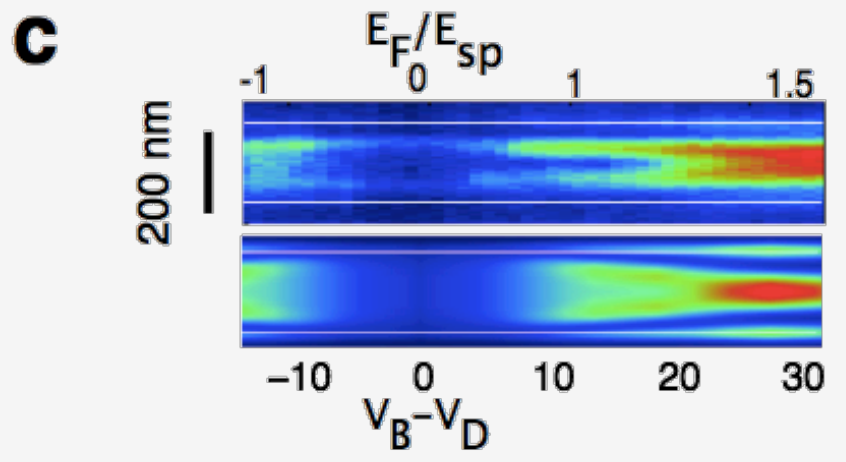
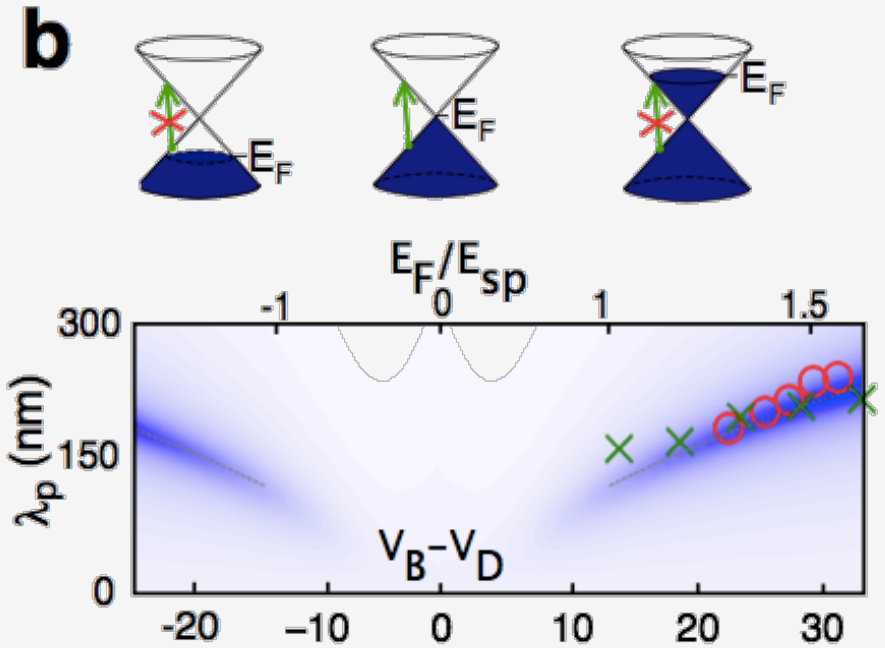
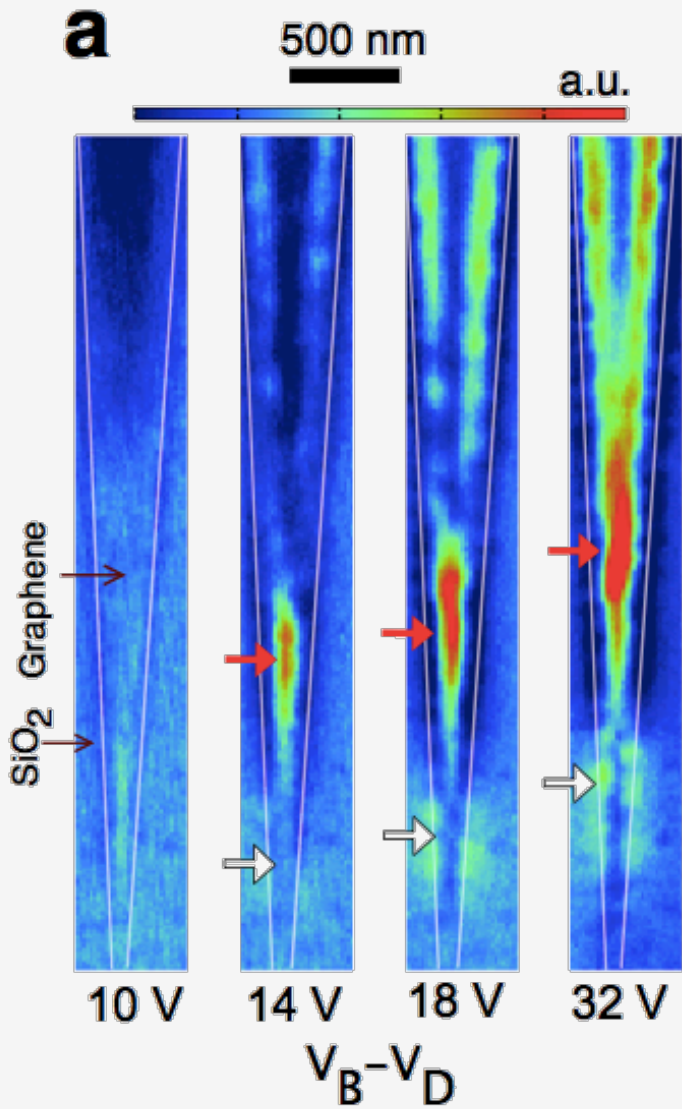
Controlling optical fields by electric fields



SiO₂

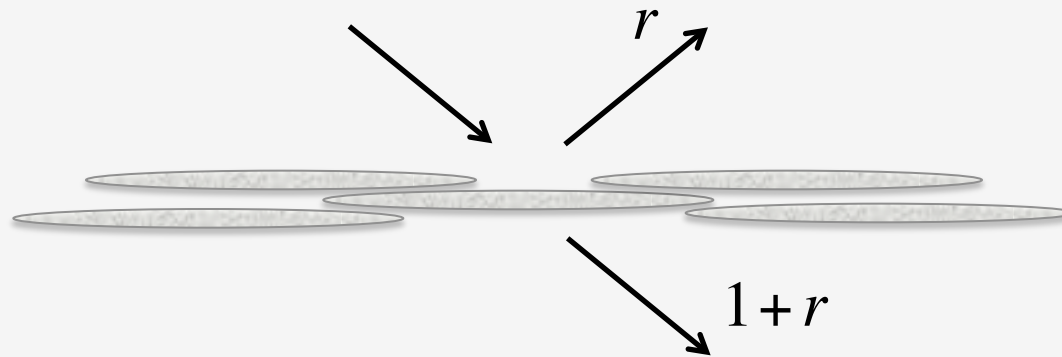
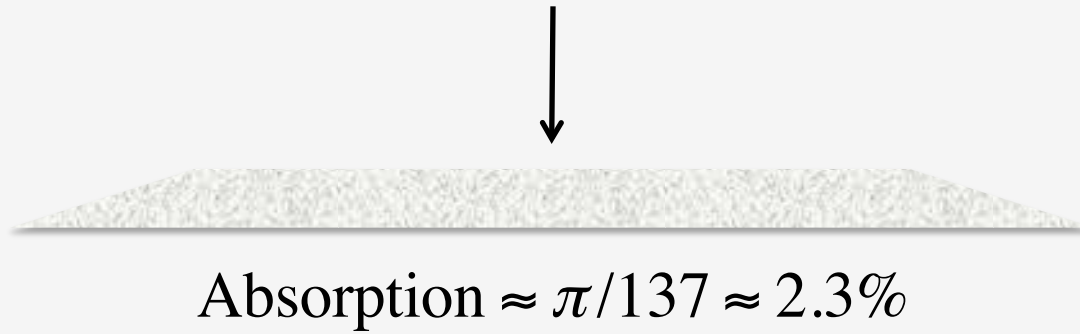


Experimental mapping of graphene plasmons



Chen *et al.*, Nature (2012)

Complete optical absorption in graphene



$$\text{Absorption} = 1 - |r|^2 - |1+r|^2 \quad \rightarrow \quad 50\% \text{ maximum}$$

Partially disordered silver films

- O. Hunderi and H. P. Myers, J. Phys. F: Metal Phys. 3, 683 (1973)

Diffraction in gratings, double-period metal gratings, and metamaterials

- M. C. Hutley and D. Maystre, Optics Communications 19, 431 (1976)
- D. Maystre and R. Petit, Optics Communications 17, 196 (1976)
- W.-C. Tan, J. R. Sambles, and T. W. Preist, Phys. Rev. B 61, 13177 (1999)
- E. Popov and L. Tsonev, Surface Science Letters 271, L378 (1992)
- N. I. Landy et al., Phys. Rev. Lett. 100, 207402 (2008)
- N. Liu et al., Nano Lett. 10, 2342 (2010)

Doped silicon lamellar grating

- F. Marquier, M. Laroche, R. Carminati, J.-J. Greffet, Journal of Heat Transfer 129, 11 (2007)
- J.-J. Greffet, R. Carminati, K. Joulain, J.-P. Mulet, S. Mainguy, Y. Chen, Nature 416, 61 (2002)

Semiconductor and metal-semiconductor-metal nanostructures

- S. Collin, F. Pardo, R. Teissier, and J.-L. Pelouard, Appl. Phys. Lett. 85, 194 (2004)
- T.V. Teperik, F.J. García de Abajo, V.V. Popov, and M.S. Shur, Appl. Phys. Lett. 90 251910 (2007)

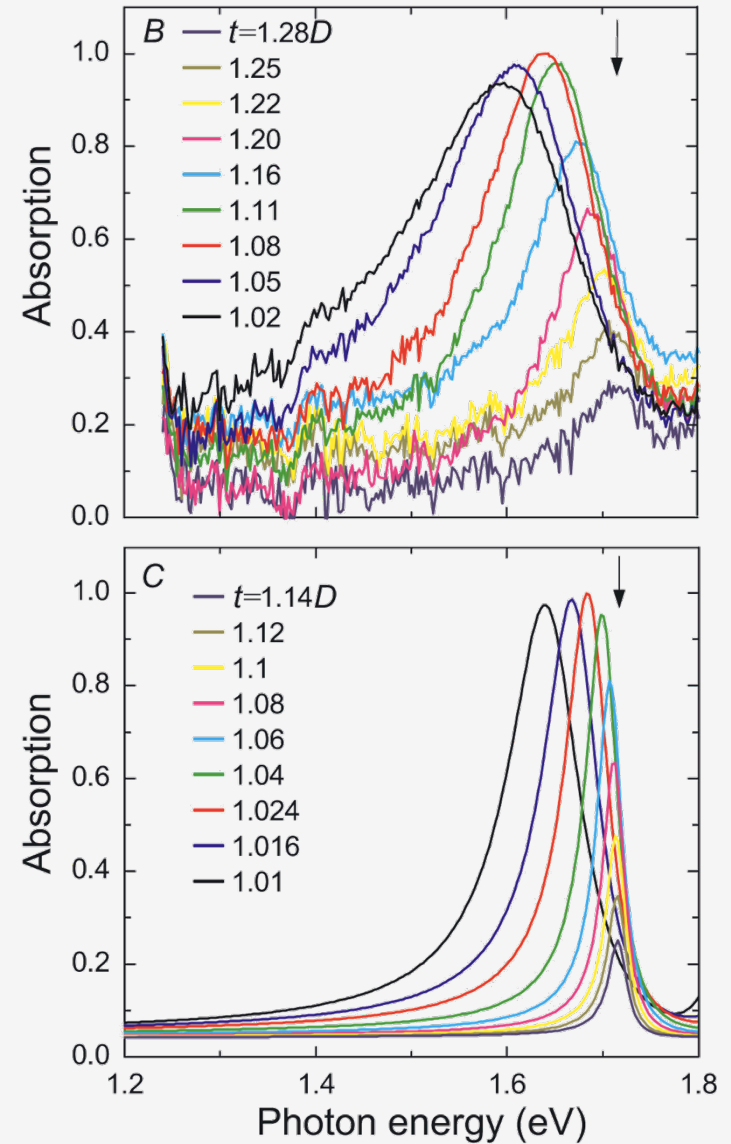
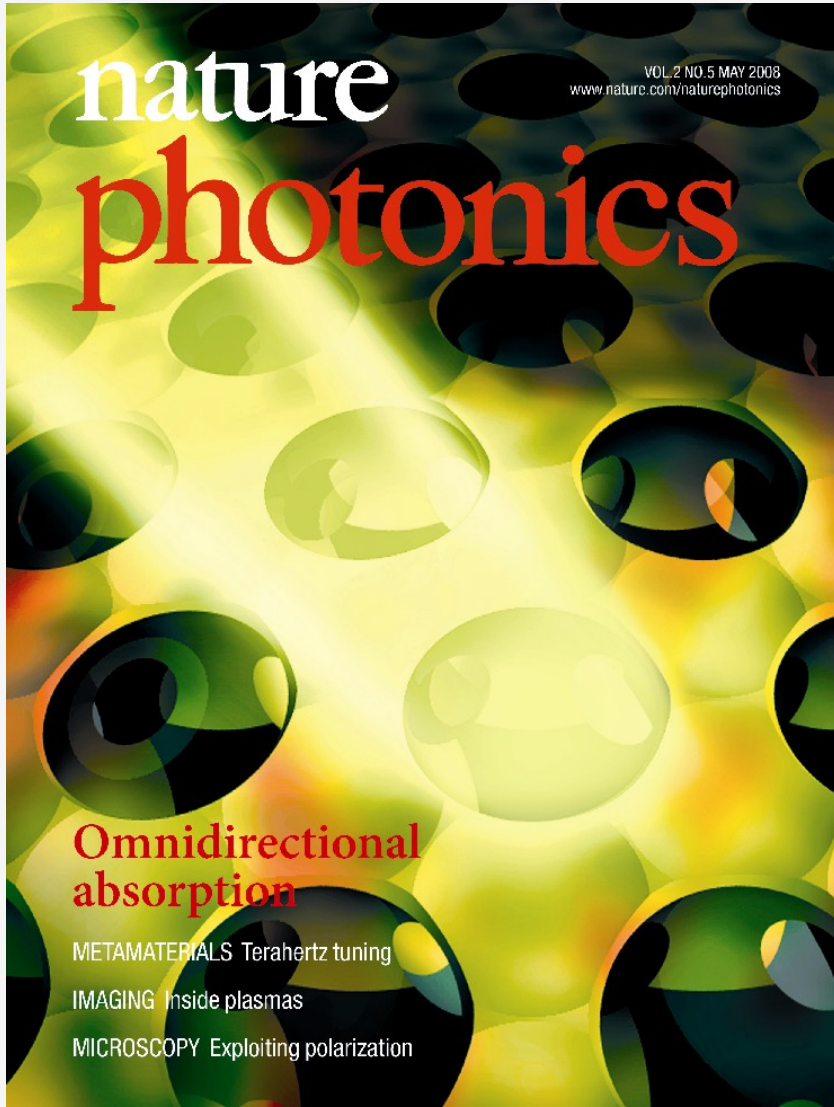
Multiplayer of metallic nanoparticles and nanopores in metal

- T. V. Teperik, V. V. Popov, and F. J. García de Abajo, Phys. Rev. B 71, 085408 (2005)
- T. Teperik, V. Popov, and F. Garcia de Abajo, J. Opt. A: Pure Appl. Opt. 0, 0 (2007)
- S.Kachan, O. Stenzel, and A. Ponyavina, Appl. Phys. B 84, 281 (2006)

Overdense plasma slab (in the microwave frequency range)

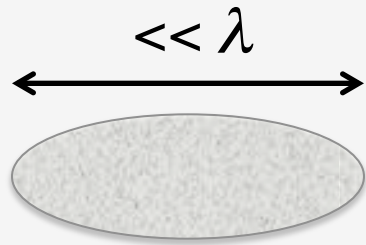
- Y. P. Bliokh, J. Felsteiner, and Y. Z. Slutsker, Phys. Rev. Lett. 95, 165003 (2005)

Nanovoids as perfect absorbers



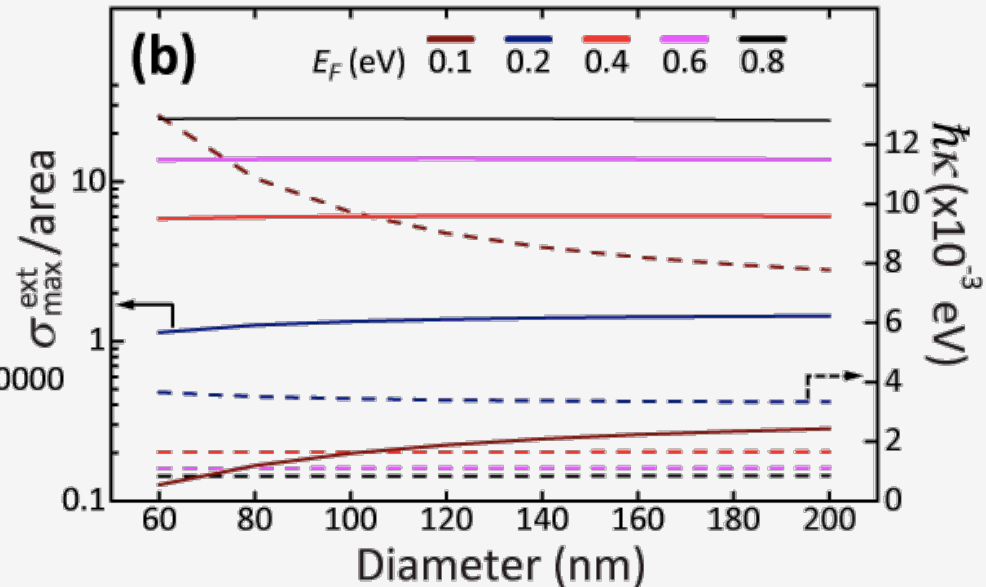
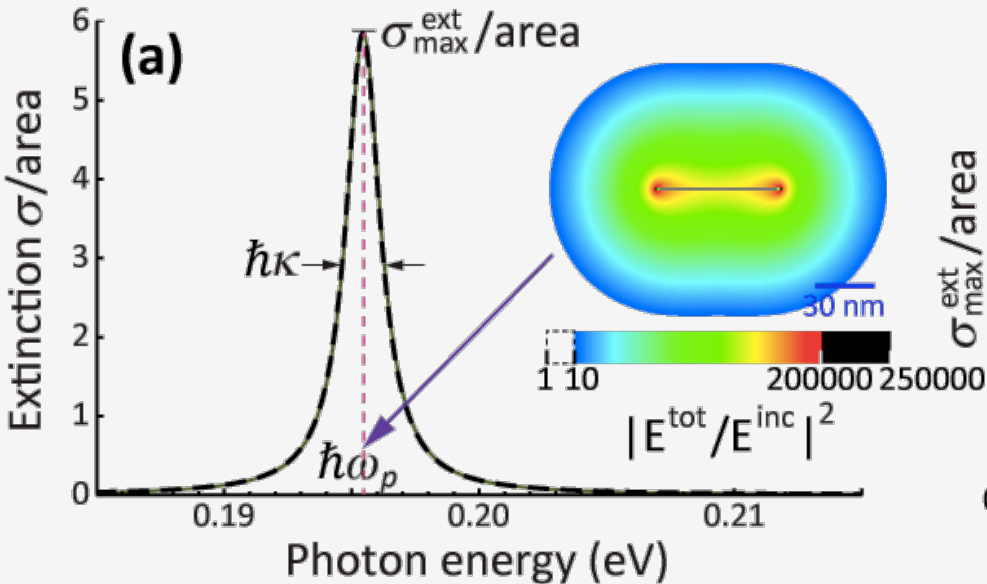
T.V.Teperik *et al.*, Nature Phot. (2008)

Plasmons in graphene disks



$$\alpha(\omega) = \frac{3c^3}{4\omega^3} \frac{\kappa_r}{\omega_p - \omega - i\kappa/2}$$

$$\sigma^{\text{ext}}(\omega) = \frac{4\pi\omega}{c} \text{Im}\{\alpha\} \approx \frac{3\lambda^2}{2\pi} \frac{\kappa_r}{\kappa}, \quad \kappa_r \ll \kappa$$



Thongrattanasiri *et al.*, Phys. Rev. Lett. (2012)

Maximum absorption by a small particle

$$\alpha(\omega) = \frac{3c^3}{4\omega^3} \frac{\kappa_r}{\omega_p - \omega - i\kappa/2}$$

$$\alpha(\omega_p) = \frac{3ic^3}{2\omega_p^3} \frac{\kappa_r}{\kappa}$$

$$\sigma^{\text{ext}}(\omega) = \frac{4\pi\omega}{c} \text{Im}\{\alpha\}$$

$$\sigma^{\text{ext}}(\omega_p) = \frac{3\lambda^2}{2\pi} \frac{\kappa_r}{\kappa}$$

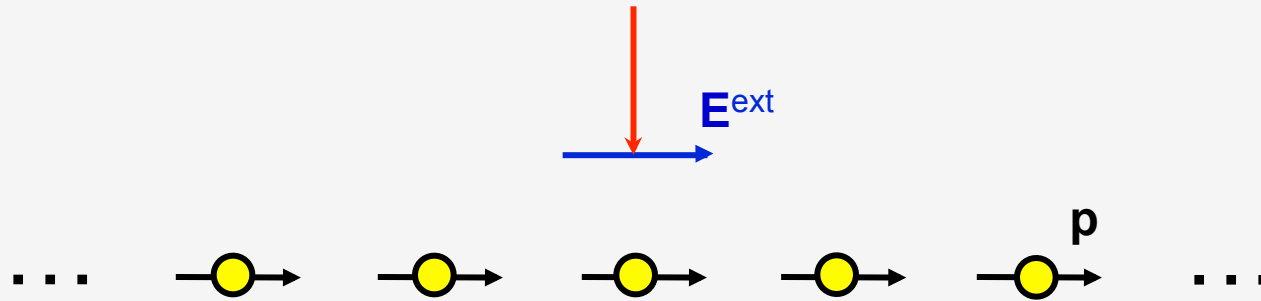
$$\sigma^{\text{abs}}(\omega) = \frac{4\pi\omega}{c} \left(\text{Im}\{\alpha\} - \frac{2\omega^3}{3c^3} |\alpha|^2 \right)$$

$$\sigma^{\text{abs}}(\omega_p) = \frac{3\lambda^2}{2\pi} \frac{\kappa_r}{\kappa} \left(1 - \frac{\kappa_r}{\kappa} \right)$$

$$\sigma_{\text{max}}^{\text{ext}}(\omega_p) = \frac{3\lambda^2}{2\pi} \quad \text{for } \kappa = \kappa_r$$

$$\sigma_{\text{max}}^{\text{abs}}(\omega_p) = \frac{3\lambda^2}{8\pi} \quad \text{for } \kappa = 2\kappa_r$$

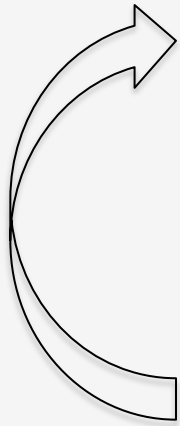
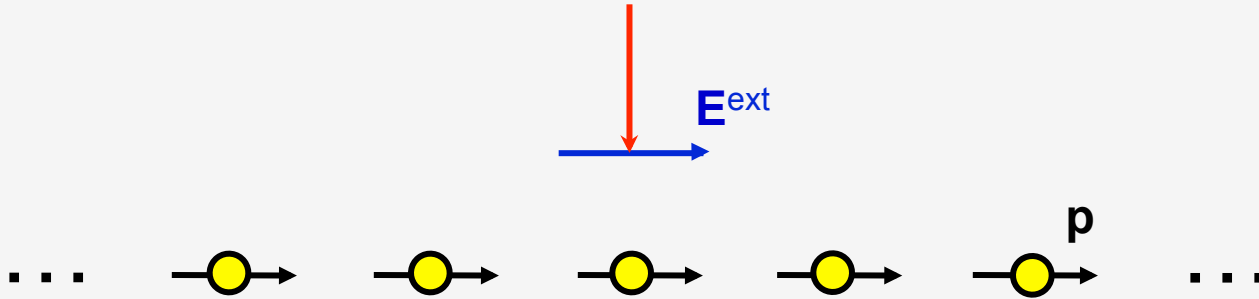
$$\boxed{\kappa_r = \kappa_a}$$



$$\mathbf{p} = \alpha \left(\mathbf{E}^{ext} + \sum_{\text{lattice}} G_{\text{dip-dip}} \mathbf{p} \right), \quad G_{\text{dip-dip}} \approx \frac{3\hat{\mathbf{r}}\hat{\mathbf{r}} - 1}{r^3}$$

$$\mathbf{p} = \frac{1}{\alpha^{-1} - G} \mathbf{E}^{ext} \quad G = \sum_{\text{lattice}} G_{\text{dip-dip}}$$

Description of graphene-disk array



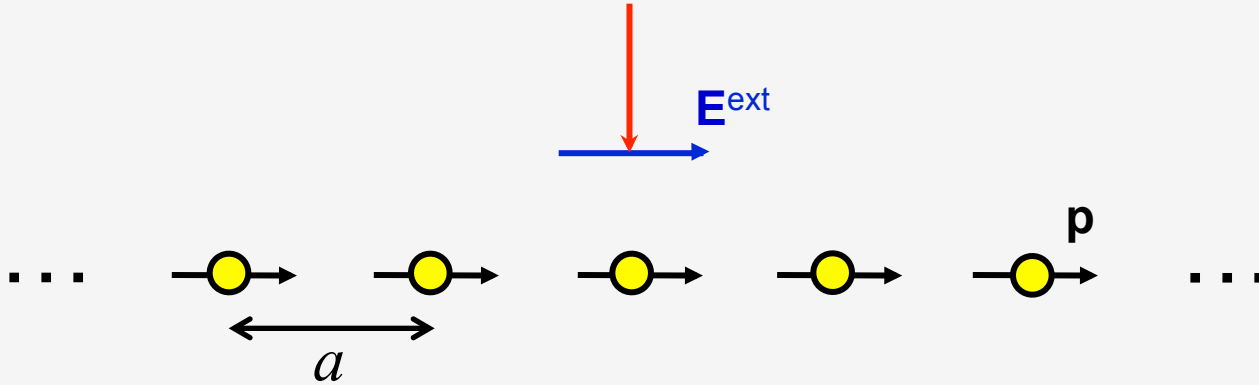
$$r = \frac{\pm iS}{\alpha^{-1} - G}$$

$$S = \frac{2\pi\omega}{Ac} (\cos\theta)^{\mp 1}$$

$$\mathbf{p} = \frac{1}{\alpha^{-1} - G} \mathbf{E}^{ext}$$

$$G = \sum_{\text{lattice}} G_{\text{dip-dip}}$$

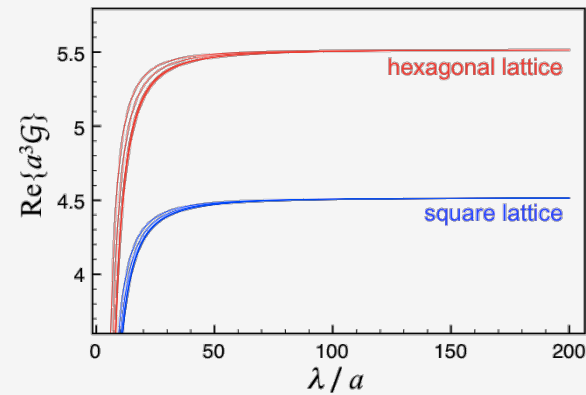
Description of graphene-disk array



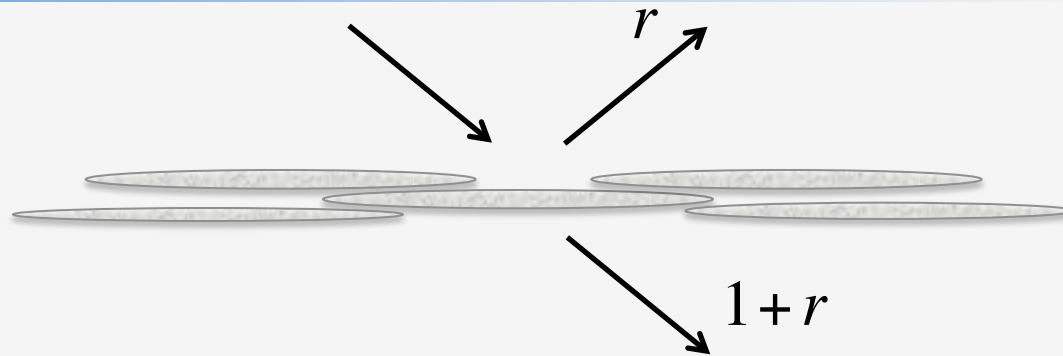
$$r = \frac{\pm iS}{\alpha^{-1} - G}$$

$$S = \frac{2\pi\omega}{Ac} (\cos\theta)^{\mp 1}$$

$$G \approx \frac{g}{a^3} + i \left(S - \frac{2\omega^3}{3c^3} \right)$$



Maximum absorption in graphene



$$\text{Absorption} = 1 - |r|^2 - |1+r|^2 \quad \rightarrow \quad 50\% \text{ maximum for } r = -1/2$$

$$r = \frac{\pm iS}{\alpha^{-1} - G}$$

$$S = \frac{2\pi\omega}{Ac} (\cos\theta)^{\mp 1}$$

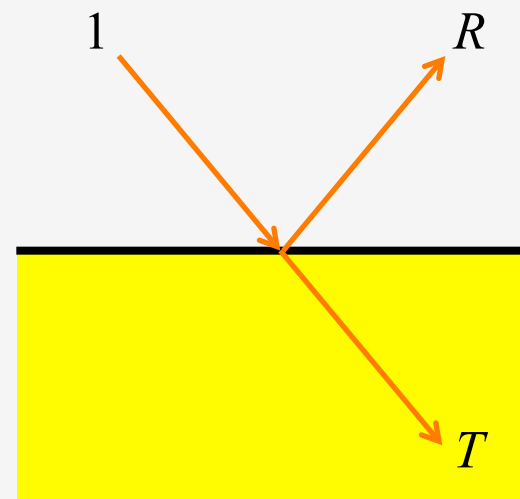
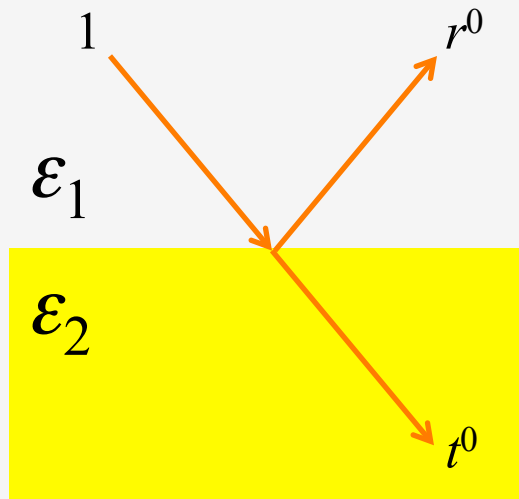
$$G \approx \frac{g}{a^3} + i \left(S - \frac{2\omega^3}{3c^3} \right)$$

$$\alpha(\omega) = \frac{3c^3}{4\omega^3} \frac{\kappa_r}{\omega_p - \omega - i\kappa/2}$$

$$\omega \approx \omega_p - 3g\kappa_r / 4(\omega a/c)^3$$

$$\sigma_{\max}^{\text{ext}} = 2A \times \begin{cases} \cos\theta, & s \text{ polarization,} \\ \cos^{-1}\theta, & p \text{ polarization.} \end{cases}$$

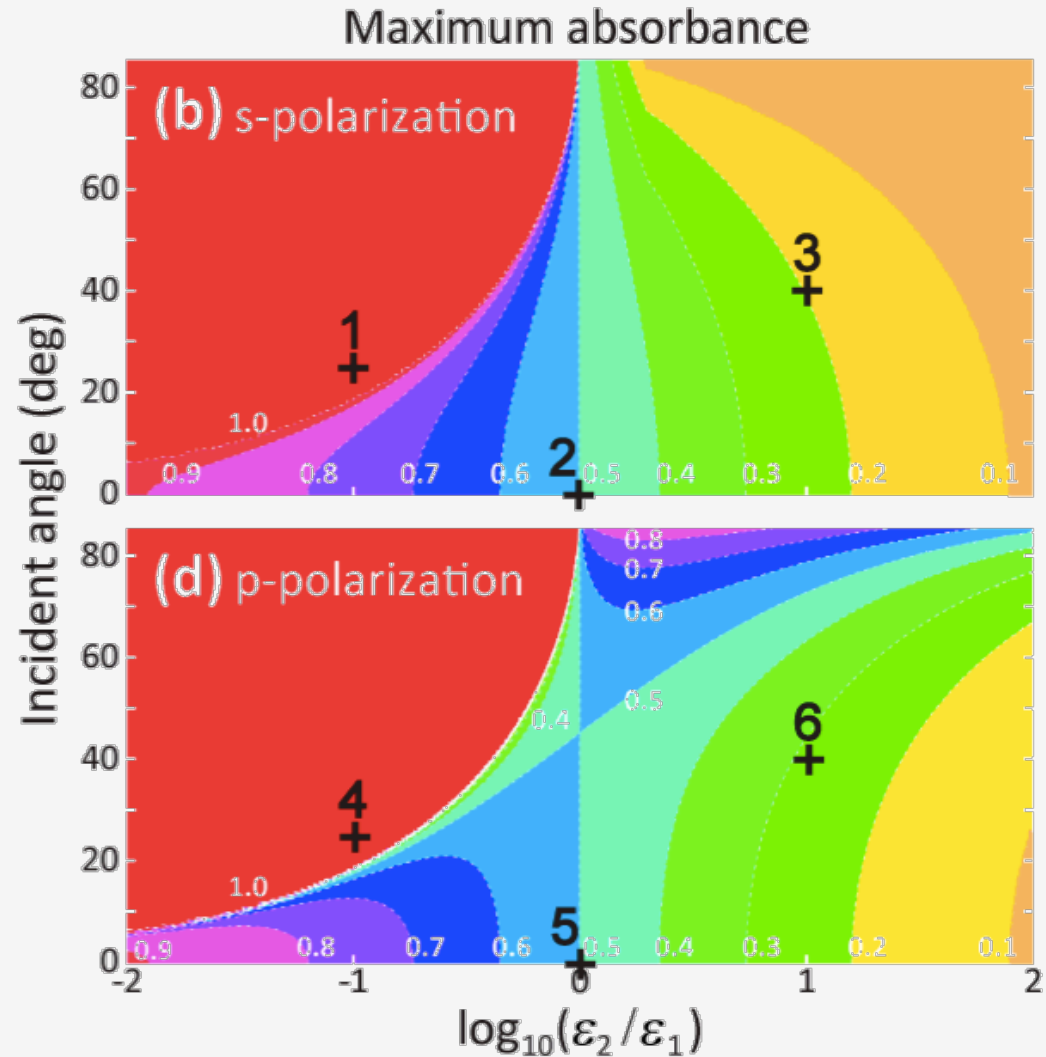
Thongrattanasiri *et al.*, Phys. Rev. Lett. (2012)



$$\eta = -\frac{r^0 \pm |r^0|^2 \pm \text{Re}\{f\}|t^0|^2}{|1 \pm r^0|^2 + \text{Re}\{f\}|t^0|^2}$$

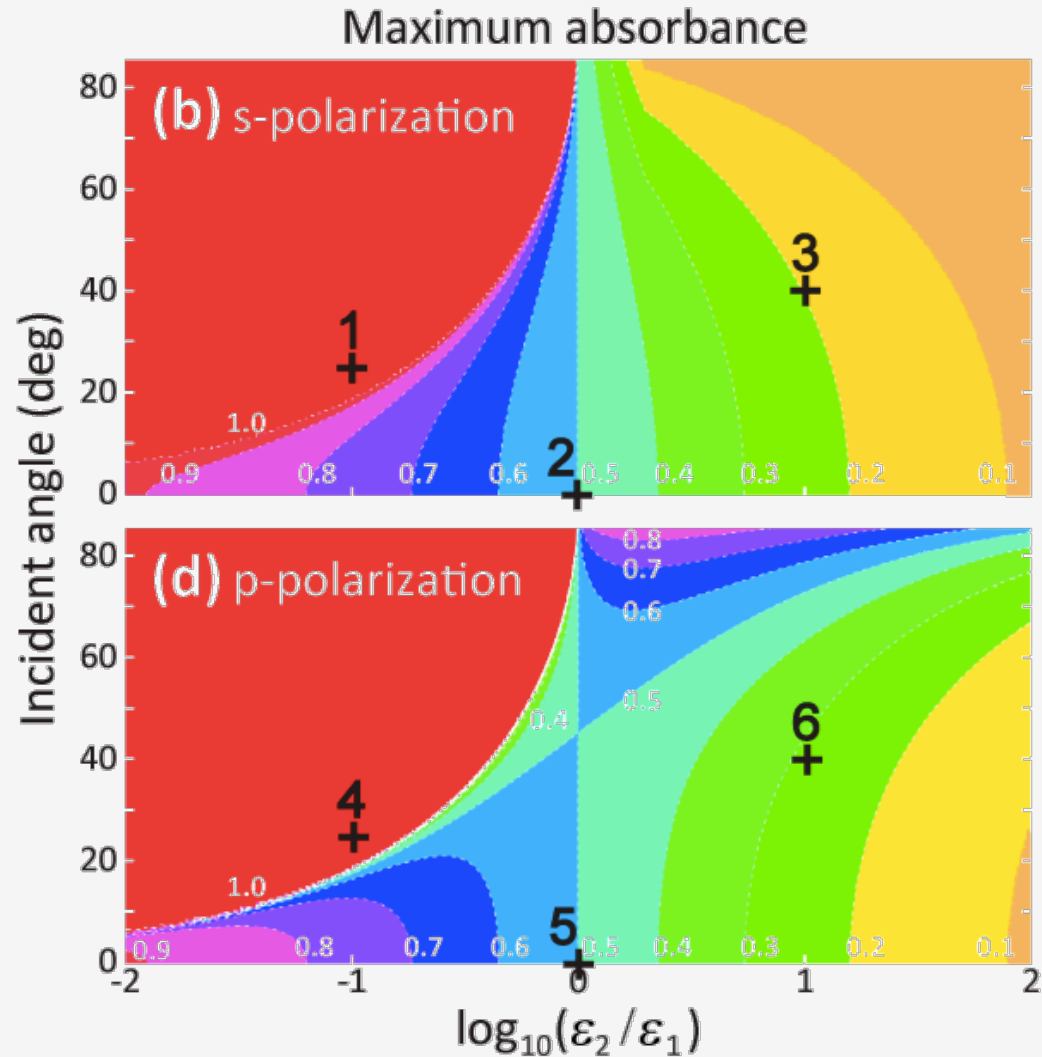
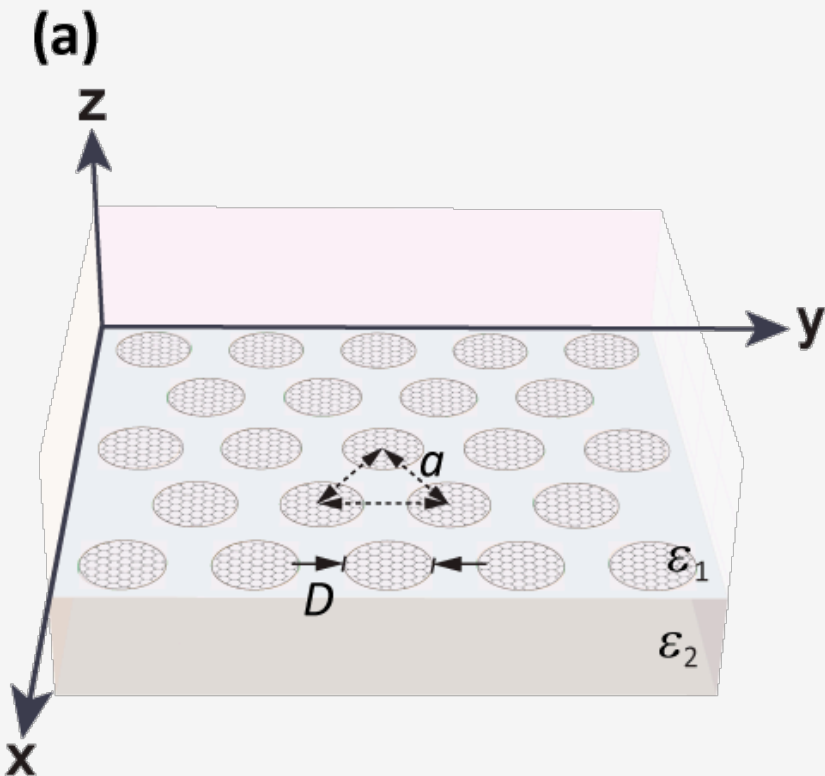
$$\mathcal{A} = 1 - |r^0 + (1 \pm r^0)\eta|^2 - \text{Re}\{f\}|t^0|^2|1 \pm \eta|^2$$

$$f = (\epsilon_2/\epsilon_1 - \sin^2 \theta)^{1/2} / \cos \theta$$



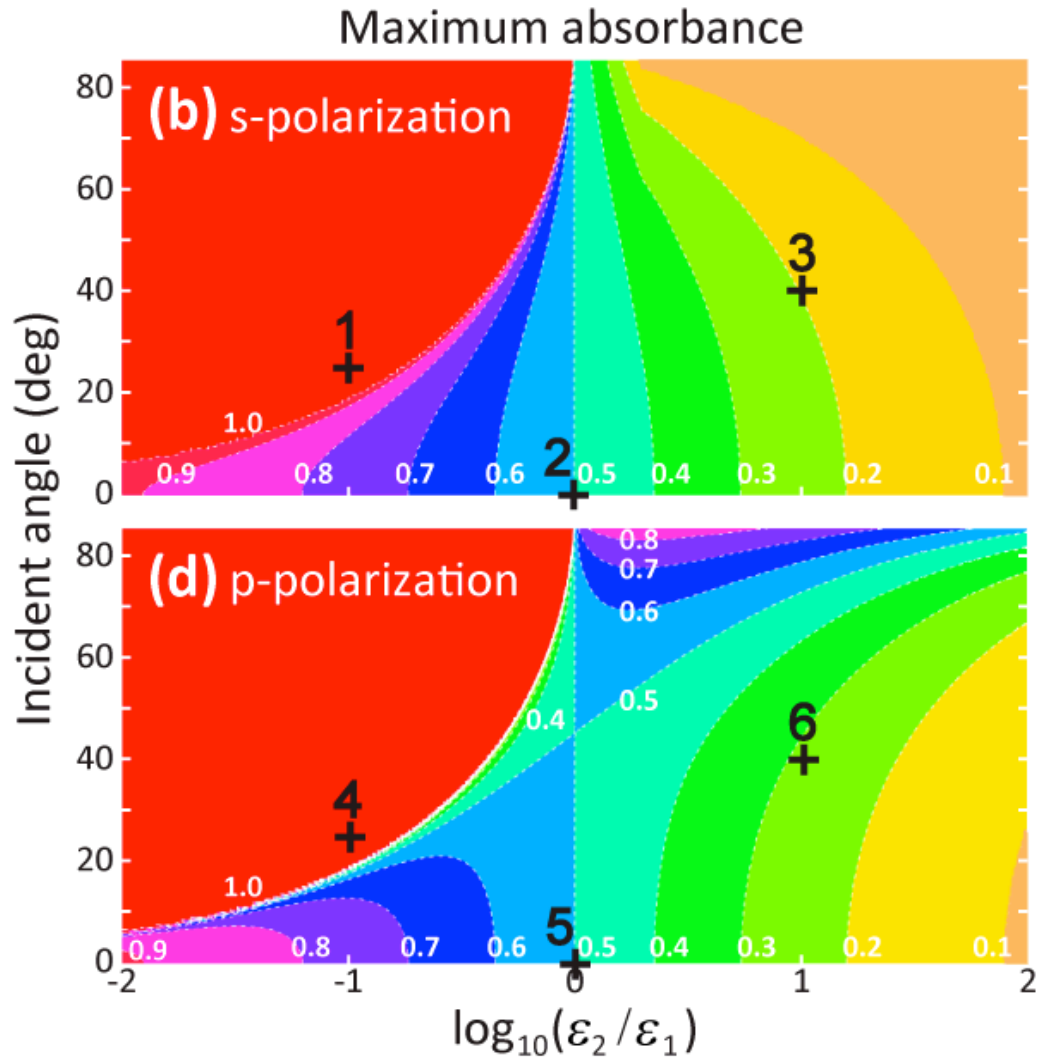
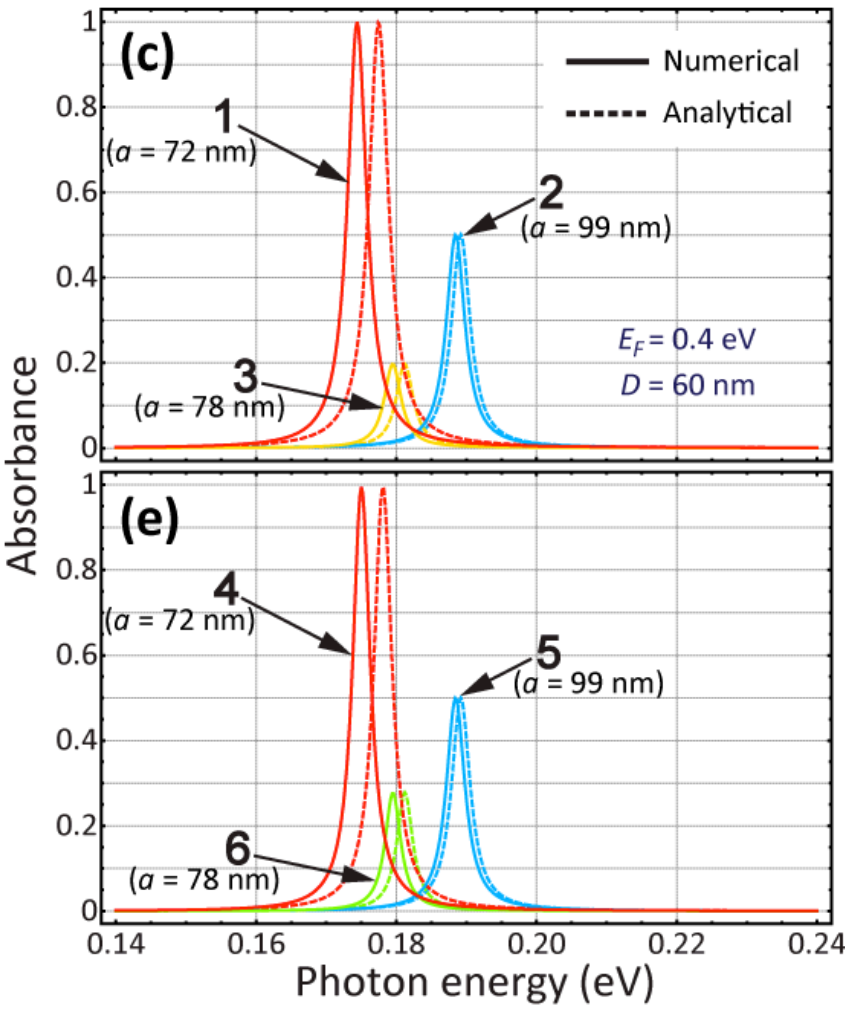
Thongrattanasiri *et al.*, Phys. Rev. Lett. (2012)

Absorption in asymmetric environments

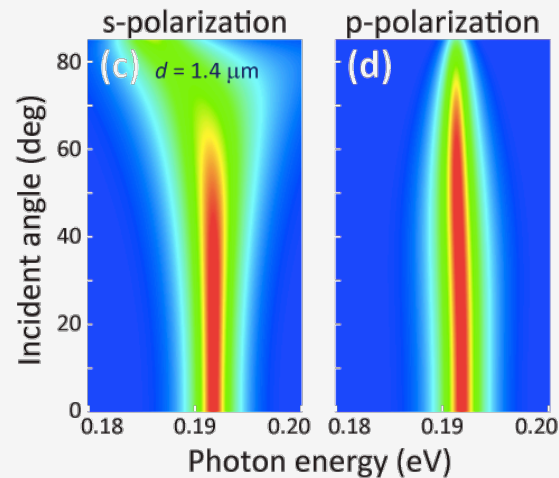
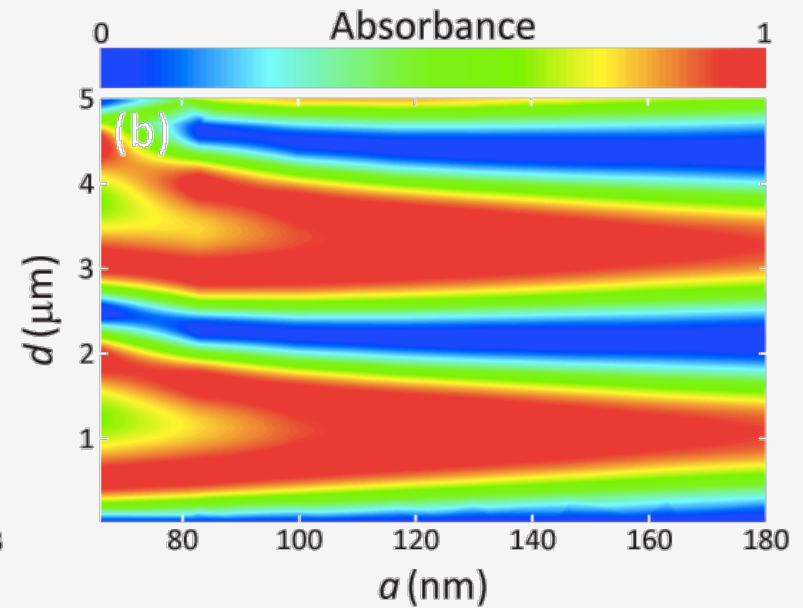
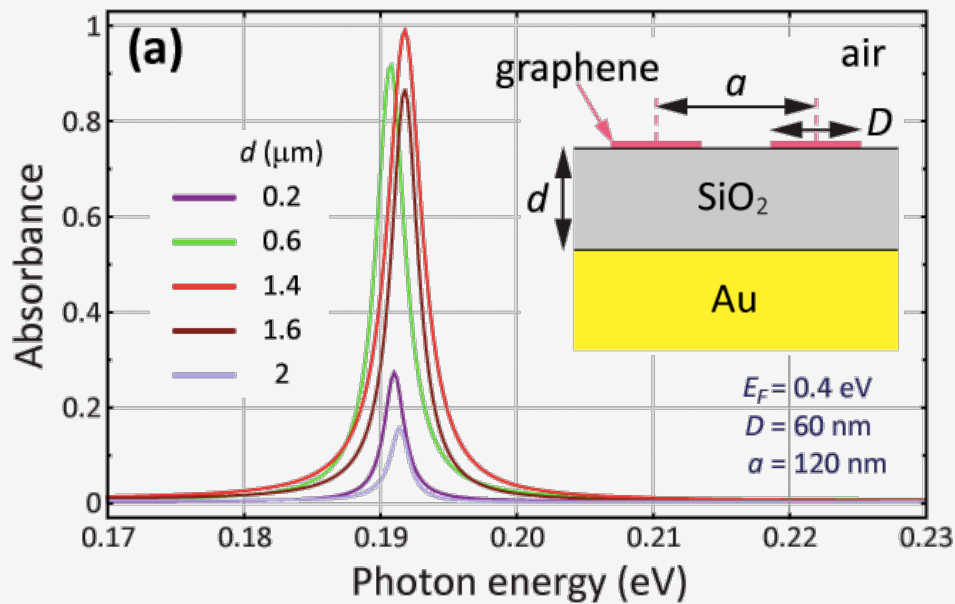


Thongrattanasiri *et al.*, Phys. Rev. Lett. (2012)

Absorption in asymmetric environments



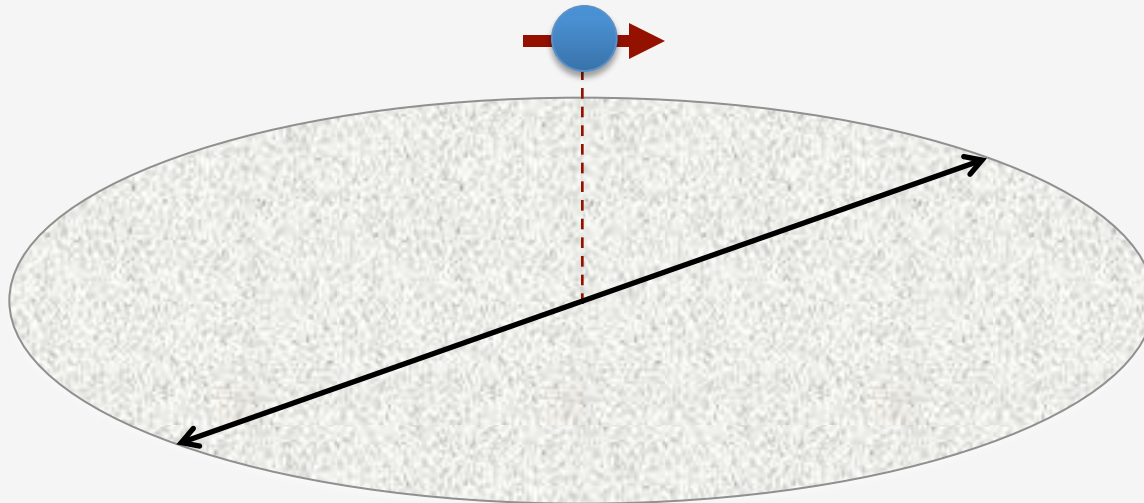
Thongrattanasiri *et al.*, Phys. Rev. Lett. (2012)



Thongrattanasiri *et al.*, Phys. Rev. Lett. (2012)

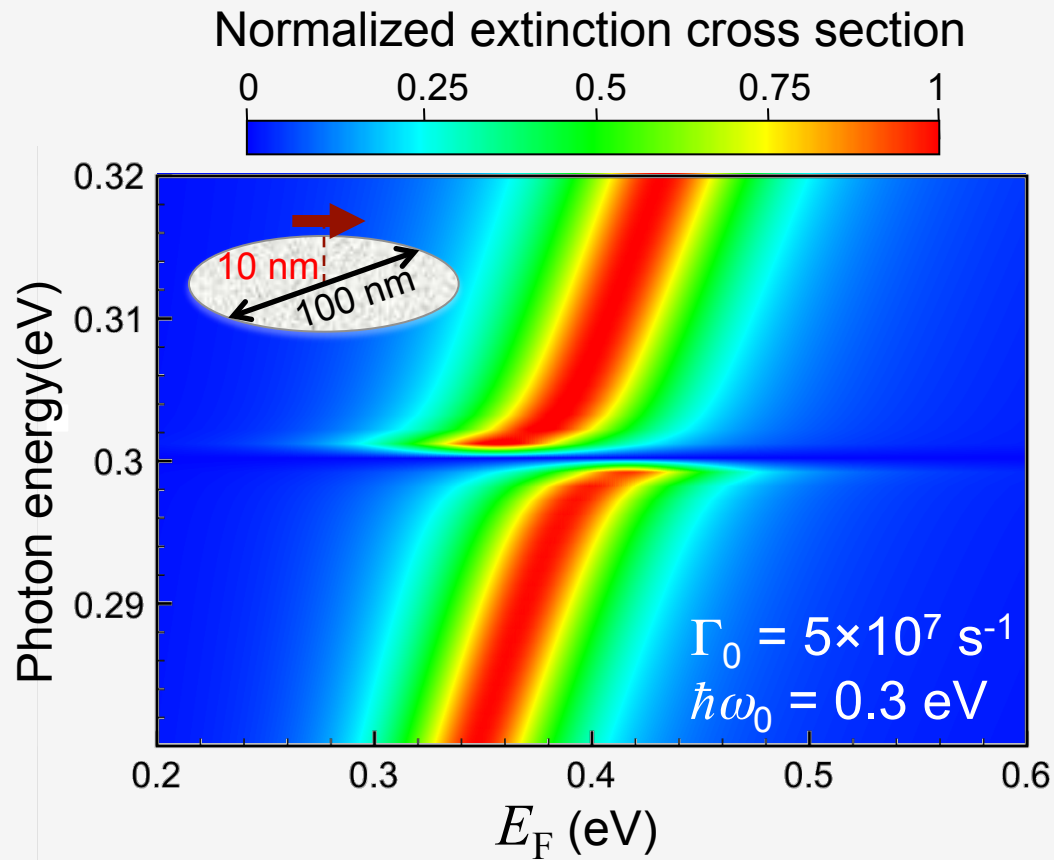
Strong light-matter interaction:
quantum plasmonics with graphene

Vacuum Rabi splitting



Koppens, Chang, and García de Abajo, *Nano Lett.* (2011)

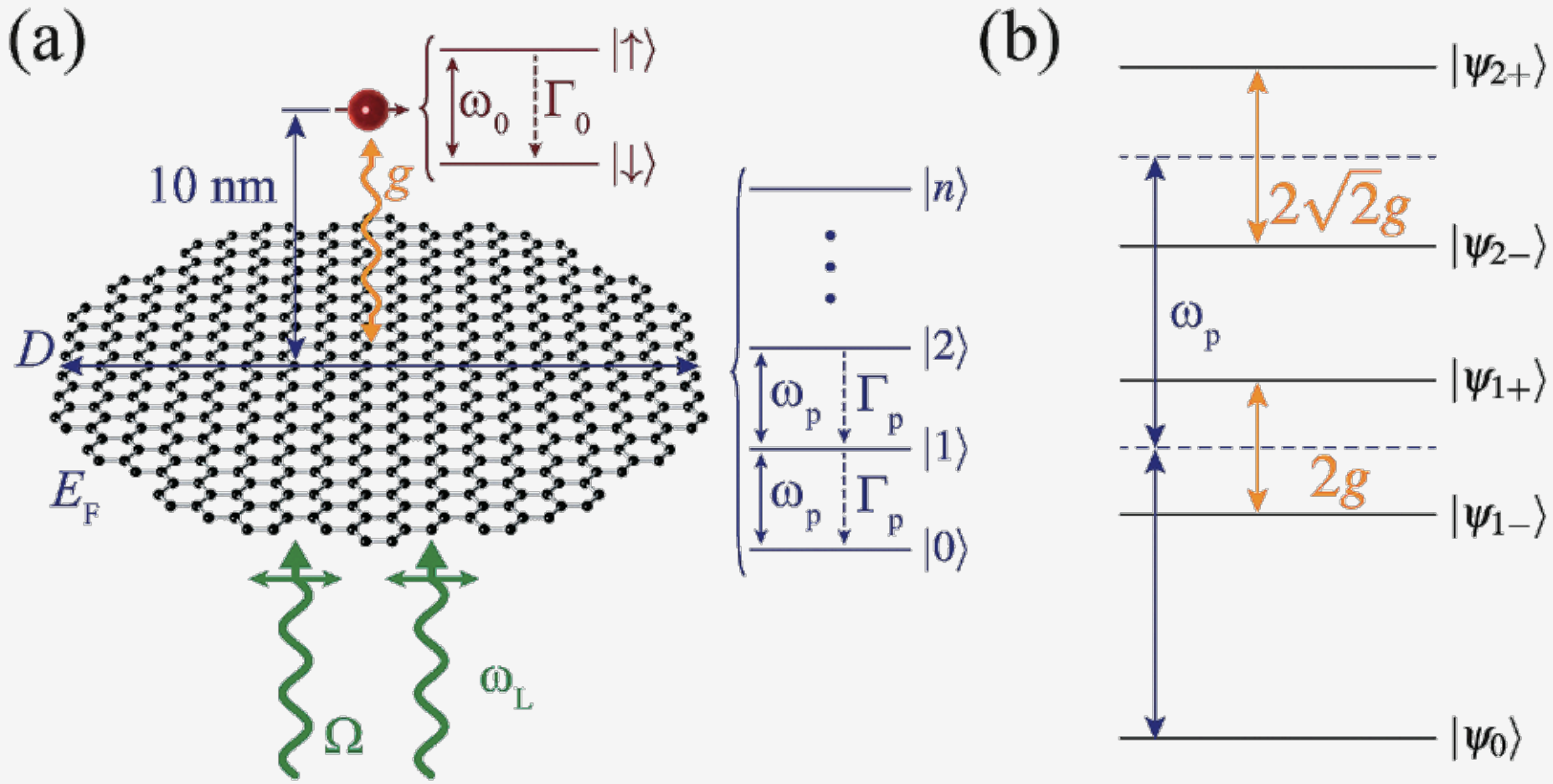
Vacuum Rabi splitting



Koppens, Chang, and García de Abajo, Nano Lett. (2011)

Strong light-matter interaction

Jaynes-Cummings ladder



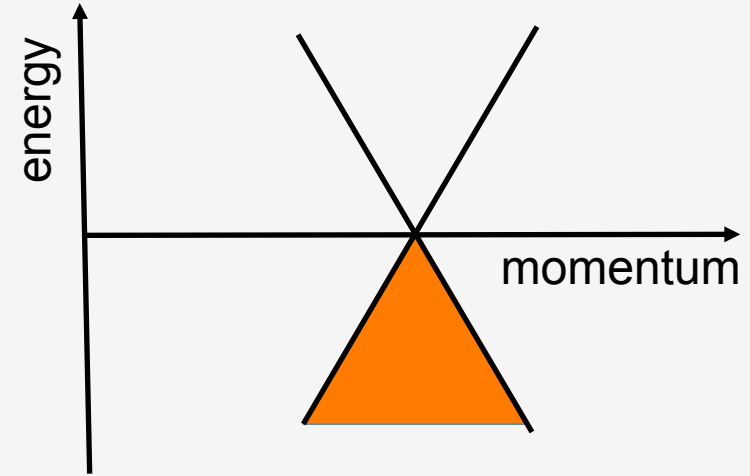
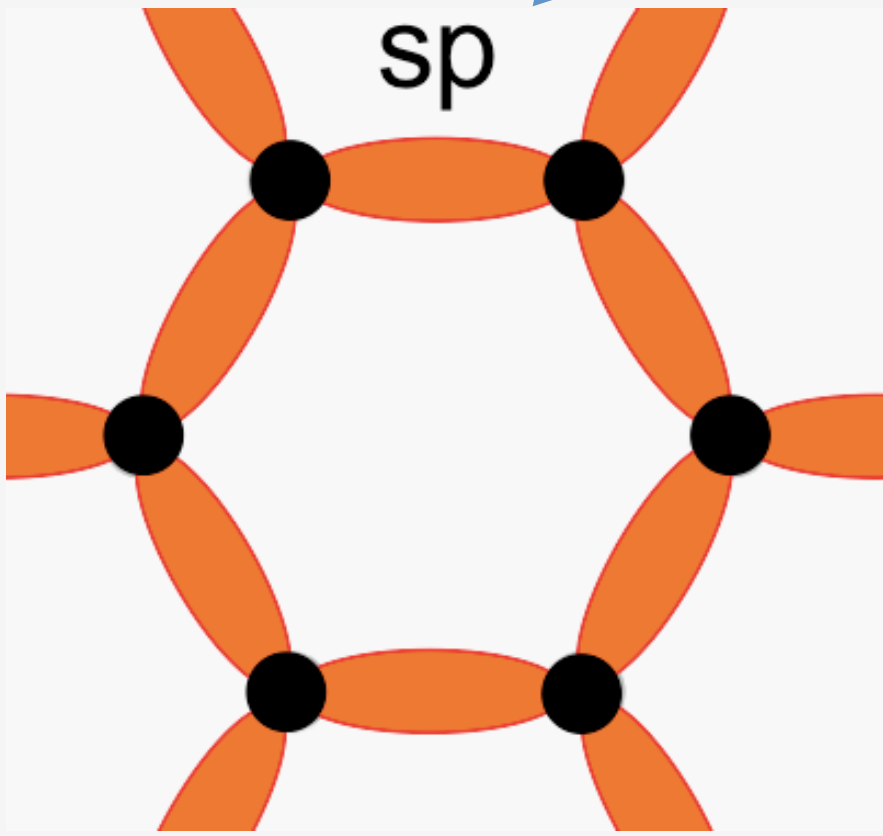
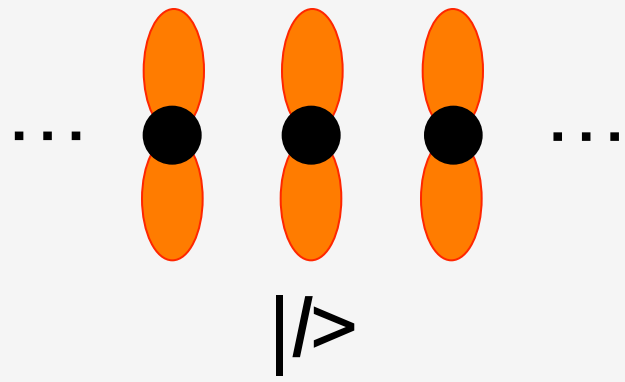
Manjavacas *et al.*, ACS Nano (2012)

Quantum effects in graphene plasmons

Quantum effects in graphene plasmons



$2p_z$



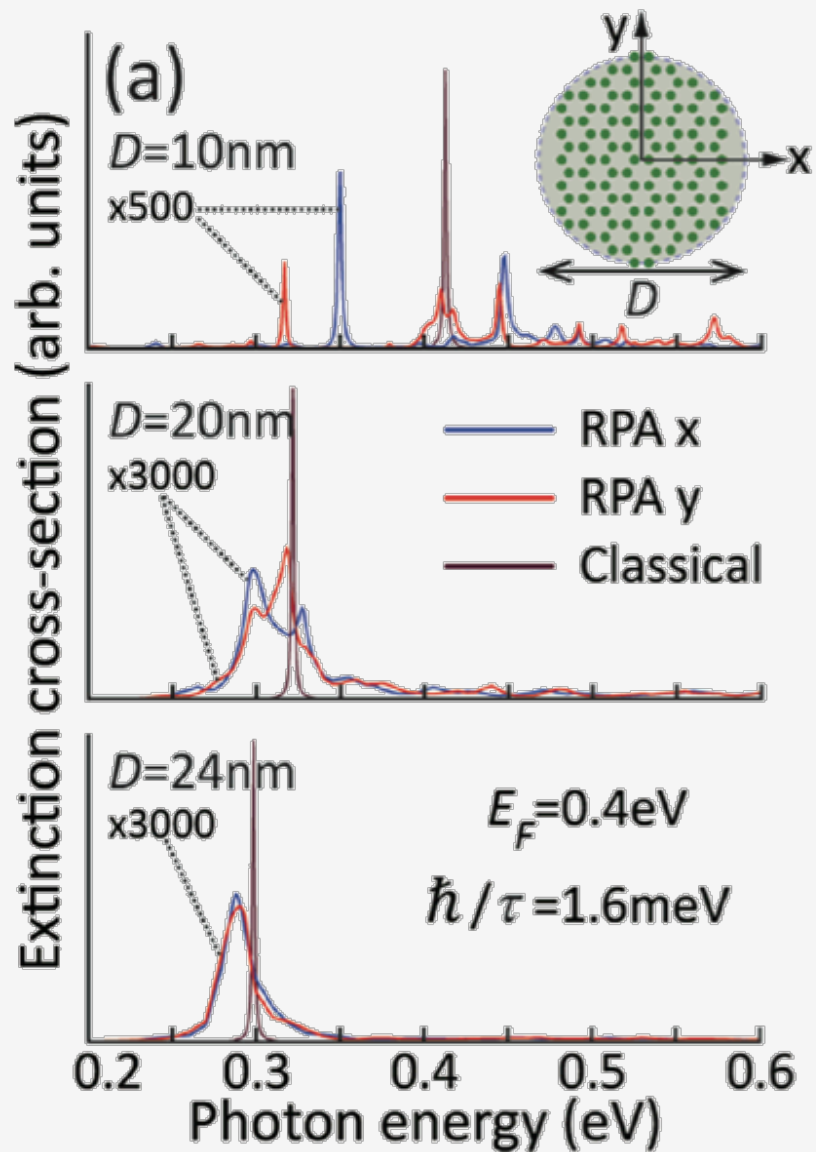
Electron state j
of energy ε_j

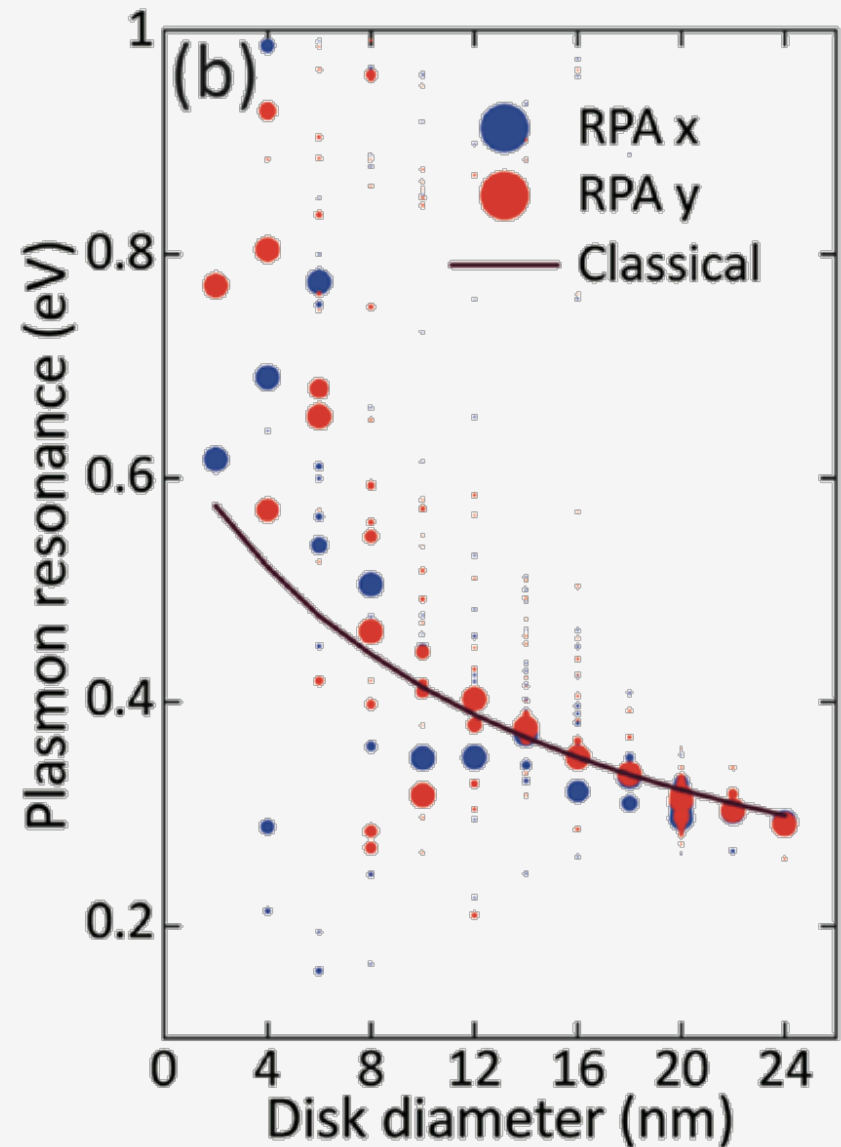
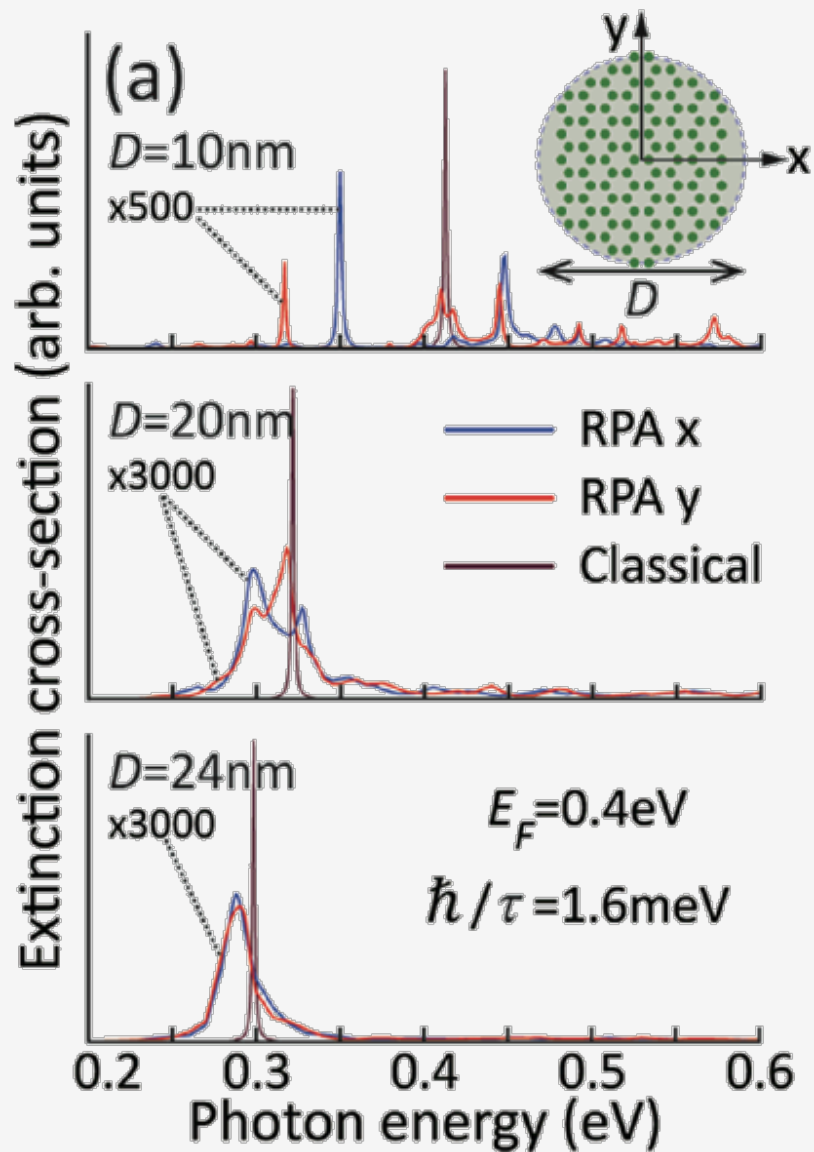


RPA response

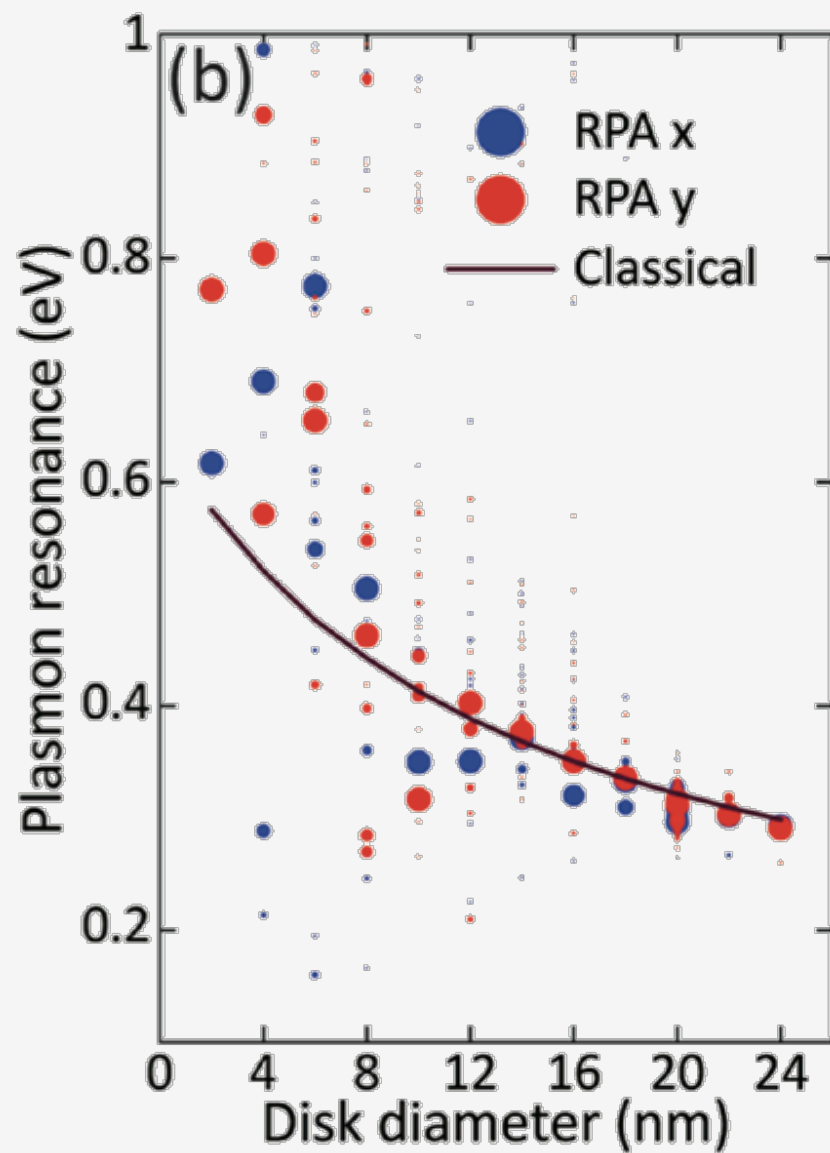
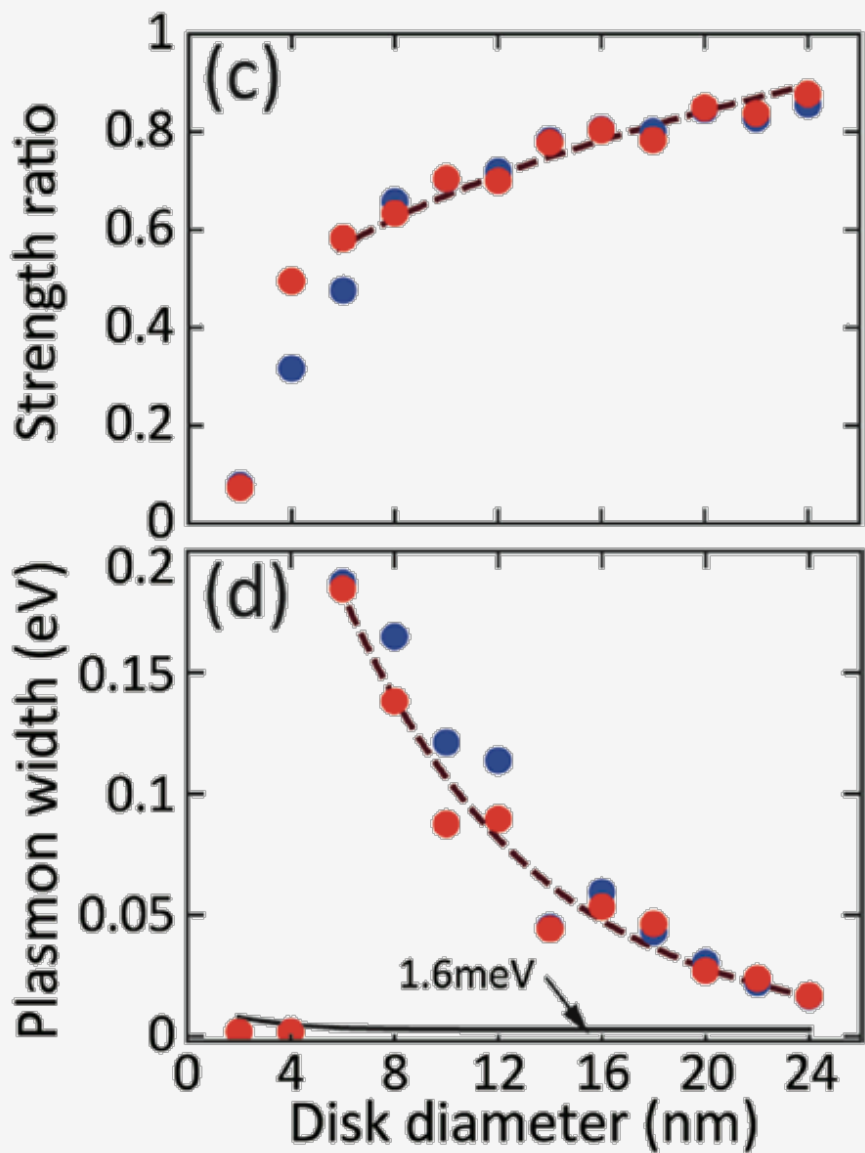
$$\chi_{ll'}^0(\omega) = \frac{2e^2}{\hbar} \sum_{jj'} (f_{j'} - f_j) \frac{a_{jl} a_{j'l'}^* a_{j'l}^* a_{j'j}}{\omega - (\varepsilon_j - \varepsilon_{j'}) + i/2\tau}$$

Thongrattanasiri *et al.*, ACS Nano (2012)

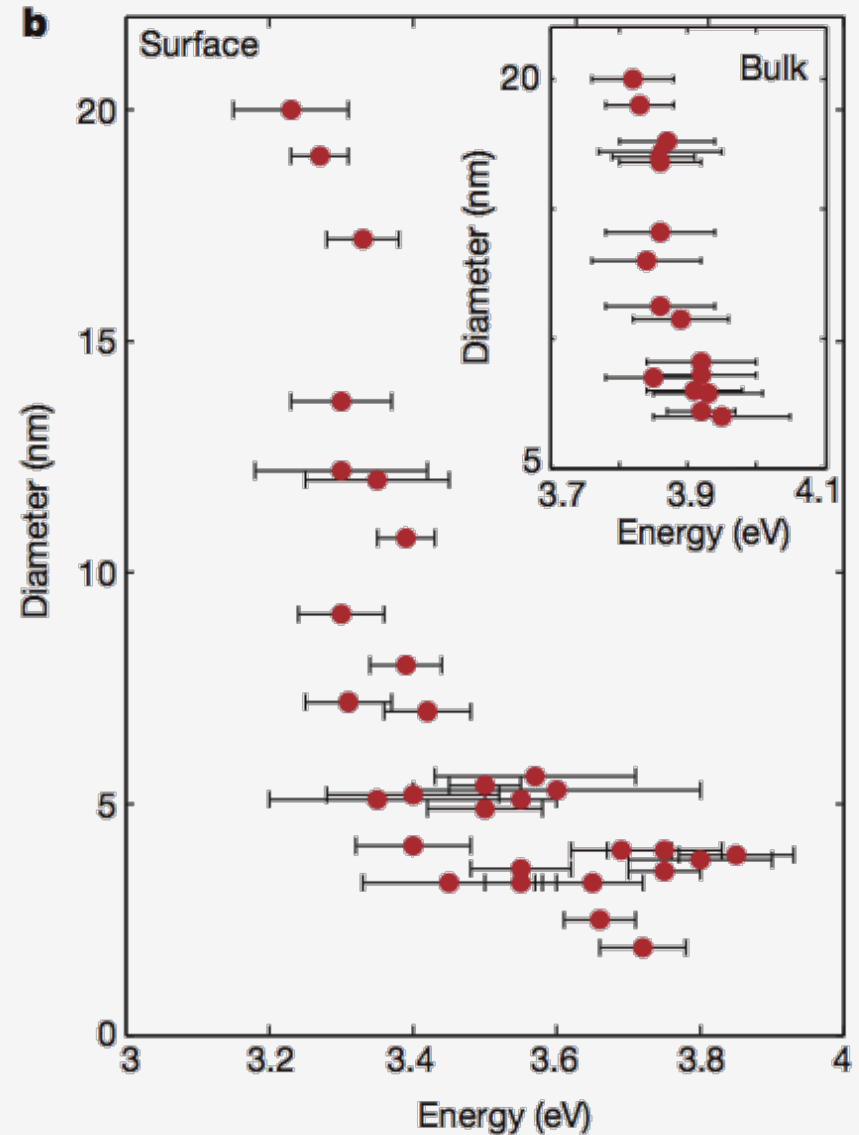
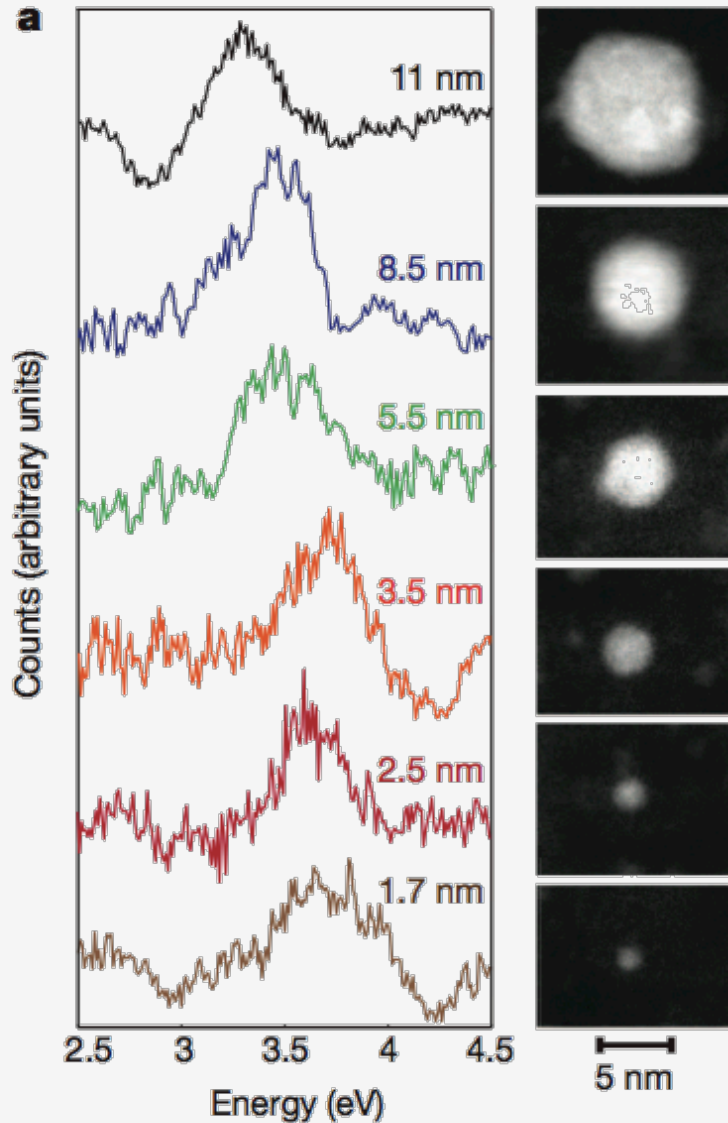




Thongrattanasiri *et al.*, ACS Nano (2012)



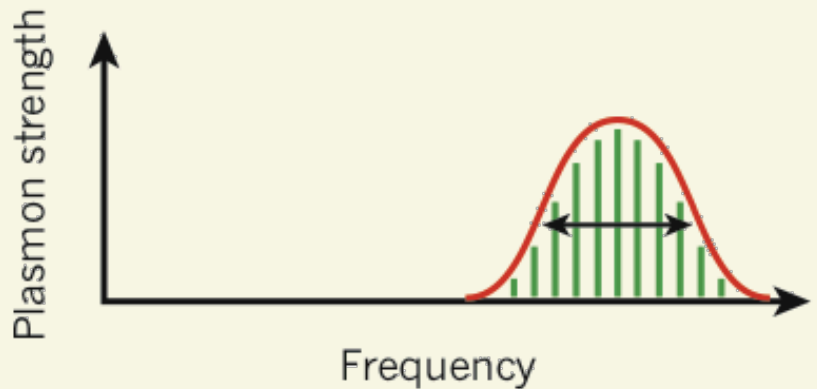
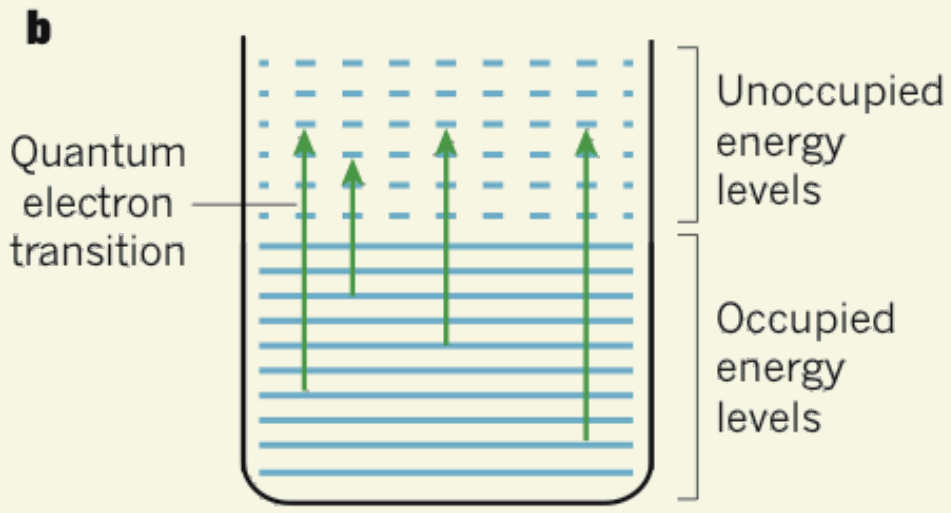
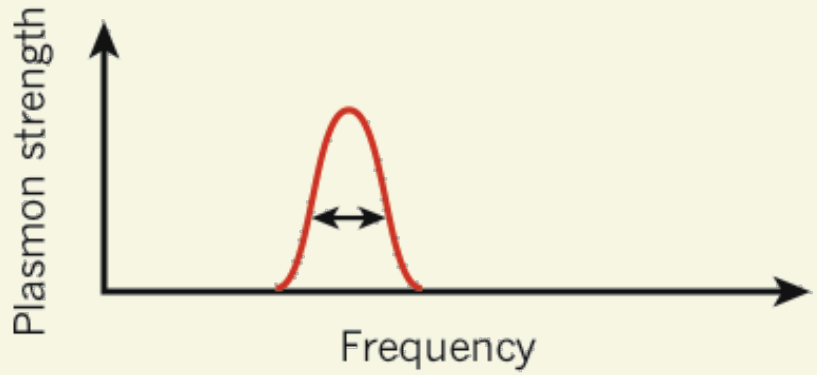
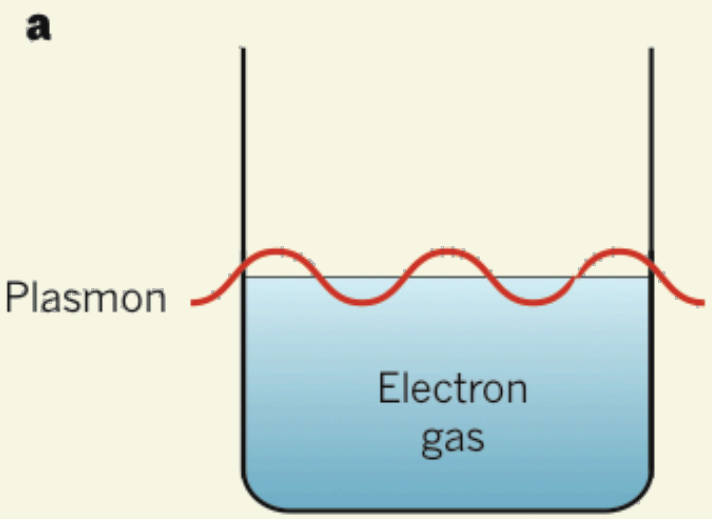
Thongrattanasiri *et al.*, ACS Nano (2012)

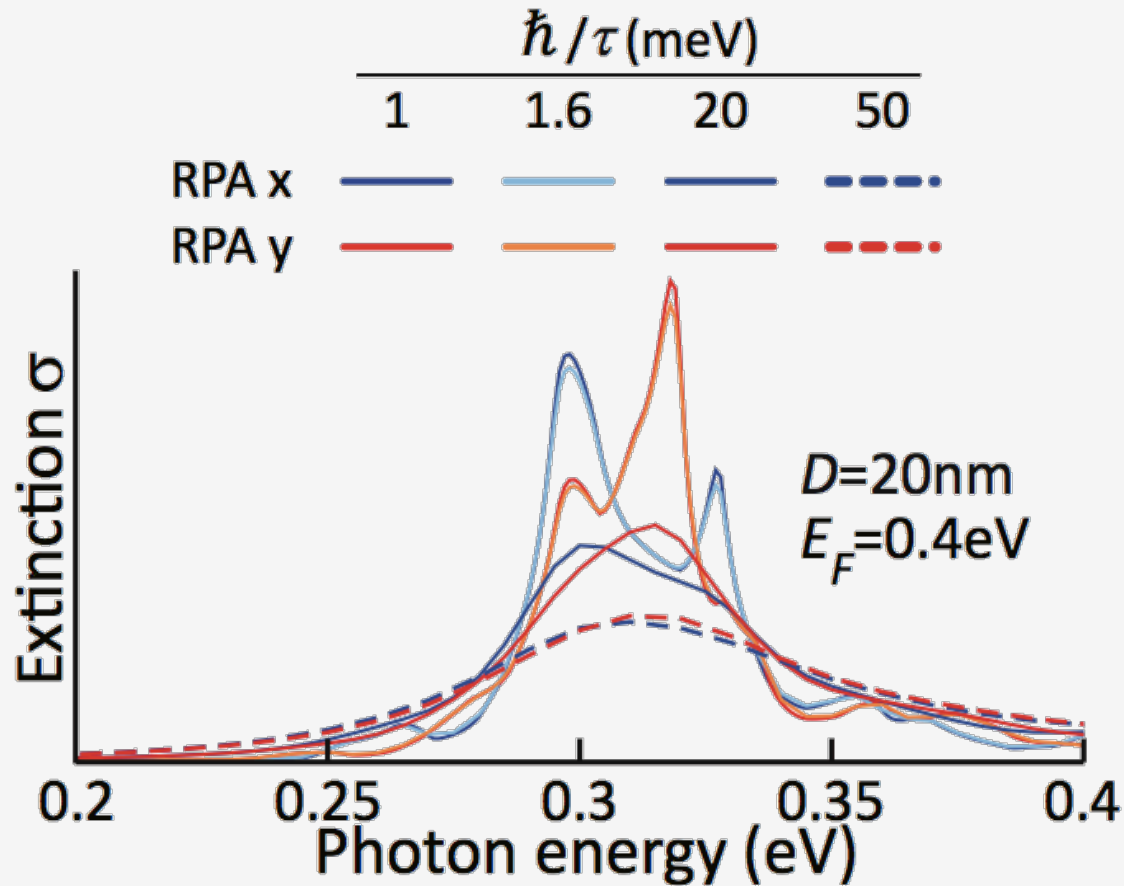


Scholl, Koh, and Dionne, *Nature* (2012)

Quantum effects in graphene plasmons

News and Views, Nature **483**, 861 (2012)



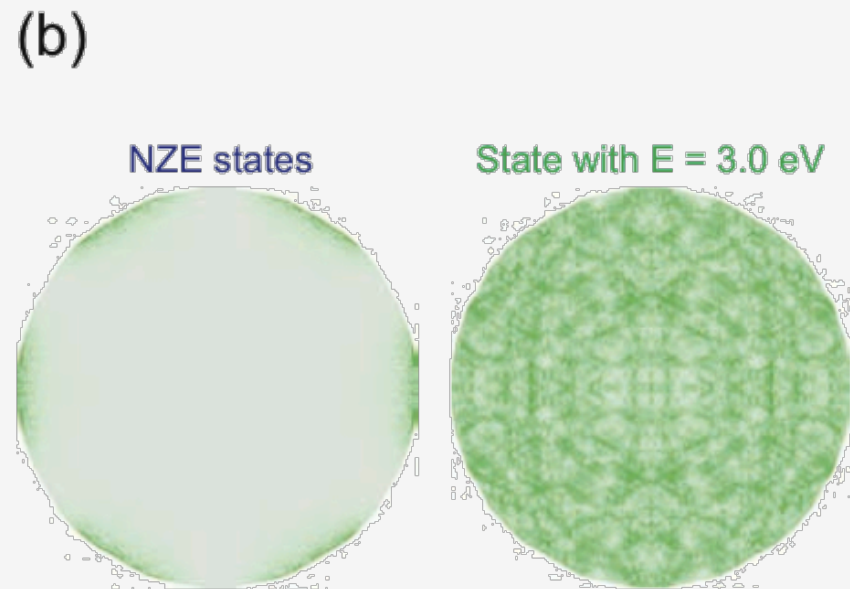
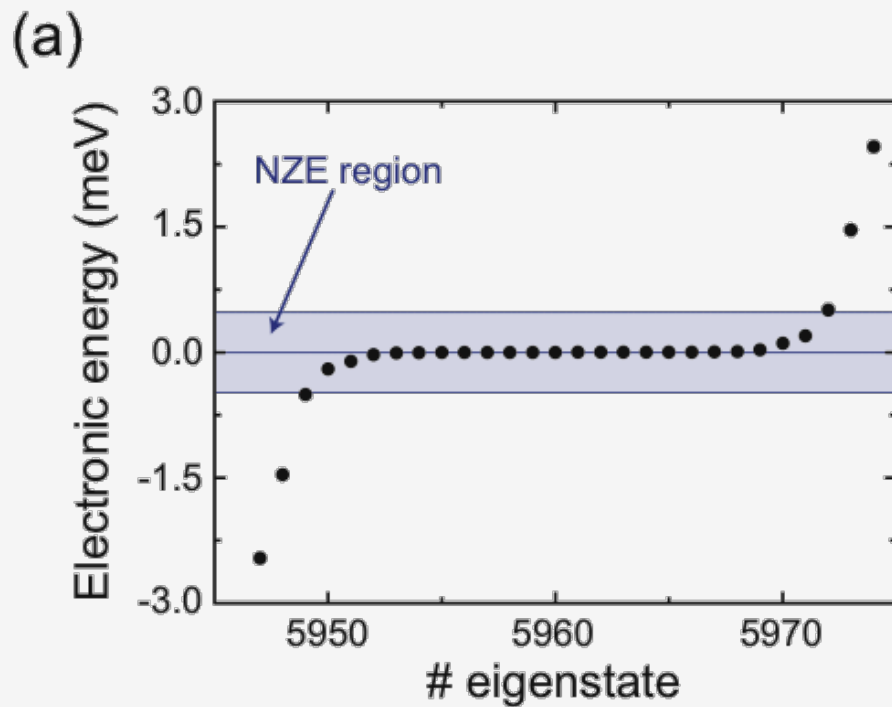


$$\chi_{ll'}^0(\omega) = \frac{2e^2}{\hbar} \sum_{jj'} (f_{j'} - f_j) \frac{a_{jl} a_{j'l}^* a_{j'l}^* a_{j'l}}{\omega - (\epsilon_j - \epsilon_{j'}) + i/2\tau}$$

Thongrattanasiri *et al.*, ACS Nano (2012)

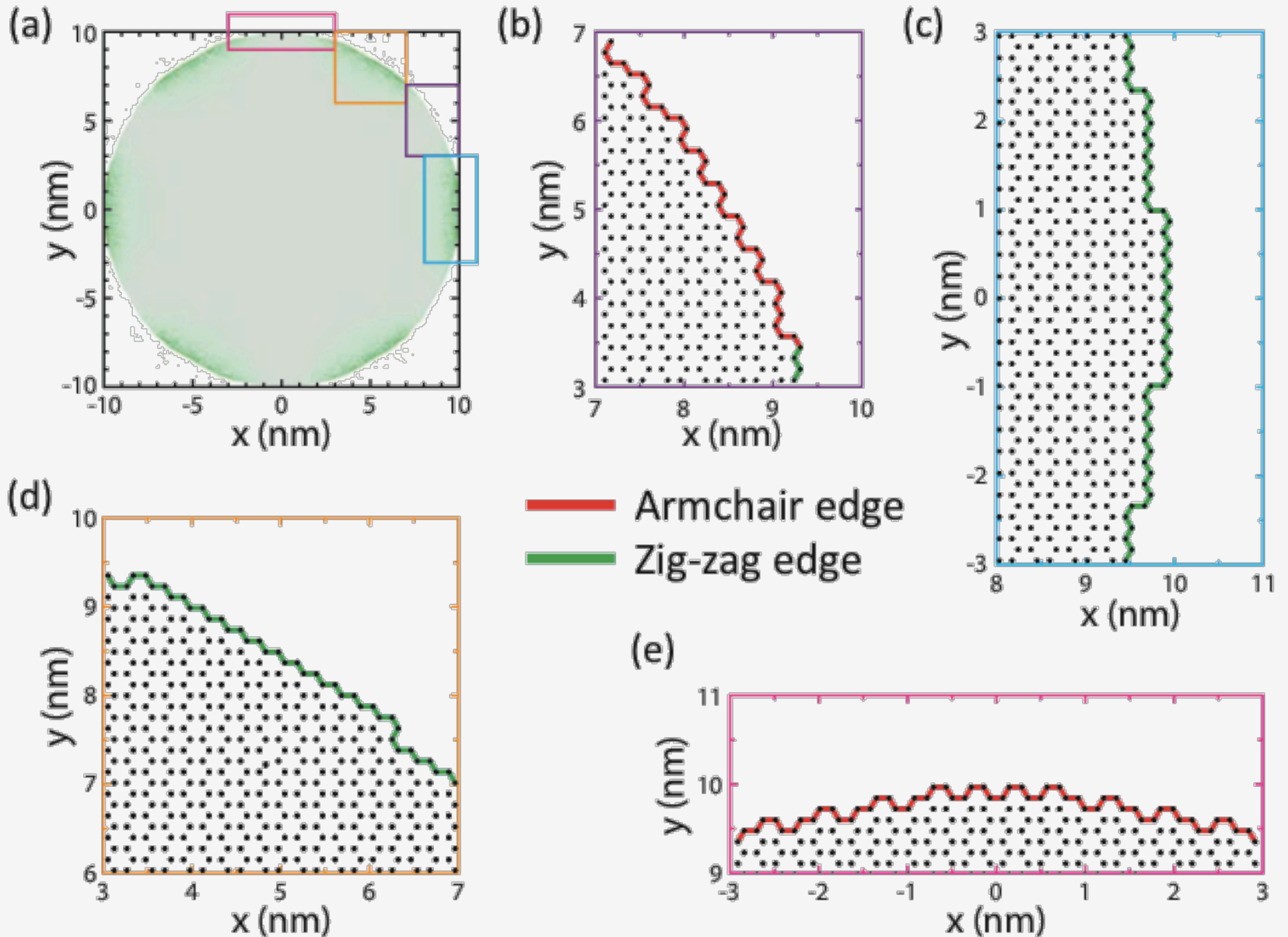
Electron state j
of energy ε_j

$$\sum_l a_{jl} |\Lambda\rangle$$

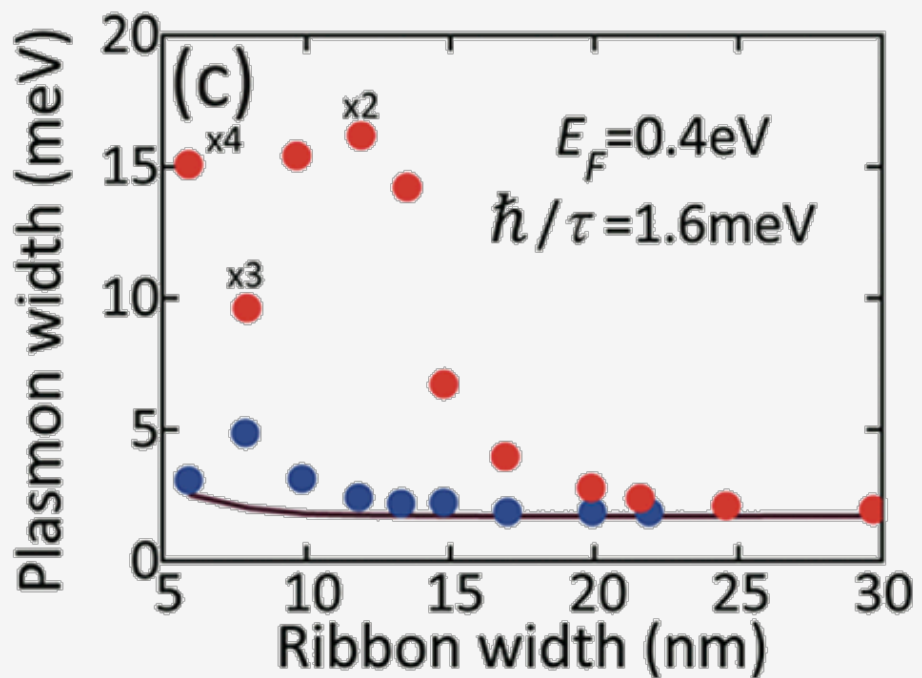
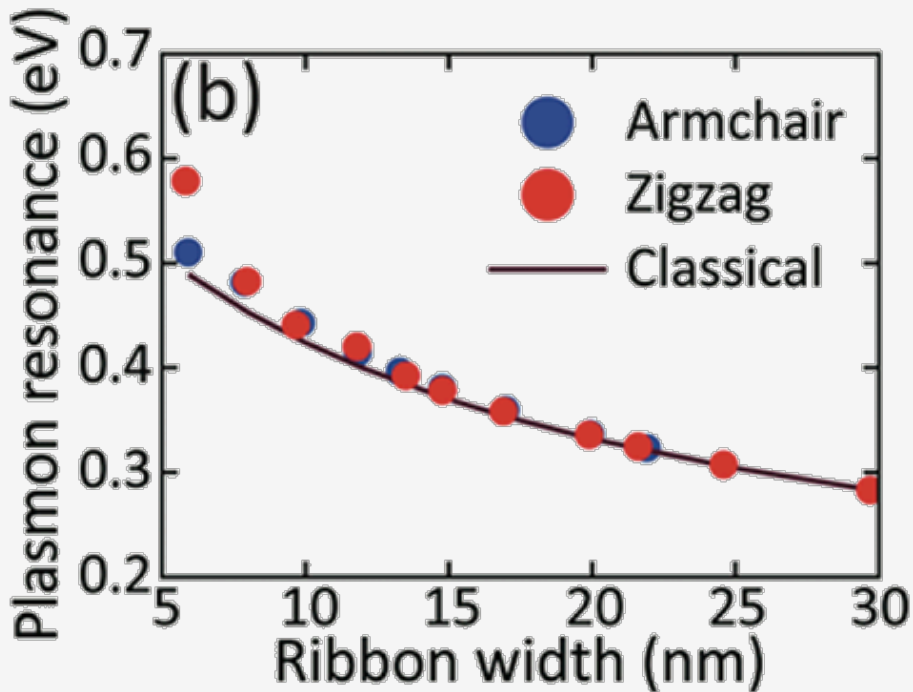
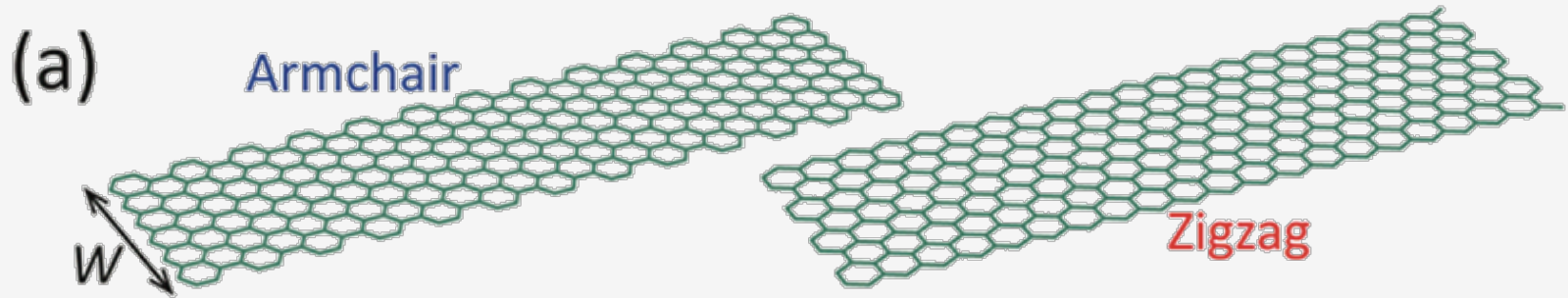


Thongrattanasiri *et al.*, ACS Nano (2012)

Quantum effects in graphene plasmons

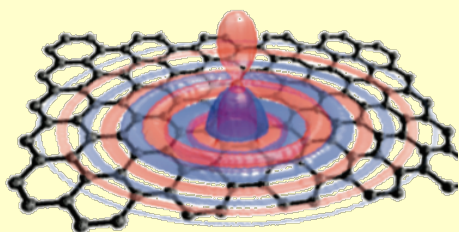


Thongrattanasiri *et al.*, ACS Nano (2012)



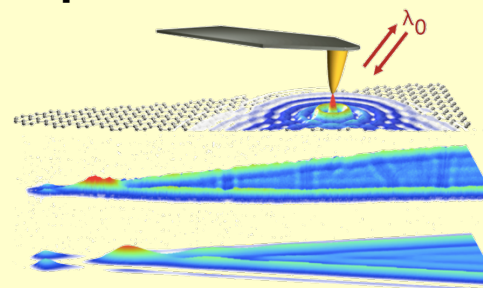
Thongrattanasiri *et al.*, ACS Nano (2012)

Strong light-matter interaction



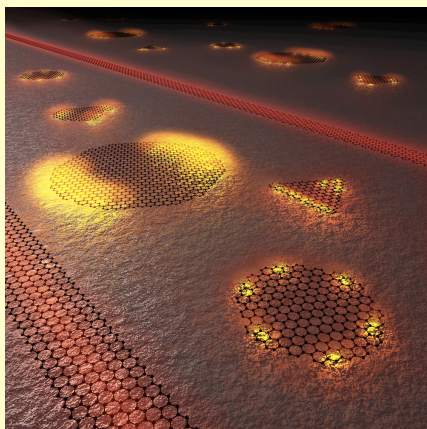
Koppens, Chang & García de Abajo, Nano Lett. (2011)

Experimental observations



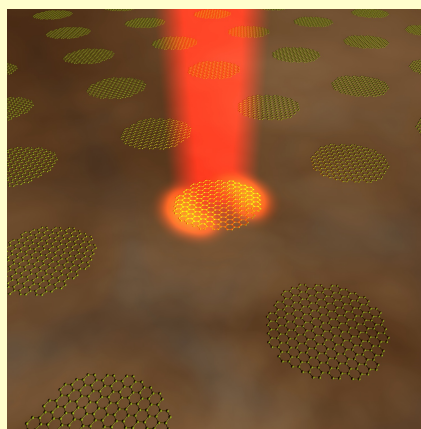
Basov's group, Nature (2012)
Koppens, Hillenbrand, García de Abajo's group, Nature (2012)
Zheyu et al, in preparation

Intrinsic quantum effects



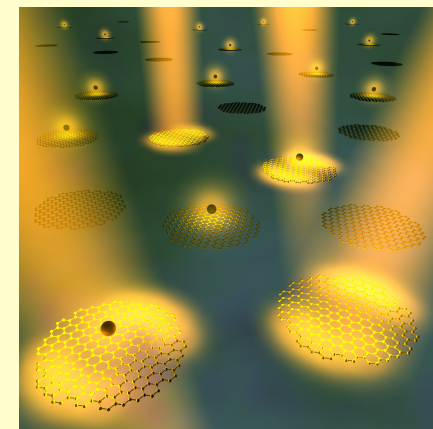
Thongrattanasiri, Manjavacas & García de Abajo, ACS NANO (2011)

Extraordinary metamaterials: Complete optical absorption



Thongrattanasiri, Koppens & García de Abajo, PRL (2011)

Quantum optics with graphene plasmons



Manjavacas, Nordlander & García de Abajo, ACS NANO (2011)